



The Mega-MUSCLES Spectral Energy Distribution of TRAPPIST-1

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Abstract

We present a 5 Å–100 μm spectral energy distribution (SED) of the ultracool dwarf star TRAPPIST-1, obtained as part of the Mega-MUSCLES Treasury Survey. The SED combines ultraviolet and blue-optical spectroscopy obtained with the Hubble Space Telescope, X-ray spectroscopy obtained with XMM-Newton, and models of the stellar photosphere, chromosphere, transition region, and corona. A new differential emission measure model of the unobserved extreme-ultraviolet spectrum is provided, improving on the Lyα–EUV relations often used to estimate the 100–911 Å flux from low-mass stars. We describe the observations and models used, as well as the recipe for combining them into an SED. We also provide a semiempirical, noise-free model of the stellar ultraviolet spectrum based on our observations for use in atmospheric modeling of the TRAPPIST-1 planets.

Unified Astronomy Thesaurus concepts: M dwarf stars (982); Exoplanet astronomy (486); Exoplanet atmospheres (487); Ultraviolet astronomy (1736); X-ray astronomy (1810); Spectral energy distribution (2129); Spectroscopy (1558); Stellar atmospheres (1584)

1. Introduction

Among the thousands of planetary systems that have been discovered over the past two and a half decades, TRAPPIST-1 is a standout case. Discovered by the Transiting Planets and PlanetIsimals Small Telescope (TRAPPIST) survey in 2015 (Gillon et al. 2016), the system is comprised of an M8 ultracool dwarf star orbited by seven planets, all of which have similar masses and radii to Earth and Venus (Gillon et al. 2017; Wang et al. 2017). The planets are almost exactly coplanar, have orbital periods ranging between 1.5 and 18.8 days, and are all in an orbital resonance with at least one other planet (Luger et al. 2017). At a distance of 12.43 ± 0.02 pc and $r_{\text{mag}} = 17.87 \pm 0.01$ (Chambers et al. 2016), the system presents a challenging but achievable target for transit spectroscopy observations of the planets, both now with the Hubble Space Telescope (HST; de Wit et al. 2016) and in the future with JWST (Barstow & Irwin 2016; Morley et al. 2017). Three or four of the planets orbit at distances where the energy received from the star is such that liquid water might persist on their surfaces. The TRAPPIST-1 system therefore offers opportunities for comparative planetology to test models of planetary habitability, biosignatures, and even, given the small orbital

separations between the planets, transfer of material, and/or life between the planets (Veras et al. 2018).

A complete evaluation of the potential habitability of the TRAPPIST-1 planets requires comprehensive knowledge of the parent star. This has proven challenging, with uncertainties remaining over, for example, the star's age (Burgasser & Mamajek 2017; Gonzales et al. 2019), activity (Vida et al. 2017), and rotation period (Roettenbacher & Kane 2017). Of particular importance, given the close proximity of the planets to the star, is the stellar magnetic activity and the resulting X-ray and ultraviolet emission. High-energy radiation can influence the retention and chemistry of planetary atmospheres as well as surface survival conditions (Miguel et al. 2015; Rugheimer et al. 2015; O'Malley-James & Kaltenegger 2017). However, TRAPPIST-1 is extremely faint at short wavelengths, making detailed characterization of the high-energy environment in the system challenging. Wheatley et al. (2017) observed TRAPPIST-1 with XMM-Newton (XMM), finding variable X-ray luminosity with intensity, similar to the modern quiescent Sun. Because of their proximity to the host stars, the planets would therefore experience XUV intensities much higher than Earth, sufficient to significantly alter their atmospheres and strip away hydrogen from water in their

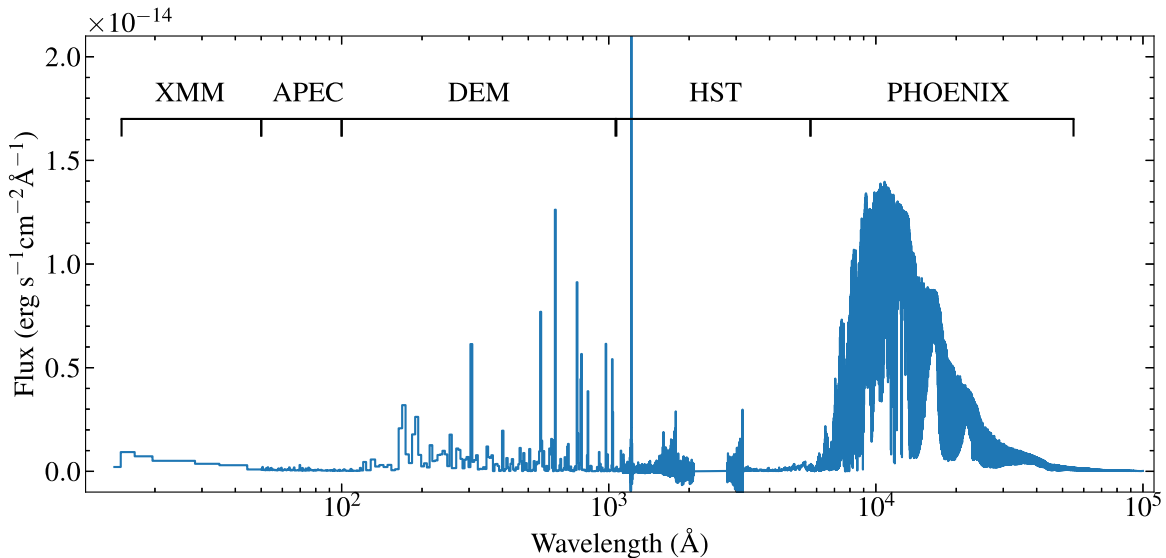


Figure 1. SED of TRAPPIST-1 with all data and models at native resolutions. The sources for each section of the spectrum are labeled above it.

atmospheres and (if present) oceans (Ribas et al. 2016; Airapetian et al. 2017). Bourrier et al. (2017a, 2017b) obtained time-series observations of the 1215.67 Å Ly α hydrogen emission line with HST, finding that it evolved over a three-month timescale but with no evidence for hydrogen (and thus water) escape from TRAPPIST-1 c, which transited during their observations. Peacock et al. (2019) used the PHOENIX stellar atmosphere code to model the chromosphere and transition region of TRAPPIST-1, scaling it to the Bourrier et al. (2017b) Ly α measurement and to distance-adjusted GALEX observations of stars with a similar spectral type. They found that the flux emitted between 100–912 Å varies by an order of magnitude depending on which calibrator was used. These studies demonstrate that accurately accounting for the effects of high-energy radiation on the TRAPPIST-1 planets requires spectroscopic observations at all accessible wavelengths.

Mega-MUSCLES (Measurements of the Ultraviolet Spectral Characteristics of Low-Mass Exoplanetary Systems) is an HST Treasury program obtaining 5 Å–100 μ m spectral energy distributions (SEDs) of a representative sample of 12 M dwarfs, covering a wide range of stellar mass, age, and planetary system architecture, and extending the original 11 star MUSCLES program (France et al. 2016; Loyd et al. 2016; Youngblood et al. 2016) to stars with lower masses, higher activity, and/or faster rotation rates. Here we present the Mega-MUSCLES SED of TRAPPIST-1 (Figure 1), composed of ultraviolet and X-ray spectroscopy with HST and XMM, along with state-of-the-art model spectra. We discuss the changes made to the data processing and stellar emission modeling implemented for Mega-MUSCLES compared with MUSCLES, present a semiempirical model for use in model atmosphere simulations, and compare the observed SED to the Peacock et al. (2019) models.

2. Observations

We observed TRAPPIST-1 with the Cosmic Origins Spectrograph (COS; Green et al. 2012) and the Space Telescope Imaging Spectrograph (STIS; Woodgate et al. 1998) on board the Hubble Space Telescope (HST) on 2017 December 15, 2018 December 8–12, and 2019 June–8,

for a total exposure time of 36,379 s. The COS gratings used were G160M (8608 s), G130M (12,404 s), and G230L (2731 s), and the STIS gratings were G430L (1795 s) and G140M (10,841 s). Combined, these spectra cover the wavelength range 1130–5700 Å except for a gap between 2080 and 2790 Å, which is not covered by the COS NUV detector and is too faint for STIS NUV observations. The HST observations are summarized in Table 1. With the exception of the STIS/G430L exposure, the observations were obtained using photon-counting detectors in TIME-TAG mode. We extracted light curves from each spectrum to search for and potentially remove contributions from flares or other stellar activity, but found no significant variation.

For the COS G160M and G130M observations, variations in target position in the aperture between each orbit induce slight differences in wavelength calibration. To remove this effect we cross-correlated known emission lines in each x1d spectrum to shift each spectrum onto a single wavelength scale before coadding. Doing so provides a small increase in signal-to-noise ratio (S/N) and resolution compared with the x1dsum files produced by the CALCOS pipeline.

Four spectra were obtained with the STIS G140M grating, covering the Ly α hydrogen line with a spectral range of 1195–1249 Å and $R \sim 10,000$. The automated reduction pipeline failed to identify the spectral trace, so we used the fit images to visually identify the spectrum in each exposure, before re-extracting the spectrum with the STISTOOLS x1d routine, fixing the “a2center” keyword to the identified spectrum position. In one spectrum, the trace could not be visually identified. Ly α line fluxes in the spectra obtained immediately before and after the nondetection were similar, so the nondetection is unlikely to be due to intrinsic variability and is probably an instrumental effect such as inaccurate slit positioning. The nondetection was therefore discarded, and the remaining three spectra coadded into the final G140M spectrum used hereafter.

TRAPPIST-1 has been observed multiple additional times with the STIS G140M grating both prior to and since our observations (Bourrier et al. 2017a, 2017b). We initially intended to combine all of the available spectra, but we found the detected Ly α flux to be variable, with the final coadded

Table 1
Summary of Observations

Date	Instrument	Grating	Central Wavelength (Å)	Start Time (UT)	Total Exposure Time (s)	Data Set
HST						
2017-12-15	COS	G230L	2950.000	07:19:18	2731	LDLM42010
2018-12-08	COS	G160M	1577.000	06:53:42	1658	LDLM39010
2018-12-08	COS	G160M	1611.000	08:07:13	5352	LDLM39020
2018-12-09	STIS	G430L	4300.000	01:49:27	1795	ODLM41010
2018-12-09	STIS	G140M	1222.000	04:46:12	2710	ODLM41030
2018-12-09	STIS	G140M	1222.000	06:21:32	2710	ODLM41040
2018-12-09	STIS	G140M	1222.000	07:56:53	2710	ODLM41050
2018-12-10	COS	G130M	1291.000	03:22:34	12403	LDLM40010
2019-06-07	COS	G160M	1577.000	19:54:25	1598	LDLMZ6010
XMM						
2018-12-10	EPIC	...	...	10:42:57	24600	0810210101

Note. Data set numbers are given for retrieval from MAST (<https://archive.stsci.edu/hst/>) or the XMM-Newton Science Archive (<http://nxs.esac.esa.int/nxsa-web/#home>) for HST and XMM, respectively.

spectrum highly dependent on which subspectra were chosen for inclusion. Epoch-to-epoch changes are somewhat beyond the scope of this paper, and we do not wish to preempt the teams leading the ongoing G140M observations, so we decided to use only the data obtained as part of the Mega-MUSCLES program. This has the advantage of ensuring that all of the data in the SED is contemporaneous, at the cost of improved S/N for the Ly α measurement. We will reevaluate this decision when a full analysis of the G140M data is available and update the Mega-MUSCLES High-Level Science Products accordingly.

We further observed TRAPPIST-1 with XMM-Newton (XMM) using the EPIC instrument with thin filters for 23 ks on 2018 December 10, overlapping in time with the COS G130M observations (Figure 2). TRAPPIST-1 was detected with an average count rate of 3 counts ks⁻¹, less than the average count rate of ≈ 10 –20 counts ks⁻¹ found by Wheatley et al. (2017). Analysis of all available archival data suggested that the Wheatley et al. (2017) result was dominated by a strong flare(s), thus we take our observation to represent the “typical” or quiescent X-ray flux.

3. SED

Figure 1 shows the full SED and the wavelength ranges covered by the different observational and model sources. The methods used to produce the SED for TRAPPIST-1 and the rest of the Mega-MUSCLES sample are based on those used for the MUSCLES survey (Lloyd et al. 2016), but with multiple differences reflecting developments in modeling techniques since MUSCLES and, in the case of TRAPPIST-1, the faintness of the target in the ultraviolet. Below we describe in full the production of the TRAPPIST-1 SED. When discussing specific wavelength regions, we follow the definitions used for the MUSCLES program by France et al. (2016): X-ray: 5–100 Å; EUV: 100–911 Å; FUV: 911–1700 Å; NUV: 1700–3200 Å and optical/IR: 3200 Å–100.0 μ m.

3.1. Variability

Late-type M dwarfs are notably active stars, with regular flares and high photometric variability (Paudel et al. 2018). Even stars that appear photometrically quiet in optical surveys have been shown to be active in the ultraviolet (Lloyd et al. 2018; France et al. 2020). Vida et al. (2017) identified 42

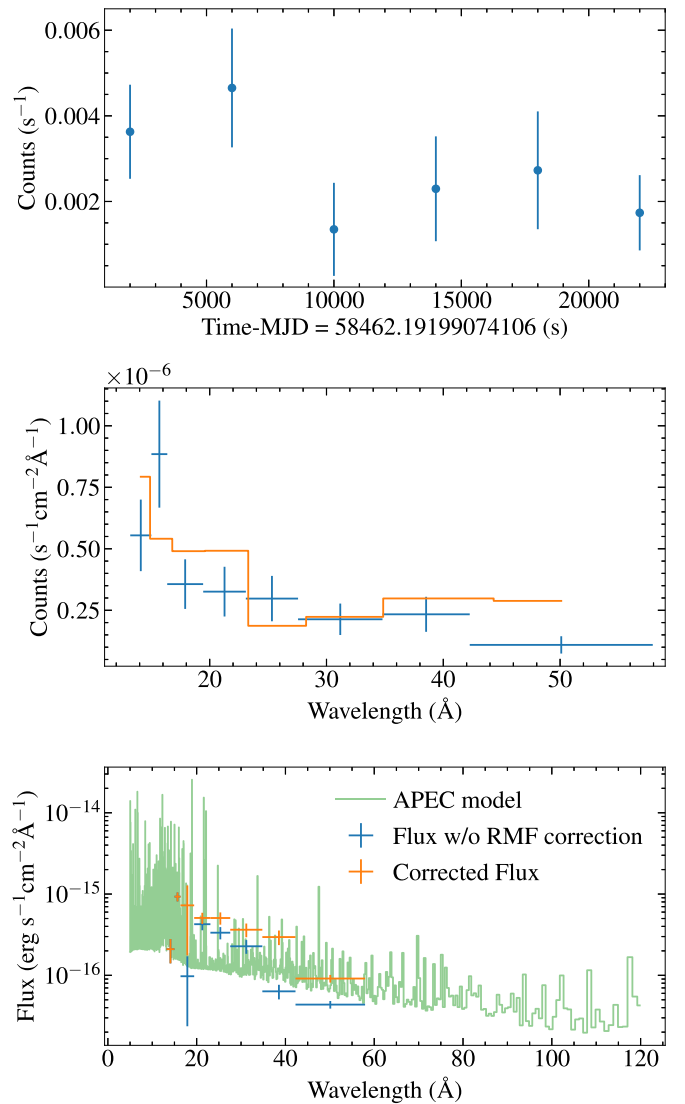


Figure 2. Top: X-ray light curve of TRAPPIST-1 observed on 2018 December 10. Middle: XMM spectrum of TRAPPIST-1 (blue) compared with the APEC model (orange). The model has been binned into the same bins as the spectrum. Bottom: full APEC model used to correct the X-ray spectrum and provide the 50–120 Å range of the SED, along with the raw and RMF-corrected flux conversions described in the text.

white-light flaring events in the *K2* light curve of TRAPPIST-1, approximately 0.5 day^{-1} , although Gillon et al. (2017) only detected two flares in 20 days of infrared Spitzer observations taken at a different epoch, indicating that the flare rate is time and/or wavelength dependent. We must therefore consider activity when assessing the validity of the SED presented here. We detected no flares in any of our ultraviolet and X-ray observations and thus consider the SED to represent TRAPPIST-1 in a “quiescent,” non-flaring state. Flares have been detected in previous X-ray observations by Wheatley et al. (2017), and potentially in broadband NUV Swift photometry by Becker et al. (2020). Determining how representative the ultraviolet spectra presented here are of the time-averaged emission of TRAPPIST-1 will require sustained monitoring at multiple wavelengths, a significant challenge considering the limitations of current ultraviolet observatories. Therefore, although we are confident that the Mega-MUSCLES SED is an accurate measurement of the quiescent emission of TRAPPIST-1, efforts to model the atmospheres of the TRAPPIST-1 planets should acknowledge that we have poor constraints on how that emission changes during a flare, and how often such flares occur.

3.2. X-Ray

The XMM spectrum was fit using the XSPEC v12.10 package (Arnaud 1996) with models generated using the Astrophysical Plasma Emission Code v3.0.9 (APEC; Smith et al. 2001; Foster et al. 2012), with the data binned in variable energy bin widths such that each bin contained the same number of counts. The data were fit using a two-temperature model ($kT = 0.2, 0.4$) and a metallicity of 0.4 solar. Due to the low number of photons, the only free parameters used were the normalization of the two temperature components. We used the standard XSPEC background subtractions, although we found that ignoring the background changed the derived flux by less than 0.1 dex.

Ordinarily, the data could be converted from counts N into physical units F by dividing by the effective area A of the instrument as a function of energy E , i.e.,

$$F(E) \propto \frac{N(E)}{A(E)T_{\text{exp}}}. \quad (1)$$

However, for cool spectra such as TRAPPIST-1, issues arise at the low-energy boundary (0.2 keV), because more photons are discarded when detected below the energy cutoff than photons being recorded at higher than nominal energies due to the highly asymmetric energy response matrix (RMF). A priori, we do not know the magnitude of this effect, but with a reasonable model spectrum, we can estimate the energy-dependent fraction of photons being lost to this effect. We use this (mildly) model-dependent knowledge to correct Equation (1) to

$$F(E) \propto \frac{N(E)}{A(E)T_{\text{exp}} C(E)}, \quad (2)$$

where the correction C captures the asymmetry of the spectrum and the RMF. For a symmetric RMF and a flat spectrum, $C(E)$ would be 1.0. Consider a spectrum consisting of a single emission line at 0.3 keV and a low-energy cutoff for the spectrum of 0.2 keV: the recorded photon rate will be lower than expected, because a sizable fraction of the photons will

have reconstructed energies below 0.2 keV and will thus be missing from the spectrum. This is not an effective area but an RMF effect. In fact, the shape of the XMM-Newton RMF at low energies (0.1–0.3 keV) is such that more photons will be detected below their nominal energy than above, so that even for a flat spectrum, one would need to correct for the RMF effect. In essence, this procedure ensures that the corrected spectrum reproduces the model flux, which would not be the case when using Equation (1). However, we caution that this “corrected” spectrum still includes RMF effects and is meant for illustrative purposes in physical units while the best-fit model spectrum captures the emitted spectrum much better.

3.3. Extreme-ultraviolet and DEM Modeling

The most significant departure from the MUSCLES procedure is in the extreme-ultraviolet (EUV), where we have replaced the Linsky et al. (2014) empirical scaling relations with a differential emission measurement (DEM) model, which estimates the chromospheric, transition region, and coronal emission based on the strength of the detected lines in the FUV spectrum and the X-ray flux. A model spectrum is required as the region spanning 120–1100 Å is unobservable, both physically due to absorption from interstellar hydrogen between 400 and 900 Å, and a lack of currently operating instruments that can observe the ranges 120–400 Å and 900–1100 Å (HST/COS and Chandra do have some modes that extend down to 900 Å and up to 175 Å, respectively, but these were not sensitive enough to be practical for this program). The DEM model is described in detail in Section 5.

3.4. Far- and Near-ultraviolet

3.4.1. Ly α Reconstruction

The H I 1215.67 Å Ly α line is clearly visible in our coadded STIS G140M spectrum, but is heavily affected by both interstellar absorption and terrestrial airglow. We reconstructed the full Ly α profile with a technique derived from Youngblood et al. (2016), where we simultaneously fit a model of the interstellar absorption and a model of the intrinsic stellar emission with a Markov Chain Monte Carlo method (emcee; Foreman-Mackey et al. 2013). In order to benefit from the increased S/N offered by spectra taken before the Mega-MUSCLES observation while simultaneously accounting for possible intrinsic stellar variability, we fit both the new Mega-MUSCLES observation and the observation from Bourrier et al. (2017b) simultaneously. We use a Voigt profile for the stellar emission as well as for the ISM absorption, but require the ISM absorption to be the same for both observations. To account for wavelength solution errors between the two observations, we do not fix the radial velocities to the same value, but rather require the radial velocity offset between the emission and ISM absorption to be the same between the two observations. Based on the measured stellar radial velocity $-56.3 \pm 1 \text{ km s}^{-1}$ (Reiners & Basri 2009a) and the predicted ISM radial velocity along the line of sight towards TRAPPIST-1 ($-1.25 \pm 1.37 \text{ km s}^{-1}$; Redfield & Linsky 2008), we give the radial velocity offset parameter a Gaussian prior with mean $+55.05 \pm 1.70 \text{ km s}^{-1}$, obtained by adding the stellar and ISM radial velocities and adding their uncertainties in quadrature. We fix the Doppler b value of the ISM absorption profile to 11.5 km s^{-1} and $D/H = 1.5 \times 10^{-5}$, both standard values for

TRAPPIST-1

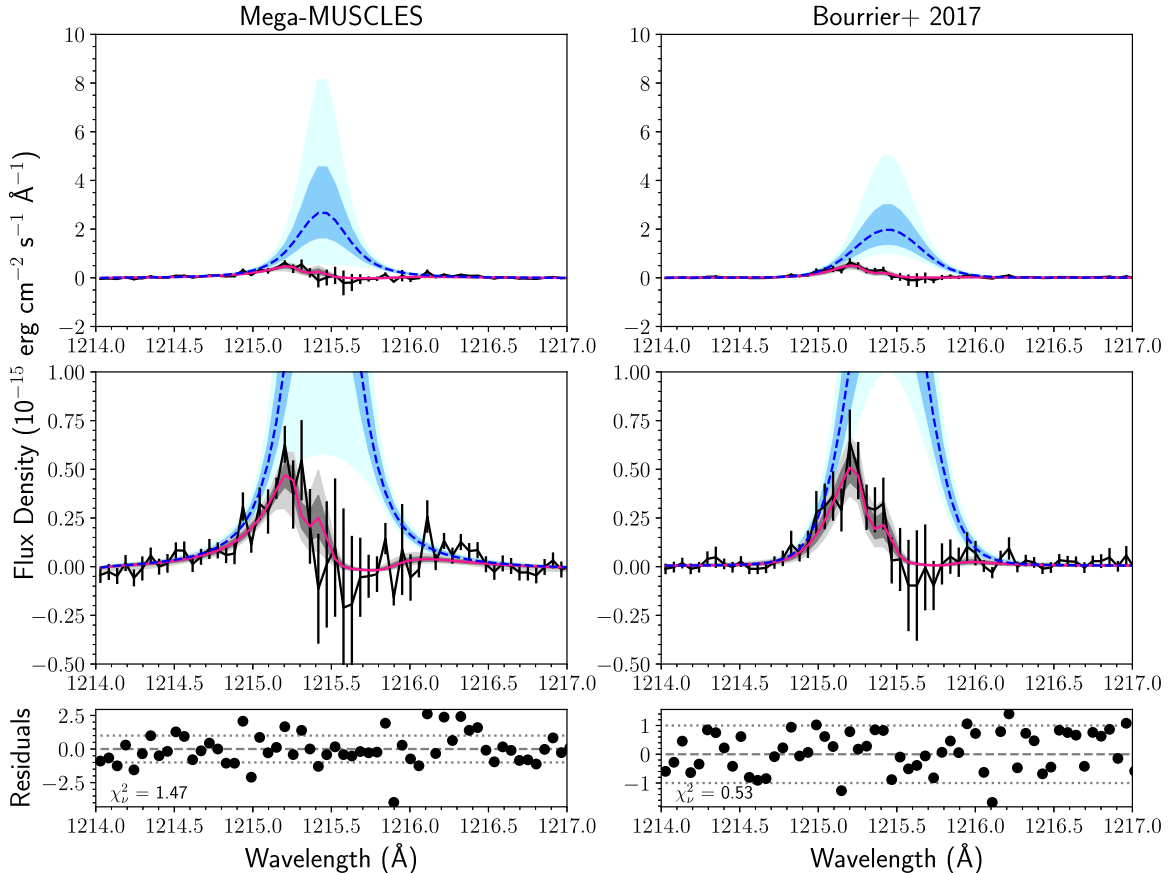


Figure 3. Best-fit models and intrinsic Ly α profiles are shown from a simultaneous fit to the new spectrum presented in this work and the spectrum presented in Bourrier et al. (2017b). The ISM properties were forced to be the same for the two spectra, but the intrinsic emission profiles were allowed to vary to account for either intrinsic stellar variability or differences in the flux calibration between the two data sets. The solid pink lines show the best fits to the data (black lines with error bars), with the dark and light shaded regions corresponding to the 68% and 95% confidence intervals, respectively. The dashed blue lines show the intrinsic emission profiles corresponding to the best-fit models, and the dark blue and light blue shaded regions represent the 68% and 95% confidence intervals, respectively. The bottom panels show the residuals of the fit, or the best-fit models subtracted from the data and divided by the data uncertainty. The reduced chi-squared values are printed in the same panels.

the local ISM (Wood 2004). All other parameters were varied with uniform priors.

Figure 3 shows the spectrum and reconstructed line profile. We report the median and the 68 per cent confidence interval as our best-fit values and 1σ error bars. We found an integrated flux of $F_{\text{Ly}\alpha} = (1.40^{+0.60}_{-0.36}) \times 10^{-14} \text{ erg s}^{-1} \text{ cm}^{-2}$ for the Mega-MUSCLES observation, and $F_{\text{Ly}\alpha} = (1.09^{+0.40}_{-0.27}) \times 10^{-14} \text{ erg s}^{-1} \text{ cm}^{-2}$ for the Bourrier et al. (2017b) observation. This value is larger but within 1σ of the flux $((8.18^{+1.64}_{-3.27}) \times 10^{-15} \text{ erg s}^{-1} \text{ cm}^{-2})$ found by Bourrier et al. (2017b) and is explained by our fit’s larger column density ($\log_{10} N(\text{H I}) = 18.4 \pm 0.1$, compared to 18.3 ± 0.2 from Bourrier et al. 2017b). However, these column densities and reconstructed fluxes are consistent with the result from Bourrier et al. (2017a) within 1σ , and we do not detect any statistically significant intrinsic stellar variability between the two observations.

3.4.2. COS Spectra

The COS FUV spectrum was contaminated by airglow from Ly α and O I over the wavelength ranges 1214–1217 Å and 1301–1307 Å respectively. Both ranges were removed and

replaced by the reconstructed Ly α profile in the first case and by a polynomial fit to the spectrum on either side in the second. As we therefore have no information at these wavelengths, for the rest of this paper we assume that there is zero flux from O I at TRAPPIST-1.

The COS NUV observations covered the wavelength ranges 1700–2100 Å and 2800–3200 Å leaving a 700 Å gap. This gap is partially covered by a second-order spectrum spanning 1950–2150 Å, but the signal was so weak that we chose not to include it. The gap was filled with a polynomial fit to the two wavelength regions, with the range 2790–2805 Å masked out to remove contributions from the Mg II 2800 Å lines.

3.5. Optical to Infrared

3.5.1. STIS G430L

Figure 4 shows the STIS G430L spectrum of TRAPPIST-1. The spectrum was obtained with a single exposure on a CCD detector and thus has a slightly different treatment from the rest of the HST data obtained with photon-counting detectors. Although the spectrum in principle covers the wavelength range 2900–5700 Å, in practice the combination of low detector throughput and decreasing flux from the target at

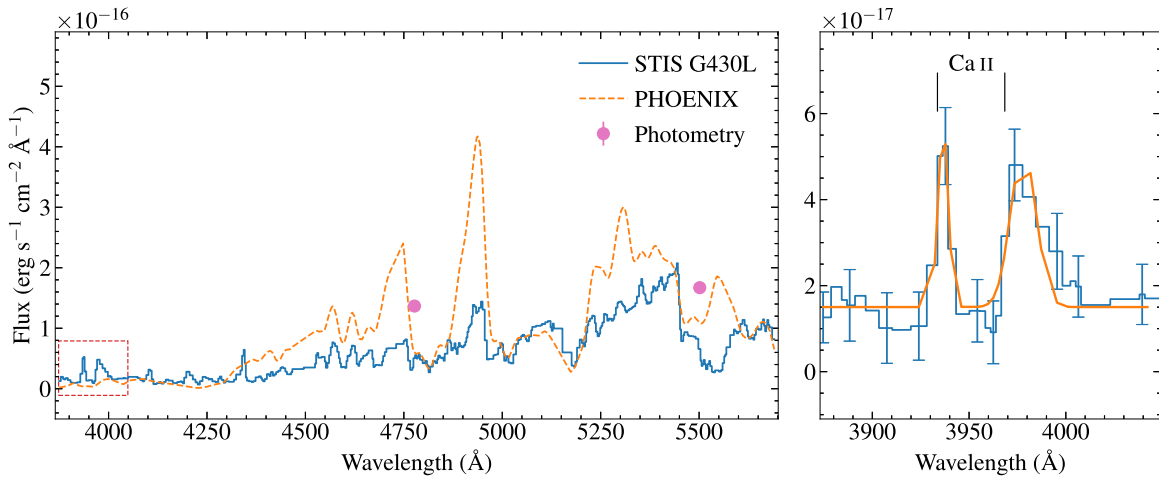


Figure 4. Left: STIS G430L spectrum (blue) compared with Pan-STARSS *g* and Gaia *B_p* photometry (Gonzales et al. 2019) and the PHOENIX model (orange). The spectrum has been smoothed by a factor 2 for clarity, and the PHOENIX model convolved to the resolution of the observed spectrum. Right: enlarged view of the region in the red box showing Ca II H&K emission lines, along with the Gaussian fits (orange) used to measure the emission line fluxes (Table 2).

short wavelengths results in an effective nondetection of the blue end of this range, with almost half of the pixel bins containing negative flux values. We therefore removed all points where the mean flux/flux error ratio of the 30 surrounding bins was less than 1. This results in a cutoff at 3872 Å, similar to the fixed 3850 Å cutoff used by Loyd et al. (2016) for the MUSCLES SEDs.

Loyd et al. (2016) also found that the STIS G430L spectra obtained for the MUSCLES program had systematically lower integrated fluxes than photometric measurements of their respective stars, which they suggested was due to imperfect alignment on the STIS slit. In an attempt to correct this, the Mega-MUSCLES G430L spectra were obtained with the slit width set to $0''.2$, twice that used for MUSCLES. To validate that the fluxes of the Mega-MUSCLES G430L spectra are correct, we have obtained extensive flux-calibrated spectroscopy with various ground-based instruments, which are in good agreement with Mega-MUSCLES spectra for the same stars (Mega-MUSCLES Collaboration 2021, in preparation). However, no comparison spectra are available for TRAPPIST-1 itself, so we have additionally compared the G430L spectra with available photometry. Of the filters for which photometric data are available, the *g* and *B* filters have bandpasses that fall nearly entirely in the G430L range. We retrieved *B*, SDSS *g*, and PanSTARS *g* magnitudes from VizieR and compared them with synthetic photometry calculated by integrating the spectrum over the respective bandpasses. We found that the spectrum underpredicted the photometry by ratios of 0.5 ± 0.1 , 0.85 ± 0.02 , and 0.6 ± 0.01 for *B*, SDSS *g*, and PanSTARS *g*, respectively. Given that the differences are small and that photometry at short optical wavelengths is more likely to be affected by stellar activity than at longer wavelengths (Paudel et al. 2018), we chose not to scale the spectra based on the photometry. An additional test is provided by the PHOENIX model, described in more detail below. As Figure 4 shows, the PHOENIX model is in good agreement with the TRAPPIST-1 data at regions of relatively low flux, but over-predicts the regions of higher flux. We speculate that this is due to incomplete treatment of opacity at these wavelengths in the model. Lançon et al. (2020) reported similar discrepancies between PHOENIX models and blue spectra of cool stars. For

Table 2
Integrated Fluxes for Ultraviolet Emission Line Lists Compiled by Linsky (2017) and Peacock et al. (2019)

Species	λ_{rest} (Å)	Flux ($10^{-18} \text{ erg s}^{-1} \text{ cm}^{-2}$)
C III ^m	1174	101 ± 27
Si III	1206.51	25.0 ± 6.3
H I (Ly α)	1215.67	$1.40^{+0.60}_{-0.36} \times 10^4$
N V	1238.821	36 ± 6.1
N V	1242.804	16 ± 10.0
Si II	1264.73	≤ 1.1
Si II	1298.95	4.4 ± 3.5
Si II	1309.28	5.7 ± 2.6
C II	1334.532	30 ± 20
C II	1335.708	55 ± 19
Cl I	1351.657	≤ 3.1
[O I]	1355.598	≤ 7.6
O V	1371.292	≤ 7.1
Fe II	1391.08	≤ 2.5
Si IV	1393.755	25 ± 19
[O IV]	1401.156	≤ 1.8
Si IV	1402.77	14.0 ± 7.6
Si II	1526.71	≤ 3.2
Si II	1533.43	≤ 3.2
C IV	1548.195	100 ± 46
C IV	1550.77	53 ± 17
C I	1561.0	≤ 6.2
He II ^m	1640	≤ 10.1
[O III]	1666.153	≤ 21.5
Al II	1670.787	≤ 79
Al I	1766.39	≤ 44
Mg II	2795.523	620 ± 29
Mg II	2802.697	≤ 340
Ca II	3933.6614	170 ± 50
Ca II	3968.4673	$\leq 350^*$

Note. Lines at wavelengths not covered by our observations were omitted.
*Detected, but blended with H ϵ . ^m Multiplet.

our purposes here, the disagreement cannot be fixed by scaling the spectrum, so we leave the flux calibration as is.

The G430L spectrum also contains features consistent with emission from the Ca II H&K lines at 3968.4673 and 3933.6614 Å. The Ca H line is blended with emission from the H ϵ line and so is given as an upper limit in Table 2.

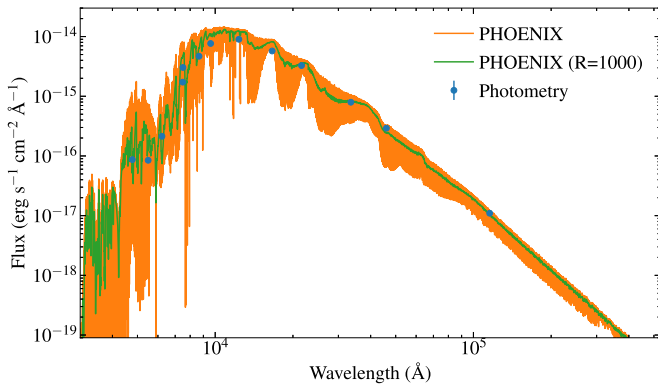


Figure 5. PHOENIX model compared with the photometry from Table 1 of Gonzales et al. (2019). The model in green has been convolved to a resolution of $R = 1000$ for easier comparison with the photometry.

3.5.2. PHOENIX

Wavelengths from 5700 Å to 100 μm are filled with a PHOENIX photospheric model spectrum from the Lyon BT-Settl [DIR] CIFIST2011_2015 grid¹⁸ (Baraffe et al. 2015; Allard 2016) as no ground-based spectroscopy of TRAPPIST-1 contemporaneous with our HST and XMM observations is available. Gonzales et al. (2019) used distance-calibrated spectra and photometry to obtain atmospheric parameters of TRAPPIST-1 of $T_{\text{eff}} = 2628 \pm 42$ K and $\log g = 5.21 \pm 0.06$ dex. We obtained the four closest models ($T_{\text{eff}}, \log g = 2600, 5.0; 2700, 5.0; 2600, 5.5; 2700, 5.5$) and linearly interpolated them onto the measured parameters using the SCIPY griddata routine. The model flux is then scaled by the square of the ratio of the measured radius and distance of TRAPPIST-1 ($1.16 \pm 0.03 R_{\text{Jup}}$ and 12.43 ± 0.02 pc, respectively, Gonzales et al. 2019). The BT-Settl models extend to 1000 μm, but to avoid the SED file size becoming too large, we truncated the model at 100 μm. The removed flux contributed less than 1% of the total integrated flux of the model. Figure 5 compares the PHOENIX model with the data from Gonzales et al. (2019) used to measure the atmospheric parameters.

4. Emission Line Flux Measurements

Producing the DEM and semiempirical models discussed below required identifying and measuring the fluxes of emission lines in the COS spectra. Table 2 provides fluxes for all of the lines in the lists compiled by Linsky (2017) and Peacock et al. (2019) that are covered by our spectra. Where lines from the list were detected, we fit them using a Gaussian profile combined with a linear fit to the surrounding continuum to account for incorrect background subtraction (Figure 6). The Gaussian profile does not always exactly recreate the apparent shape of the line profiles (for example, the C II lines), but, given the low S/N ratio, we cannot be confident that using more complex profiles would not lead to overfitting. We therefore do not claim to recover the shape of the line profile, but provide a reasonable measurement of the integrated flux. The flux is given as the integral of the Gaussian adjusted by the y value of the linear fit, along with the propagated statistical error of the fit. Where the fit was unsuccessful, we computed a 3σ upper limit by treating the line as a Gaussian with a width fixed to the average of nearby lines and an amplitude equal to three times

the error on the linear fit, adjusted by the linear fit as before. Al II 1670.787 Å is formally detected by this method, but as the noise in the spectrum surrounding the line is of comparable amplitude we give its flux as an upper limit. Due to the complexity of the line profile of the C III 1176 Å multiplet, the flux was measured as the integral of the flux between 1174.5–1176.5 Å with the uncertainty estimated as the sum of the rms of two adjacent 1 Å wide regions.

5. Differential Emission Measure

The differential emission measure (DEM) describes the amount of emitting plasma as a function of temperature for an optically thin plasma in coronal equilibrium (Warren et al. 1998; Loudén et al. 2017), allowing us to predict the fluxes of emission lines formed in the coronal atmosphere if we have a functional form to describe the DEM. The DEM combined with the emissivity of a particular emission line determines the emission flux in that line emitted by a star. By measuring the flux of FUV emission lines formed at $\sim 10^5$ K and the X-ray flux formed at $\gtrsim 10^6$ K, we can constrain the DEM by assuming it is well described by a smooth low-order polynomial across the temperature range relevant to the chromosphere, transition region, and corona. This operates under the same principle as the Linsky et al. (2014) empirical scaling relations but is tailored more specifically to the star by using more lines to resolve the structure of the star’s upper atmosphere at a finer temperature resolution. Combining the DEM with atomic data from CHIANTI v8.0 (Dere et al. 1997; Del Zanna et al. 2015), we estimate the EUV spectrum of TRAPPIST-1 between 10 to 912 Å, incorporating the errors in fitting the DEM forward to the predicted EUV spectrum. Our implementation is similar to the method employed by Loudén et al. (2017); details of the Mega-MUSCLES DEM prescription will be described in a future paper (Duvvuri et al. 2021). Figure 7 shows a comparison of the DEM model with the XMM data and the Linsky et al. (2014) EUV scaling relations.

6. Semiempirical Model

Given the unavoidable low S/N of ultraviolet observations of late M dwarfs in general and TRAPPIST-1 in particular, the Mega-MUSCLES data products will also include semiempirical models constructed from the observed spectra, which we recommend as inputs for model planetary atmosphere studies as they avoid potentially biasing the models with a large amount of noise in the observed spectra. The semiempirical model for TRAPPIST-1 is shown in Figure 8. Four models are used, including the APEC, DEM, and PHOENIX models already discussed. The fourth model replaces the HST spectra covering 1100–4200 Å. The model is constructed by first fitting a polynomial to the DEM model and the blue end ($\lambda < 4200$) of the STIS G430L spectrum to create a baseline flux. Note that this does not represent just the chromospheric/coronal continuum emission, but rather the combination of a forest of weak emission lines plus any continuum emission in that wavelength range. We then added the reconstructed Ly α line and the fits to the emission lines shown in Figure 6, as well as the PHOENIX spectrum at those wavelengths, to produce the final spectrum section.

Figure 8 compares our final semiempirical model SED with the quiet solar spectrum (Woods et al. 2009). The solar spectrum is scaled by the ratio of the blackbody luminosities of the photospheres of the two stars, such that the photospheric flux for

¹⁸ <https://phoenix.ens-lyon.fr/Grids/BT-Settl/>

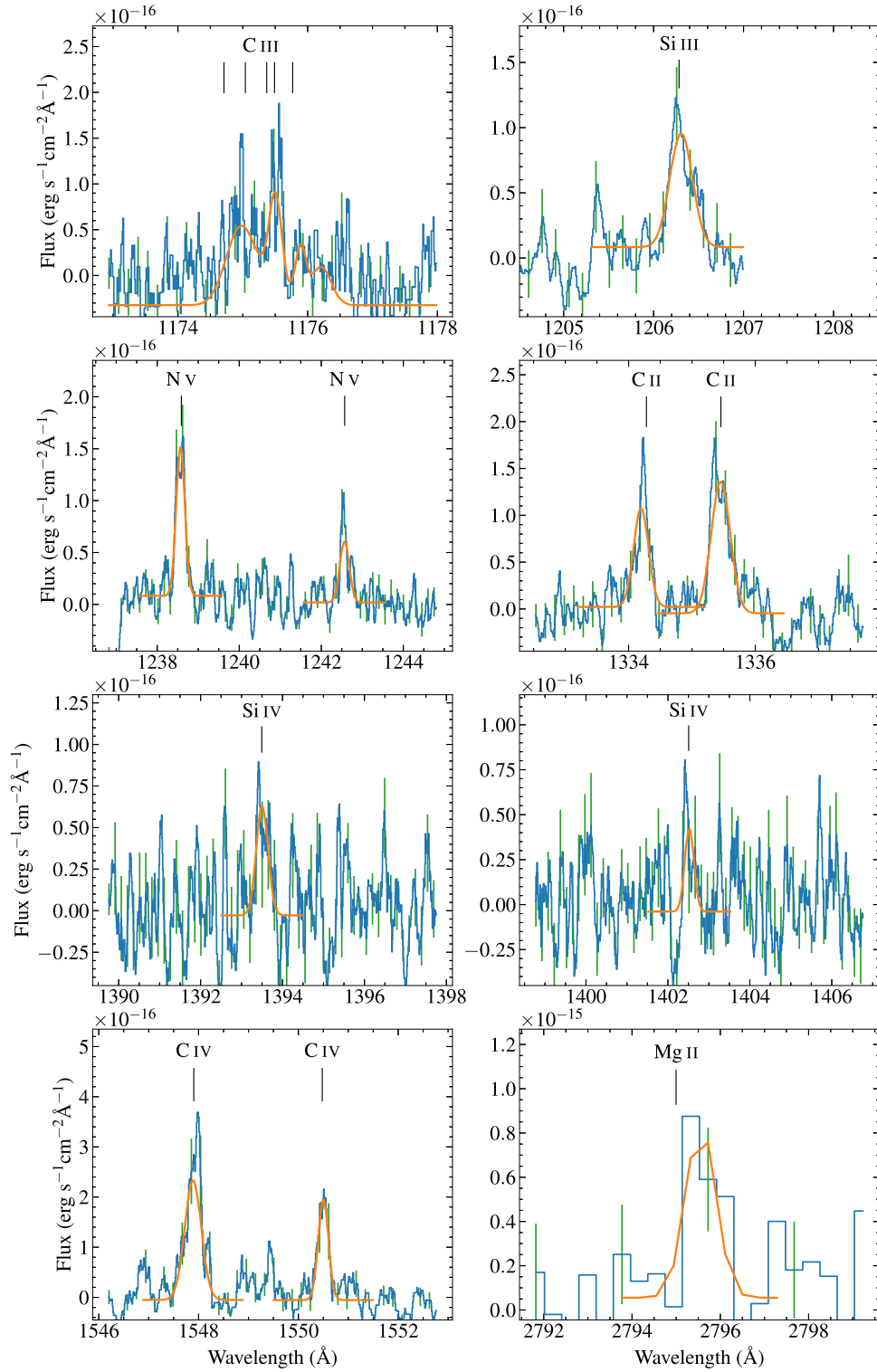


Figure 6. Detected emission lines in COS spectra of TRAPPIST-1 (blue) and the fit used to measure the integrated fluxes (orange). Example error bars on the flux values are shown in green. The data were smoothed with a 10 point boxcar for clarity, with the exception of the region around the Mg II lines. Line positions were taken from NIST and shifted by the -56.3 km s^{-1} radial velocity of TRAPPIST-1 (Reiners & Basri 2009b).

each spectrum is the same as the irradiation Earth receives. The comparison clearly demonstrates the relative difference in high-energy flux between TRAPPIST-1 and the Sun, implying that any planet receiving the same photospheric flux from TRAPPIST-1 as (for example) Earth receives from the Sun is experiencing high-energy flux levels that are several orders of magnitude higher (Figure 9), even when the star is not flaring.

7. Comparison with Model Spectra

Peacock et al. (2019) extended the PHOENIX stellar atmosphere code into the ultraviolet by adding contributions from the chromosphere and transition regions, providing model SEDs of the TRAPPIST-1 ultraviolet regions against which we can compare our observations. The three PHOENIX models

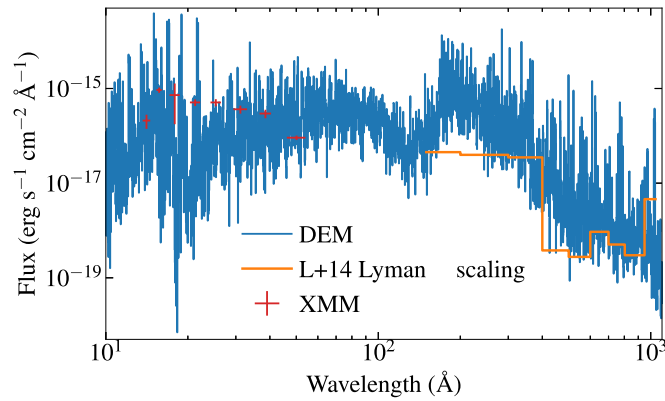


Figure 7. Comparison of the DEM model with the EUV/Ly α relationships of Linsky et al. (2014), and the XMM spectrum.

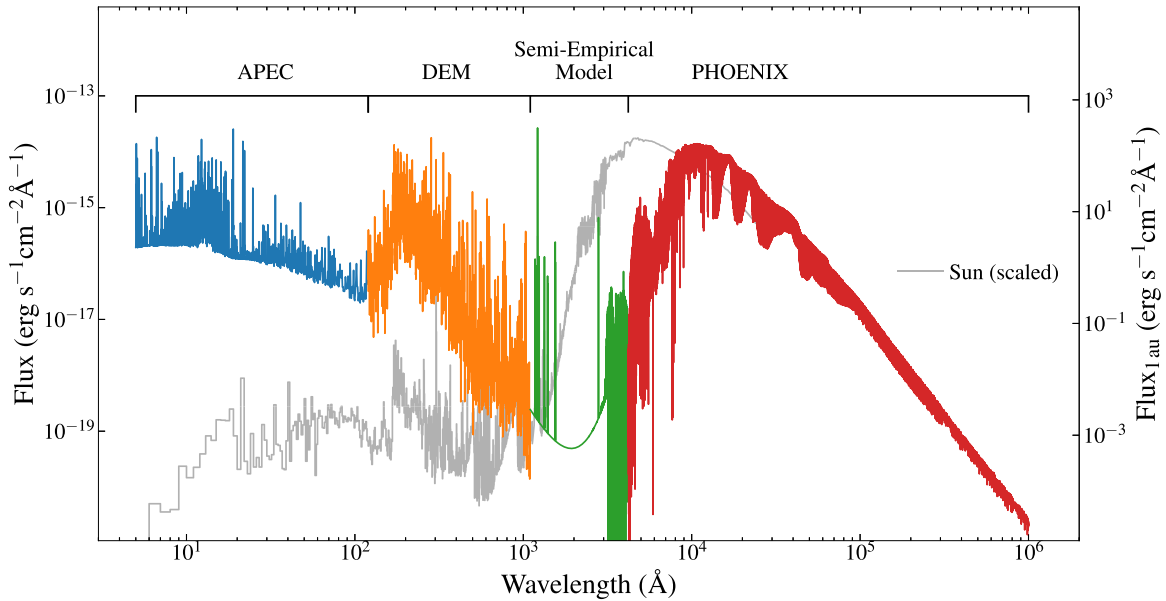


Figure 8. Semiempirical model spectrum of TRAPPIST-1. The four different model components used are labeled above the spectrum, and the PHOENIX model has been rebinned to 1 Å for clarity. The spectrum is compared with the quiet solar spectrum (Woods et al. 2009), scaled to have the same blackbody photospheric flux as TRAPPIST-1. The left axis shows the flux received from TRAPPIST-1 at Earth, whereas the right shows the flux received at the 1 au equivalent distance, i.e., the spectrum of TRAPPIST-1 at a distance receiving the same photospheric flux as at 1 au from the Sun.

were calibrated to the Bourrier et al. (2017b) Ly α flux (model 1A) and distance-adjusted GALEX photometry of stars with similar spectral types (models 2A and 2B).

Figure 10 compares the predicted line strengths given in Table 5 of Peacock et al. (2019) with the line fluxes measured from our COS data (Table 2) as a function of formation temperature and wavelength. For all three models, the agreement between the predicted and measured line fluxes is in general poor. The Ly α -scaled model 1A accurately predicts the measured fluxes of the C II and Ca II lines, but predicts multiple lines at values ~ 10 times higher than the upper limits placed on their fluxes here. We find no trend between the accuracy of the predicted fluxes and either formation temperature or wavelength.

8. Swift Photometry

As this paper was under review, Becker et al. (2020) presented a deep Swift uvm2 observation of TRAPPIST-1, finding a flux of $(8.4 \pm 1.4) \times 10^{-15}$ erg s $^{-1}$ cm $^{-2}$ around 1900 Å. As a large fraction of the uvm2 wave band is covered

by the gap in the COS G230L detector, an exact comparison between the SED and the Swift measurement is challenging. Integrating the semiempirical model over the uvm2 wave band, we find a flux of 1.5×10^{-15} erg s $^{-1}$ cm $^{-2}$, suggesting that the model may underpredict the flux in the NUV, although there is no way to further test this with the available data. We note that, although the uvm2 peaks at ≈ 1900 Å, the effective area is still $\approx 10\%$ of the maximum around the Mg II 2800 Å lines. As these are the only NUV features detected in our COS spectroscopy, it is possible that they are contributing the majority of the flux detected by Swift. Similar photometric measurements of emission lines may be an avenue to study the ultraviolet SEDs of low-mass stars too faint to observe spectroscopically.

9. High-level Science Products

The TRAPPIST-1 SED will be made available at or before the publication of this paper on the MUSCLES Treasury Survey Page at the Mikulski Archive for Space Telescopes (MAST):

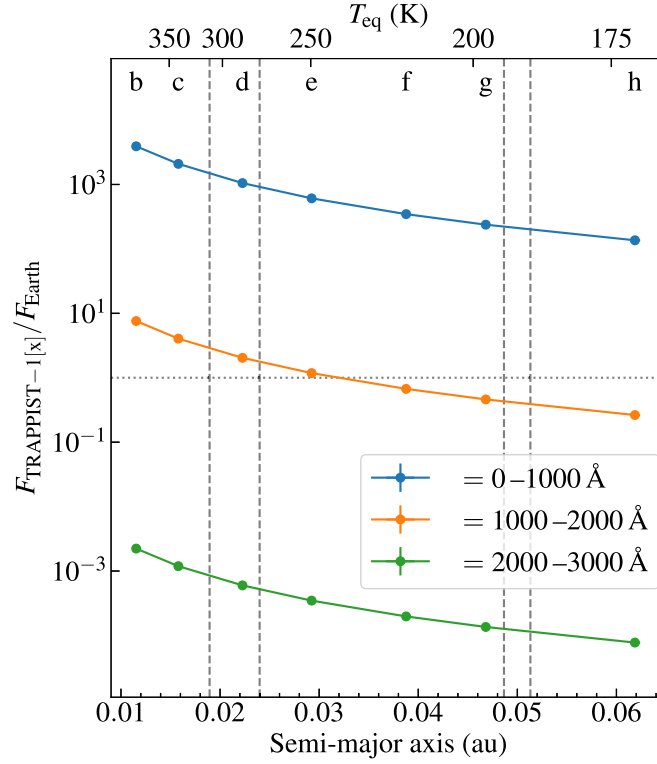


Figure 9. Fluxes at different wave bands experienced by the TRAPPIST-1 planets compared with Earth. [x] in the y-axis label refers to planets b, c, etc. The dashed lines show the empirical (including TRAPPIST-1 d) and conservative (excluding TRAPPIST-1 d) edges of the temperate zone for $1 M_{\oplus}$ planets as modeled by Kopparapu et al. (2014). The equilibrium temperature as a function of distance from the star is shown on the top axis.

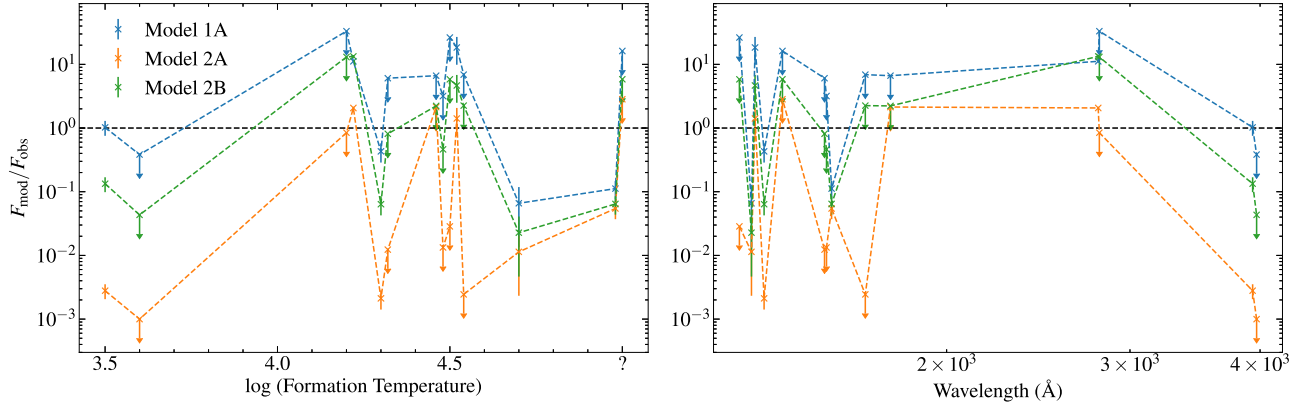


Figure 10. Comparison of the integrated flux of detected emission lines with predicted fluxes given in Table 5 of Peacock et al. (2019) as a function of formation temperature (left) and wavelength (right). We were unable to find a measured formation temperature for Fe II 1391 Å.

<https://archive.stsci.edu/prepds/muscles/>, 10.17909/T9DG6F. Available products will include the standard data products provided by the MUSCLES survey (i.e., the SED at native and 1 Å resolutions along with the component observations and models), along with the new semiempirical model SED at native and 1 Å resolutions. The remainder of the Mega-MUSCLES targets, listed at <http://cos.colorado.edu/~kevinf/muscles.html>, will be added to the database in the coming months.

10. Conclusion

We have constructed a panchromatic SED of the M8 star TRAPPIST-1, the first data product from the Mega-MUSCLES survey.

TRAPPIST-1 is the faintest target in the Mega-MUSCLES survey, and the SED presented here both represents the state-of-the-art for observation of the high-energy flux of low-mass stars and demonstrates the limits of our current observing facilities. Obtaining the ultraviolet spectroscopy pushed the capabilities of COS to their limits, with many expected emission lines remaining below the noise limit. The EUV spectrum cannot be observed with any currently operating facility. Improving on these observations, which is desirable given the continued importance of low-mass stars for exoplanet science, will require the launch of large-aperture space telescopes with ultraviolet capabilities and a dedicated EUV observatory (Youngblood et al. 2019).

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Facilities: HST (STIS and COS), XMM-Newton.

Software: astropy Astropy Collaboration et al. (2018), xspec (Arnaud 1996), stistools¹⁹, scipy (Virtanen et al. 2020), numpy (Harris et al. 2020), matplotlib (Hunter 2007).

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References

- Airapetian, V. S., Gloer, A., Khazanov, G. V., et al. 2017, *ApJL*, **836**, L3
 Allard, F. 2016, in SF2A-2016: Proc. Annual Meeting of the French Society of Astronomy and Astrophysics, ed. C. Reylé et al. (Paris: French Society of Astronomy and Astrophysics), 223
 Arnaud, K. A. 1996, in ASP Conf. Ser. 101, *Astronomical Data Analysis Software and Systems V*, ed. G. H. Jacoby & J. Barnes (San Francisco, CA: ASP), 17
 Astropy Collaboration, Price-Whelan, A. M., Sipőcz, B. M., et al. 2018, *AJ*, **156**, 123
 Baraffe, I., Homeier, D., Allard, F., & Chabrier, G. 2015, *A&A*, **577**, A42
 Barstow, J. K., & Irwin, P. G. J. 2016, *MNRAS*, **461**, L92
 Becker, J., Gallo, E., Hodges-Kluck, E., Adams, F. C., & Barnes, R. 2020, *AJ*, **159**, 275
 Bourrier, V., de Wit, J., Bolmont, E., et al. 2017a, *AJ*, **154**, 121
 Bourrier, V., Ehrenreich, D., Wheatley, P. J., et al. 2017b, *A&A*, **599**, L3
 Burgasser, A. J., & Mamajek, E. E. 2017, *ApJ*, **845**, 110
 Chambers, K. C., Magnier, E. A., Metcalfe, N., et al. 2016, arXiv:1612.05560
 Del Zanna, G., Dere, K. P., Young, P. R., Landi, E., & Mason, H. E. 2015, *A&A*, **582**, A56
 Dere, K. P., Landi, E., Mason, H. E., Monsignori Fossi, B. C., & Young, P. R. 1997, *A&AS*, **125**, 149
 de Wit, J., Wakeford, H. R., Gillon, M., et al. 2016, *Natur*, **537**, 69
 Duvvuri, G. M., Pineda, J. S., & Berta-Thompson, Z. K. 2021, arXiv:2102.08493
 Foreman-Mackey, D., Hogg, D. W., Lang, D., & Goodman, J. 2013, *PASP*, **125**, 306
 Foster, A. R., Ji, L., Smith, R. K., & Brickhouse, N. S. 2012, *ApJ*, **756**, 128
 France, K., Duvvuri, G., Egan, H., et al. 2020, *AJ*, **160**, 237
 France, K., Loyd, R. O. P., Youngblood, A., et al. 2016, *ApJ*, **820**, 89
 Gillon, M., Jehin, E., Lederer, S. M., et al. 2016, *Natur*, **533**, 221
 Gillon, M., Triaud, A. H. M. J., Demory, B.-O., et al. 2017, *Natur*, **542**, 456
 Gonzales, E. C., Faherty, J. K., Gagné, J., et al. 2019, *ApJ*, **886**, 131
 Green, J. C., Froning, C. S., Osterman, S., et al. 2012, *ApJ*, **744**, 60
 Harris, C. R., Millman, K. J., van der Walt, S. J., et al. 2020, *Natur*, **585**, 357
 Hunter, J. D. 2007, *CSE*, **9**, 90
 Koppurapu, R. K., Ramirez, R. M., SchottelKotte, J., et al. 2014, *ApJL*, **787**, L29
 Lançon, A., Gonneau, A., Verro, K., et al. 2020, arXiv:2012.09129
 Linsky, J. L. 2017, *ARA&A*, **55**, 159
 Linsky, J. L., Fontenla, J., & France, K. 2014, *ApJ*, **780**, 61
 Loudon, T., Wheatley, P. J., & Briggs, K. 2017, *MNRAS*, **464**, 2396
 Loyd, R. O. P., France, K., Youngblood, A., et al. 2016, *ApJ*, **824**, 102
 Loyd, R. O. P., France, K., Youngblood, A., et al. 2018, *ApJ*, **867**, 71
 Luger, R., Sestovic, M., Kruse, E., et al. 2017, *NatAs*, **1**, 0129
 Miguel, Y., Kaltenegger, L., Linsky, J. L., & Rugheimer, S. 2015, *MNRAS*, **446**, 345
 Morley, C. V., Kreidberg, L., Rustamkulov, Z., Robinson, T., & Fortney, J. J. 2017, *ApJ*, **850**, 121
 O'Malley-James, J. T., & Kaltenegger, L. 2017, *MNRAS*, **469**, L26
 Paudel, R. R., Gizis, J. E., Mullan, D. J., et al. 2018, *ApJ*, **858**, 55
 Peacock, S., Barman, T., Shkolnik, E. L., Hauschildt, P. H., & Baron, E. 2019, *ApJ*, **871**, 235
 Redfield, S., & Linsky, J. L. 2008, *ApJ*, **673**, 283
 Reiniers, A., & Basri, G. 2009a, *ApJ*, **705**, 1416
 Reiniers, A., & Basri, G. 2009b, *A&A*, **496**, 787
 Ribas, I., Bolmont, E., Selsis, F., et al. 2016, *A&A*, **596**, A111
 Roettenbacher, R. M., & Kane, S. R. 2017, *ApJ*, **851**, 77
 Rugheimer, S., Segura, A., Kaltenegger, L., & Sasselo, D. 2015, *ApJ*, **806**, 137
 Smith, R. K., Brickhouse, N. S., Liedahl, D. A., & Raymond, J. C. 2001, *ApJL*, **556**, L91
 Veras, D., Armstrong, D. J., Blake, J. A., et al. 2018, *AsBio*, **18**, 1106
 Vida, K., Kővári, Z., Pál, A., Oláh, K., & Kriskovics, L. 2017, *ApJ*, **841**, 124
 Virtanen, P., Gommers, R., Oliphant, T. E., et al. 2020, *Nature Methods*, **17**, 261
 Wang, S., Wu, D.-H., Barclay, T., & Laughlin, G. P. 2017, arXiv:1704.04290
 Warren, H. P., Mariska, J. T., & Lean, J. 1998, *JGR*, **103**, 12077
 Wheatley, P. J., Loudon, T., Bourrier, V., Ehrenreich, D., & Gillon, M. 2017, *MNRAS*, **465**, L74
 Wood, B. E. 2004, *LRSP*, **1**, 2
 Woodgate, B. E., Kimble, R. A., Bowers, C. W., et al. 1998, *PASP*, **110**, 1183
 Woods, T. N., Chamberlin, P. C., Harder, J. W., et al. 2009, *GeoRL*, **36**, L01101
 Youngblood, A., Drake, J., Mason, J., et al. 2019, *BAAS*, **51**, 300
 Youngblood, A., France, K., Loyd, R. O. P., et al. 2016, *ApJ*, **824**, 101

¹⁹ <https://stistools.readthedocs.io/en/latest/>