

PAPER • OPEN ACCESS

The Role of CHPD and AIMI processing on enhancing J_c and transverse connectivity of *in-situ* MgB₂ strand

To cite this article: F Wan *et al* 2020 *IOP Conf. Ser.: Mater. Sci. Eng.* **756** 012018

View the [article online](#) for updates and enhancements.

You may also like

- [Understanding routes for high connectivity in *ex situ* MgB₂ by self-sintering](#)
Shunsuke Mizutani, Akiyasu Yamamoto, Jun-ichi Shimoyama et al.
- [Mechanism for high critical current density in *in situ* MgB₂ wire with large area-reduction ratio](#)
Motomune Kodama, Yota Ichiki, Kazuhide Tanaka et al.
- [Electromagnetic densification of MgB₂/Cu wires](#)
M Woniak and B A Glowacki



ECS
The
Electrochemical
Society
Advancing solid state &
electrochemical science & technology

DISCOVER
how sustainability
intersects with
electrochemistry & solid
state science research

The Role of CHPD and AIMI processing on enhancing J_C and transverse connectivity of *in-situ* MgB₂ strand

F Wan¹, M D Sumption¹, M A Rindfleisch², and E W Collings¹

¹Center for Superconductor and Magnetic Materials, Department of Materials Science and Engineering, The Ohio State University, Columbus, OH 43210, USA.

Corresponding author e-mail: wan.108@osu.edu

²Hyper Tech Research Inc, Columbus, OH 43228, USA

Abstract. Research into *in-situ* MgB₂ strand has been focused on improvements in J_C through reduction of porosity. Both of cold-high-pressure-densification (CHPD) and advanced-internal-magnesium-infiltration (AIMI) techniques can effectively remove the voids in *in-situ* MgB₂ strands. This study shows the nature of the reduced porosity for *in-situ* MgB₂ strands lies on increases in transverse grain connectivity as well as longitudinal connectivity. The CHPD method bi-axially applying 1.0 GPa and 1.5 GPa yielded 4.2 K $J_{C||}$ s of 9.6×10^4 A/cm² and 8.5×10^4 A/cm² at 5 T, respectively, with compared with 6.0×10^4 A/cm² for typical powder-in-tube (PIT) *in-situ* strand. Moreover, AIMI-processed monofilamentary MgB₂ strand obtained even higher J_C s and transverse grain connectivity than the CHPD strands.

1. Introduction

MgB₂ superconducting strands are promising to the practical magnetic application due to its high transition temperature T_C (39 K) [1], high coherence length [2, 3], and low anisotropy of upper critical fields (B_{c2}) [3–6]. The powder-in-tube (PIT) *in-situ* MgB₂ strands were fabricated by filling a mixture of Mg and B powder into a non-reactive metallic tube and then being cold-worked into wires or tapes. The PIT strands have large amount of voids elongated along longitudinal strand axis which were left behind by molten Mg powders after heat treatment [7]. The present of voids tends to limit the number of the continuous current path in the *in-situ* MgB₂ strand and therefore suppress the current-carrying-capacity of the strand.

The CHPD technique can effectively increase the transport properties of PIT *in-situ* MgB₂ strands by eliminating the pores [8–11]. The pre-reacted powder-in-tube composite was bi-axially cold-densified at room temperature to increase the mass density of the Mg + B mixture. In this case, the cold-densified MgB₂ strand can obtain higher grain connectivity after heat treatment. Additionally, the Mg reactive – liquid – infiltration (RLI) process, initiated by Giunchi *et al* [12], also has the ability to eliminate the pores and produce a dense MgB₂ layers in *in-situ* MgB₂ strands. For RLI process, a Mg rod is inserted axially into a boron-filled metallic tube. After wire drawing the heat treatment (H. T.), MgB₂ layer is formed through the reactive diffusion of Mg into B layer. Since Mg is totally separated



with precursor B layer before H. T., the RLI process can totally eliminate the “Mg-site porosity” from MgB_2 layer but induce the formation of a big hole at the central region of the strand [13]. Furthermore, since the molar volume of MgB_2 ($17.46 \text{ cm}^3/\text{mol}$) is twice as that of two B atoms ($9.18 \text{ cm}^3/\text{mol}$), the volume expansion associated with the reactive transformation from 2B to MgB_2 enables even better connections between MgB_2 grains during heat treatment [13]. Our group named our RLI-processed strands as advanced – internal – magnesium – infiltration (AIMI) strands due to their optimized strand architecture and high J_{CS} [13].

The previous researches mostly focused on investigating the effect of CHPD and AIMI techniques on the transport properties along longitudinal strand axis of *in-situ* MgB_2 strands, such as transport J_{CS} and longitudinal grain connectivity. However, Shi and Susner pointed out that high anisotropic grain connectivity exists in PIT *in-situ* strands, which was resulted from the elongated voids between elongated MgB_2 stringers [7]. Therefore longitudinal transport properties are different with the transport properties along transverse strand axis for MgB_2 superconducting wires. In continuing along these lines, the transport properties along transverse strand axis were investigated for the 2.0 mol% C-doped *in-situ* MgB_2 strands in this study. The anisotropic connectivity of the PIT strands results in the differences between perpendicular magnetic J_C ($J_{C\perp}$) and parallel J_C ($J_{C\parallel}$). The influence of aspect ratio ($S = \text{length/diameter}$) on transverse and longitudinal J_C were investigated for the PIT strand (P00). The CHPD-processed strands were P10 (1.0 GPa cold-pressing) and P15 (1.5 GPa cold-pressing). According to the previous results of our group, the CHPD technique increased the transport J_C of the monofilamentary PIT *in-situ* MgB_2 strand from $3.0 \times 10^4 \text{ A/cm}^2$ to $3.6 \times 10^4 \text{ A/cm}^2$ at 4.2 K and 10 T due to decreased porosity [11] and AIMI-processed MgB_2 strands attained the 4.2 K, 10 T transport layer J_{CS} of $1.0 \sim 1.5 \times 10^5 \text{ A/cm}^2$ [13-16]. In this study, we compared the $J_{C\parallel}$ s and transverse connectivity of the CHPD- and AIMI-processed strands with those of the P00 strand at 4.2 K and 20 K. The relationship between porosity and transverse flux pinning force density $F_{p\parallel}$, which is $J_{C\parallel} \times B$, for the *in-situ* MgB_2 strands was also discussed.

2. Experimental

2.1. Sample preparation

A series of pre-reacted powder-in-tube (PIT) *in-situ* strands, typically 0.834 mm diameter, with a Nb barrier and a Cu outer sheath were fabricated by Hyper Tech Research, Inc. (HTR). Two PIT strands (designated P10 and P15) were bi-axially densified with 1.0 GPa and 1.5 GPa at room temperature, respectively. The other strand (designed A00) manufactured through AIMI technique were also provided by HTR. The AIMI-processed strand, with 0.55 mm diameter, has a Nb barrier and a Monel outer sheath. The powders used for the present strands were 2 mol% C-doped amorphous B (10 – 100 nm) from Specialty Materials Inc. (SMI). The cold-densified PIT strands were heat-treated at 675 °C for 1 h and the AIMI strand was heat – treated at 625 °C for 16 h. The specification and heat treatment (H. T.) conditions of the strands are presented in table 1.

2.2. Transport and Magnetic Measurements

The transport I_C (I_{CT}) test was conducted in perpendicular magnetic field up to 13 T in a pool of liquid Helium at 4.2 K on the MgB_2 strands with a total length of 50 mm and a gauge length of 5 mm. The electric criterion used for determining I_{CT} s is $1.0 \mu\text{V/cm}$. The magnetizations versus perpendicular and parallel magnetic fields ($M - H$) loops were measured by a Quantum Design Model 6000 Physical Property Measuring System (PPMS) for all strands with a sample length of 3 – 5 mm.

Table 1. Strand specification and H. T. conditions

Strand Name	Strand Type	CHPD (GPa)	H. T. (h/°C)
P00	PIT	0.0	1/675
P10	PIT	1.0	1/675
P15	PIT	1.5	1/675
A00	AIMI	0.0	16/625

3. Results

The J_{CT} s of the PIT strands were the transport critical current normalized by MgB_2 core area. As shown in Figure 1(a) the MgB_2 core of the typical PIT *in-situ* strand is a solid cylinder. Figure1(b)-(d) shows the shape of MgB_2 cores for the CHPD- and AIMI-processed strands are cuboid and hollow cylinder, respectively. The J_{CT} of the AIMI-processed strands, which is also named as transport layer J_C , were calculated by dividing the I_{CT} by the area of annulus MgB_2 layer. Values of magnetic J_C ($J_{CM\perp}$ and $J_{CM\parallel}$) for the MgB_2 strands were extracted from the full $M - H$ loops heights ΔM , using the standard Bean model equations [7, 17]:

For the PIT *in-situ* wire P00:

$$\text{Perpendicular Magnetic } J_C: J_{CM\perp} = \frac{3\pi\Delta M}{8R_0} \quad (1)$$

$$\text{Parallel Magnetic } J_C: J_{CM\parallel} = \frac{3\Delta M}{2R_0} \quad (2)$$

Here R_0 is the radius of the cylinder MgB_2 core in the PIT wire.

For the densified wires P10 and P15:

$$\text{Parallel Magnetic } J_C: J_{CM\parallel} = \frac{2\Delta M}{b(1-\frac{b}{3a})} \quad (3)$$

Here a , b are both lengths of the transverse cross – sectional area of cuboid MgB_2 core, $a > b$.

For the AIMI wire A00:

$$\text{Parallel Magnetic } J_C: J_{C\parallel} = \frac{3\Delta M}{2} \frac{R_0^2 - R_i^2}{R_0^3 - R_i^3} \quad (4)$$

Here R_i is the inner diameter of the annulus MgB_2 layer and R_0 is the outer diameter of the annulus MgB_2 layer.

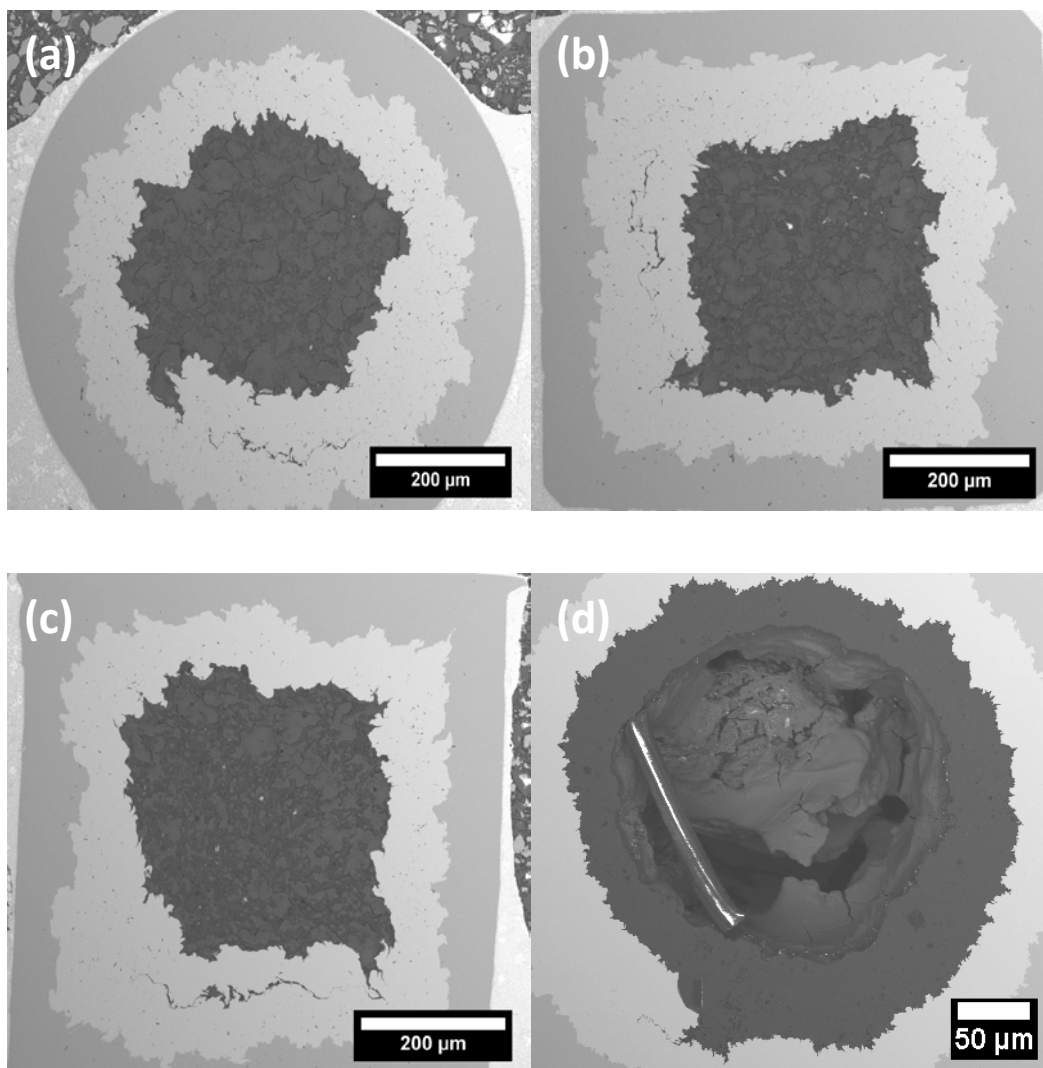


Figure 1. (a) Back scattered SEM images of (a) strand P00, (b) strand P10 (1.0 GPa cold pressure), (c) strand P15 (1.5 GPa cold pressure), and (d) AIMI strand A00.

3.1. Transport and magnetic critical current densities, J_{CT} and J_{CM}

Figure 2(a) shows the J_{CT} and $J_{CM\perp}$ versus B at 4.2 and 20 K for the strand P00. It can be seen that $J_{CM\perp}$ s agree with J_{CT} at low fields, whereas the bifurcation of J_{CT} and $J_{CM\perp}$ happened at high fields. Moreover, $J_{CM\perp}$ s were greatly affected by the aspect ratio S , especially at high magnetic fields. The relationships among J_{CT} , $J_{CM\perp}$, and aspect ratio were fully discussed in ref [7] and [18]. Therefore, the values of $J_{CM\perp}$ are not only determined by the intrinsic properties but also the extrinsic properties of the *in-situ* MgB₂ strands. As shown in Figure 2(b) $J_{CM\parallel}$ s were independent on the aspect ratio S . Based on this result, we can know the values of $J_{CM\parallel}$ s as well as J_{CT} s are merely determined by the intrinsic properties of the *in-situ* MgB₂ strands. The MgB₂ macrostructure of the reacted PIT *in-situ* strand is characterized by elongated polycrystalline MgB₂ fibers and elongated pores partially separating the fibers [7, 11, and 18]. The J_{CT} s and $J_{CM\parallel}$ s can represent the longitudinal and transverse connectivity of the MgB₂ fibers, respectively. In summary, we can investigate the effect of CHPD and AIMI technique on the current-carrying-capacity of the *in-situ* MgB₂ strands through $J_{CM\parallel}$ s as well as J_{CT} .

Figure 3 shows the J_{CT} and $J_{CM\parallel}$ versus B at 4.2 and 20 K for all strands. The CHPD technique significantly enhanced $J_{CM\parallel}$ s of PIT *in-situ* strands at 4.2 K. Strand P10 and P15 attained 9.5×10^4 A/cm² and 8.5×10^4 A/cm² at 4.2 K and 5 T, with respect to 6.0×10^4 A/cm² for strand P00. 4.2 K J_{CT} s were slightly enhanced by the CHPD technique. Therefore, the increases in J_{CT} and $J_{CM\parallel}$ indicated that the both longitudinal and transverse connections between MgB₂ fibers were enhanced. With the formation of high density MgB₂ layer, strand A00 attained 4.2 K, 10 T J_{CT} of 9.4×10^4 A/cm², which is 180% higher than those of CHPD-processed strands. On the other hand, the AIMI strand obtained 5 T $J_{CM\parallel}$ s of 3.1×10^5 A/cm² at 4.2 K and 1.7×10^4 A/cm² at 20 K, which are about 240% higher than those of CHPD strands.

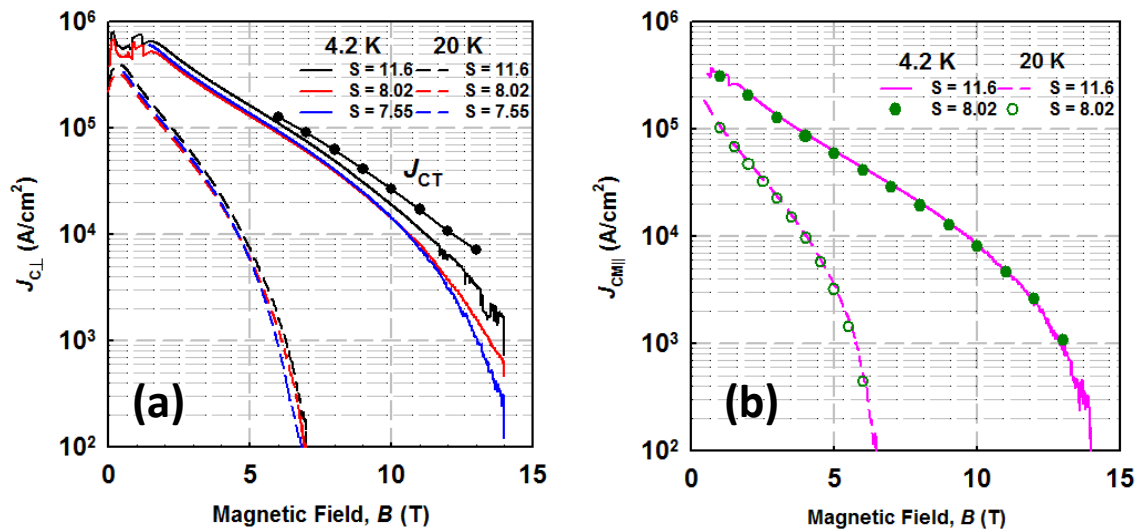


Figure 2. (a) J_{CT} and $J_{CM\perp}$ versus B at 4.2 and 20 K (b) $J_{CM\parallel}$ versus B at 4.2 and 20 K for strand P00

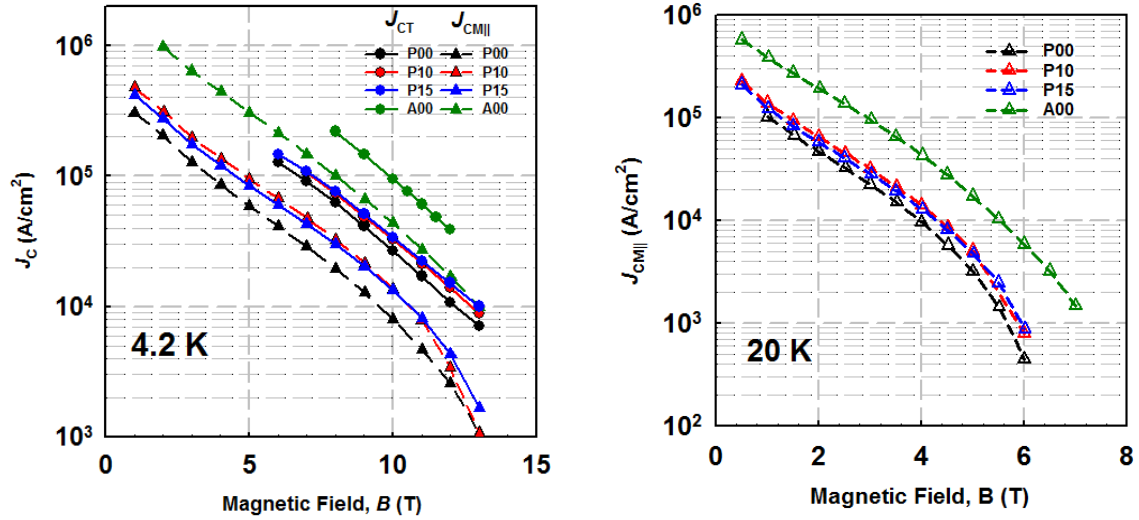


Figure 3. J_{CT} and J_{CMII} versus B at 4.2 and 20 K for all strands

3.2. Transverse Flux Pinning, Porosity and Transverse Grain Connectivity

Figure 4 shows the $F_{p||}/F_{p,max||}$ versus B/B_{irr} for all strands, where $F_{p||} = J_{CMII} \times B$. It can be seen that the peak pinning occurred at $b = B/B_{irr||}$ close to 0.2 at 4.2 and 20 K, which is in agreement with the Dew-Hughes/Kramer model [19, 20]. In other words, the dominant pinning centers for the all strands are also grain boundaries for the direction along transverse strand axis. According to the previous work [11], the densified wires have decreased porosities, the values of porosity are $\sim 50\%$ ($p = 0$) and $\sim 30\%$ ($p = 1.0$ or 1.5 GPa). Since the Mg-site porosity was totally eliminated for AIMI-processed strands, the porosity of the AIMI strand is close to 0. Therefore, both CHPD and AIMI processes enable the resulting MgB_2 phases to be denser and more connected. For the discussion, it can be concluded that the enhanced J_{CT} and J_{CMII} in densified wires and AIMI wire is correlated with lower porosity (higher grain connectivity).

The connectivity K defined by Rowell [21] can represent the grain connectivity of MgB_2 strands. The connectivity K is calculated by the equation:

$$K = \frac{\Delta\rho_{SC}}{\Delta\rho} \quad (5)$$

Here $\Delta\rho$ is the difference between the sample's resistivity at 300 K and the sample's resistivity at 40 K and $\Delta\rho_{SC}$ is the resistivity difference for an ideal single crystal. However, transverse connectivity $K_{||}$ is difficult to be determined for the MgB_2 wires. It has been reported that the maximum flux pinning force densities of fully-connected MgB_2 superconductor, where grain connectivity $K = 100\%$, are estimated to be 90 GN/m^3 at 4.2 K and 22 GN/m^3 at 20 K [22]. Therefore, we can roughly estimate the transverse grain connectivities $K_{||s}$ of the *in-situ* MgB_2 strands by normalizing the $F_{p,max||}$ with the $F_{p,max}$ of fully-connected MgB_2 . Table 2 shows the transverse maximum flux pinning force densities $F_{p,max||}$ at 4.2 K and 20 K and estimated transverse grain connectivities for all the strands. The connectivity of 5% is achieved by the 2 mol% C-doped PIT strand, P00. The cold-densification increases the

connectivity of PIT strand by 20 %. The AIMI strand A00 obtained the highest transverse grain connectivities of 20 %.

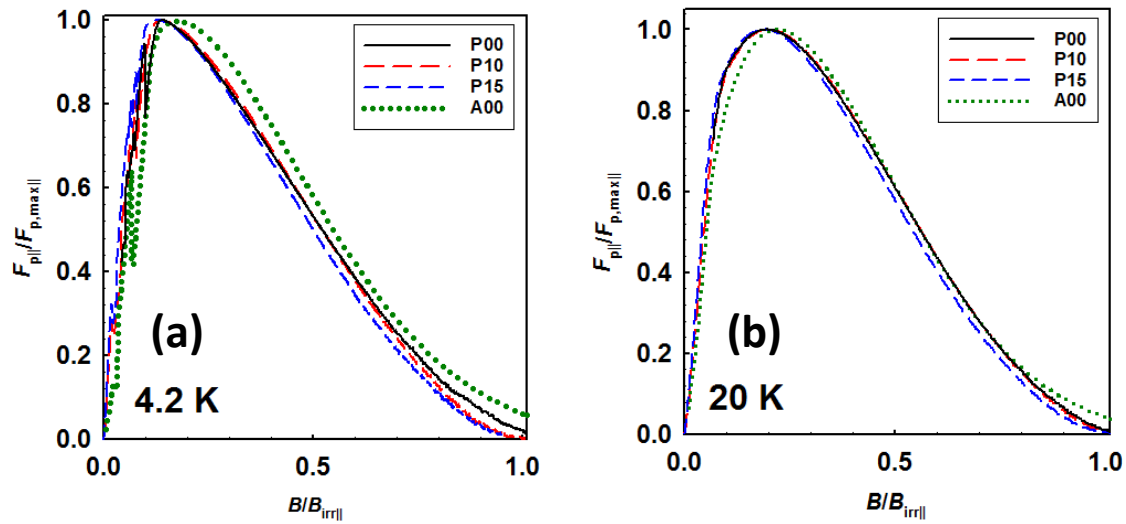


Figure 4. (a) $F_{p||}/F_{p,max||}$ versus B/B_{irr} for all strands at 4.2 K, (b) $F_{p||}/F_{p,max||}$ versus B/B_{irr} for all strands at 20K

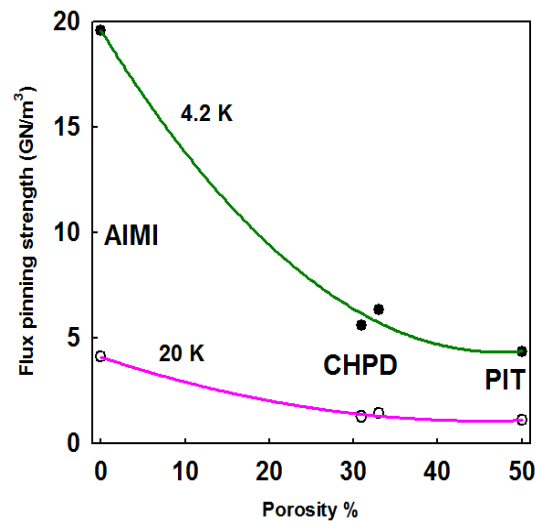


Figure 5. Relationship between porosity and transverse flux pinning force densities for *in-situ* MgB₂ strands

Table 2. Transverse flux pinning force densities and estimated connectivity of the strands.

Strand Name	$F_{p,max , 4.2\text{ K}} \text{ (GN/m}^3\text{)}$	$F_{p,max , 20\text{ K}} \text{ (GN/m}^3\text{)}$	Transverse Connectivity, $K_{ } \%$
P00	4.3	1.1	5.0
P10	6.3	1.4	6.0
P15	5.6	1.3	6.0
A00	19.6	4.1	20

4. Conclusion

$J_{CM||}$ is merely determined by the intrinsic properties of *in-situ* MgB_2 strands, so the current-carrying capacity of *in-situ* MgB_2 strands can be represented by $J_{CM||}$ as well as J_{CT} . The CHPD of 1.0 GPa and 1.5 GPa enhanced the 4.2 K, 5 T $J_{CM||}$ from $6.0 \times 10^4 \text{ A/cm}^2$ to $9.6 \times 10^4 \text{ A/cm}^2$ and $8.5 \times 10^4 \text{ A/cm}^2$, respectively. AIMI strand attained the highest J_{CT} and $J_{CM||}$ at 4.2 and 20 K due to the formation of a high dense MgB_2 layer. By eliminating the voids in *in-situ* MgB_2 strands through CHPD and AIMI technique, better connections between MgB_2 grains along transverse strand axis can be obtained in *in-situ* MgB_2 strands

5. References

- [1] Nagamatsu J, Nakagawa N, Muranaka T, Zenitani Y, and Akimitsu J 2001 *Nature* **410** 63-64
- [2] Finnemore D K, Ostenson J E, Bud'ko S L, Lapertot G and Canfield PC 2001 *Rhys. Rev. Lett* **86** 2420-2422
- [3] Vinod K, Abhilash Kumar R G and Syamaprasad U 2007 *Supercond. Sci. Technol* **20** R1-R3
- [4] Eltsev Y, Lee S, Nakao K, Chikumoto N, Tajima S, Koshizuka N and Murakami M 2002 *Phys. Rev. B* **65** 140501
- [5] Eisterer M, Zehetmayer M and Weber H W 2003 *Phys. Rev. Lett* **90** 247002
- [6] De Liam O F, Ribeiro R A, Avila M A, Cardoso C A and Coelho A A 2001 *Phys. Rev. Lett* **86** 5974
- [7] Shi Z X, Susner M A, Majoros M, Sumption M D, Peng X, Rindfleisch M, Tomsic M J and Collings E W 2010 *Supercond. Sci. Technol* **23** 045018
- [8] Flukiger R, Hossain M S A and Senatore C 2009 *Supercond. Sci. Technol* **22** 085002
- [9] Hossain M S A, Senatore C, Flukiger R, Rindfleisch M A, Tomsic M J, Kim J H and Dou S X 2009 *Supercond. Sci. Technol* **22** 095004
- [10] Hossain M S A et al 2014 *Supercond. Sci. Technol* **27** 095016
- [11] Wan F, Sumption M D, Rindfleisch M A and Collings E W 2017 *IOP Conf. Series: Materials Science and Engineering* **279** 012024
- [12] Giunchi G, Ripamonti G, Perini E, Cavallin T and Bassani E 2007 *IEEE TRANSACTIONS ON APPLIED SUPERCONDUCTIVITY* **17** 2761
- [13] Li G Z, Sumption M D, Susner M A, Yang Y, Reddy K M, Rindfleisch M A, Tomsic M J, Thong C J and Collings E W 2012 *Supercond. Sci. Technol* **25** 115023
- [14] Ye S J, Matsumoto A, Zhang Y C and Kumakura H 2014 *Supercond. Sci. Technol* **27** 085012
- [15] Li G Z, Sumption M D, Zwyer J B, Susner M A, Rindfleisch M A, Thong C J, Tomsic M J and Collings E W 2013 *Supercond. Sci. Technol* **26** 095007
- [16] Kumakura H, Hur J, Togano K, Matsumoto A, Wada H and Kimura K 2011 *IEEE Trans. Appl. Supercond.* **21** 2643
- [17] Sumption M D, Peng X, Lee X, Wu X and Collings E W 2004 *Cryogenics* **44** 711-725
- [18] Susner M A, Daniels T W, Sumption M D, Rindfleisch M A, Thong C J and Collings E W 2012 *Supercond. Sci. Technol* **25** 065002
- [19] Dew-Hughes D 1974 *Philos. Mag.* **30** 293-305
- [20] Kramer E J 1973 *J. Appl. Phys.* **44** 1360-70
- [21] Rowell J 2003 *Supercond. Sci. Technol* **16** R17-27
- [22] Matsushita T, Kiuchi M, Yamamoto A, Shimoyama J and Kishio K 2008 *Physica C* **468** 1833