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# The simplest oscillating solutions of nonlocal nonlinear models 

A O Smirnov and E E Aman<br>Advanced Mathematics and Mechanics, FTTI, State University of Aerospace Instrumentation, Saint Petersburg, 67A, Bolshaya Morskaya str., Saint Petersburg, 190000, Russia<br>E-mail: alosm@mail.ru


#### Abstract

In their recent works, Ablowitz and Musslimani proposed a new type of integrable nonlinear equations - nonlocal analogues of the nonlinear Schrödinger equation, the modified Korteweg-de Vries equation, and other nonlinear differential equations. In subsequent works, numerous researchers constructed the simplest soliton and rational solutions of these equations. In this paper, we construct the simplest oscillating solutions of some of the integrable nonlocal nonlinear differential equations associated to the nonlinear Schrödinger equation.


## 1. Introduction

Research in this area began with the work of Ablowitz and Musslimani [2], where the authors modified the Lax pair for the nonlinear Schrödinger equation. The authors also found [2] a one-soliton solution of the nonlocal nonlinear Schrödinger equation by the inverse scattering problem. Then they continued their research on this topic in works [1, 3-5, 10], the intermediate results of which are published in [5], where 16 types of integrable nonlinear nonlocal differential equations are considered and analyzed.
Research on this topic has attracted other researchers who have published at least 15 papers (see, for example, [6-8, 11-13, 16-20, 22, 23, 25]) on the subject in 2018 alone. In these works, as a rule, the authors construct soliton or rational solutions by methods of Darboux, Hirota transformation or selfsimilar substitution. Also, a large number of papers on this subject have been published in previous years, since each of the authors applied his method to several nonlocal equations proposed in [5] (note that the paper [5] was previously published in 2016 as a preprint arXiv:1610.02594). There is no doubt that research on this subject will continue (see, for example, [9, 21, 24]).
We remind that the equations of the AKNS hierarchy have the form

$$
\begin{equation*}
p_{t k}=-i^{k} H_{k+1}(p, q), q_{t k}=-i^{k} G_{k+1}(p, q) \text { or } p_{t k}+i^{k} H_{k+1}(p, q)=0, q_{t k}+(-i)^{k} H_{k+1}(p, q)=0, \tag{1}
\end{equation*}
$$

where the functions $H_{k}$ and $G_{k}$ satisfy the following equations [14, 15]

$$
\begin{gather*}
H_{1}(p, q)=-p_{x}, G_{1}(p, q)=-q_{x},\left(F_{k}(p, q)\right)_{x}=-p G_{k}(p, q)-q H_{k}(p, q), H_{k+1}(p, q)=2 p F_{k}(p, q)+  \tag{2}\\
\left(H_{k}(p, q)\right)_{x}, G_{k+1}(p, q)=-2 q F_{k}(p, q)-\left(G_{k}(p, q)\right)_{x} .
\end{gather*}
$$

In particular,

$$
\begin{gather*}
F_{1}(p, q)=p q, H_{2}(p, q)=2 p^{2} q-p_{x x}, \\
G_{2}(p, q)=-2 q^{2} p+q_{x x}, F_{2}(p, q)=p_{x} q-p q_{x},  \tag{3}\\
H_{3}(p, q)=6 p q p_{x}-p_{x x}, G_{3}(p, q)=6 p q q_{x}-q_{x x x}, \\
F_{3}(p, q)=p q_{x x}+q p_{x x}-p_{x} q_{x}-3 p^{2} q^{2},
\end{gather*}
$$

$$
\begin{gathered}
H_{4}(p, q)=-6 p^{3} q^{2}+6 q p_{x}^{2}+4 p p_{x} q_{x}+8 p q p_{x x}+2 p^{2} q_{x x}-p_{x x x x} \\
G_{4}(p, q)=6 p^{2} q^{3}-6 p q_{x}^{2}-4 q p_{x} q_{x}-8 p q q_{x x}-2 q^{2} p_{x x}-q_{x x x x} \\
F_{4}(p, q)=-6 p q^{2} p_{x}+6 p^{2} q q_{x}-q_{x} p_{x x}+p_{x} q_{x x}+q p_{x x x}-p q_{x x x}
\end{gathered}
$$

It is easy to show that the functions $F_{k}(p, q), H_{k}(p, q)$ and $G_{k}(p, q)$ have the following properties $F_{k}(q, p)=(-1)^{k-1} F_{k}(p, q), F_{k}(-p,-q)=F_{k}(p, q), G_{k+1}(p, q)=(-1)^{k} H_{k+1}(q, p), H_{k+1}(-p,-q)=-H_{k+1}(p, q)$

And

$$
\begin{align*}
& F_{k}\left(\left.p\right|_{x=-x},\left.q\right|_{x=-x}\right)=\left.(-1)^{k-1} F_{k}(p, q)\right|_{x=-x}, \\
& G_{k}\left(\left.p\right|_{x=-x},\left.q\right|_{x=-x}\right)=\left.(-1)^{k} G_{k}(p, q)\right|_{x=-x},  \tag{4}\\
& H_{k}\left(\left.p\right|_{x=-x},\left.q\right|_{x=-x}\right)=\left.(-1)^{k} H_{k}(p, q)\right|_{x=-x} .
\end{align*}
$$

Since properties (3), (4) of equations (1) depend on the equation number, then reductions of equations of the AKNS hierarchy depend on the equation number as well. In particular, equations (1) except for general reductions

$$
\begin{gathered}
q\left(x, t_{n}\right)=\sigma p^{*}\left(x, t_{n}\right) \\
q\left(x, t_{n}\right)=\sigma p^{*}\left(-x,-t_{n}\right),
\end{gathered}
$$

where $\sigma= \pm 1$, the following reductions are allowed

$$
\begin{align*}
q\left(x, t_{2 n-1}\right) & =\sigma p\left(x,-t_{2 n-1}\right),  \tag{5a}\\
q\left(x, t_{2 n}\right) & =\sigma p\left(x, t_{2 n}\right)  \tag{5b}\\
q\left(x, t_{2 n-1}\right) & =\sigma p^{*}\left(-x, t_{2 n-1}\right),  \tag{5c}\\
q\left(x, t_{2 n}\right) & =\sigma p^{*}\left(-x,-t_{2 n}\right) \tag{5d}
\end{align*}
$$

Naturally, the mixed equations [15] also allow reductions (5).
Equations of the form
$p_{t}+\sum_{k \geq 1} i^{2 k-1} \gamma_{k} H_{2 k}(p, q)=0$
allow reductions (5a), (5c), and equations
$p_{t}+\sum_{k \geq 1} i^{2 k} \gamma_{k} H_{2 k+1}(p, q)=0$
allow reductions (5b), (5d).
In this paper, we consider a class of relatively simple solutions of odd equations of the AKNS hierarchy. The solutions considered by us can be used to obtain more complex solutions using the Darboux transformation.

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## 2. Modulated plane wave

Recall that the algebraic-geometric solution of equations of the AKNS hierarchy for all values of $t_{j}$ satisfies the stationary mixed equation (see, for example, [15])

$$
H_{g}(p, q)+\sum_{k=1}^{g-1} C_{k} H_{g-k}(p, q)=C_{g} p
$$

where $g$ is the genus of the corresponding spectral curve, $C_{j}$ are some constants. Using the method described in [15], one can find the equation of the spectral curve corresponding to a given particular stationary solution.

Suggested $g=2 n-1$ and

$$
p(x, 0)=f(x) \text { and } q(x, 0)=\sigma f(x)
$$

or

$$
p(x, 0)=f(x) \text { and } q(x, 0)=\sigma f^{*}(-x)
$$

from equation (6) we find the function $f(x)$. Then adding the dependence on $t$ by formulas

$$
\begin{equation*}
p\left(x, t_{2 n-1}\right)=f(x) e^{i b t_{2 n-1}}, q\left(x, t_{2 n-1}\right)=\sigma f(x) e^{-i b t_{2 n-1}} \tag{6}
\end{equation*}
$$

or

$$
\begin{equation*}
p\left(x, t_{2 n-1}\right)=f(x) e^{i b t_{2 n-1}}, q\left(x, t_{2 n-1}\right)=\sigma f^{*}(-x) e^{-i b t_{2 n-1}}, \tag{7}
\end{equation*}
$$

we obtain the solution of the mixed equation of the AKNS hierarchy with a reduction (5a) or (5c).
Further, applying the Darboux transformation to the found solutions, it is possible to obtain more complex solutions of the considered mixed equations of the AKNS hierarchy. Naturally, each of the reductions corresponds to its Darboux transformation. Therefore, from the same solution satisfying both reductions, it is possible to obtain further different solutions that will satisfy only one reduction.

For example, making in the coupled nonlinear Schrödinger equation:

$$
\left\{\begin{array}{l}
i p_{t_{1}}+p_{x x}-2 p^{2} q=0,  \tag{8}\\
-i q_{t_{1}}+q_{x x}-2 q^{2} p=0,
\end{array}\right.
$$

replacing (8), we obtain

$$
\begin{equation*}
f_{x x}=2 \sigma f^{3}+b f \tag{9}
\end{equation*}
$$

Multiplying equation (9) by $2 f_{x}$ and integrating, we obtain

$$
\begin{equation*}
\left(f_{x}\right)^{2}=\sigma f^{4}+b f^{2}+c \tag{10}
\end{equation*}
$$

where $c$ is the integration constant. In general, equation (10) is the equation of an elliptic function. Depending on which elliptic Jacobi function is used, one of three solutions of equations (8) with reduction (5a) is obtained. If the solution of the equation (10) satisfies the condition

$$
f^{*}(-x) \equiv f(x)
$$

then the functions (6) are the solution of equations (8) with reduction (5c).

## 3. Elliptic solution I

If $f(x)=A \operatorname{sn}\left[a\left(x-x_{0}\right) ; k^{2}\right]$ then

$$
A=\frac{a k}{\sqrt{\sigma}}, b=-a^{2}\left(k^{2}+1\right), c=\frac{a^{4} k^{2}}{\sigma}
$$

and

$$
\begin{gather*}
p=\frac{a k e^{-i a^{2}\left(k^{2}+1\right) t}}{\sqrt{\sigma}} \operatorname{sn}\left[a\left(x-x_{0}\right) ; k^{2}\right], \\
q=\sqrt{\sigma} a k e^{i a^{2}\left(k^{2}+1\right) t} \operatorname{sn}\left[a\left(x-x_{0}\right) ; k^{2}\right] . \tag{11}
\end{gather*}
$$

Figure 1 shows the real part of the function $p(x, t)(11)$ when $\sigma=1, a=3, k=0,8$


Figure 1. The real part of the function $p(x, t)(12)$ when $\sigma=1, a=3, k=0,8$.

For $\quad a \in \mathrm{R}, 0<k<1, x_{0}=0, \sigma=-1$ the condition is satisfied $f^{*}(-x) \equiv f(x)$.

Therefore, for given parameter values, solution (11) satisfies equations (8) with reduction (5c).

## 4. Elliptic solution II

If $f(x)=A \operatorname{cn}\left[a\left(x-x_{0}\right) ; k^{2}\right]$ then

$$
A=\frac{i a k}{\sqrt{\sigma}}, b=a^{2}\left(2 k^{2}-1\right),{ }_{c=\frac{a^{4} k^{2}\left(k^{2}-1\right)}{\sigma}}^{\sigma}
$$

and

$$
\begin{align*}
p & =\frac{i a k e^{i a^{2}\left(2 k^{2}-1\right) t}}{\sqrt{\sigma}} \operatorname{cn}\left[a\left(x-x_{0}\right) ; k^{2}\right], \\
q & =i \sqrt{\sigma} a k e^{-i a^{2}\left(2 k^{2}-1\right) t} \operatorname{cn}\left[a\left(x-x_{0}\right) ; k^{2}\right] . \tag{12}
\end{align*}
$$

Figure 2 shows the real part of the function $p(x, t)(12)$ when $\sigma=1, a=3, k=0,8$


Figure 2. The real part of the function $p(x, t)$ (12) when $\sigma=1, a=3, k=0,8$.

Solution (12) for $a \in \mathrm{R}, 0<k<1, x_{0}=0, \sigma=-1$ also satisfies equations (8) with reduction (5c), since in this case the condition is also satisfied

$$
f^{*}(-x) \equiv f(x)
$$

## 5. Elliptic solution III

If $f(x)=A \operatorname{dn}\left[a\left(x-x_{0}\right) ; k^{2}\right]$ then

$$
A=\frac{i a}{\sqrt{\sigma}}, b=a^{2}\left(2-k^{2}\right), c=\frac{a^{4}\left(1-k^{2}\right)}{\sigma}
$$

and

$$
\begin{gather*}
p=\frac{i a e^{i a^{2}\left(2-k^{2}\right) t}}{\sqrt{\sigma}} \operatorname{dn}\left[a\left(x-x_{0}\right) ; k^{2}\right], \\
q=i \sqrt{\sigma} a e^{-i a^{2}\left(2-k^{2}\right) t} \operatorname{dn}\left[a\left(x-x_{0}\right) ; k^{2}\right] . \tag{13}
\end{gather*}
$$

Figure 3 shows the real part of the function $p(x, t)(13)$ when $\sigma=1, a=3, k=0,8$
Solution (13) for $a \in \mathrm{R}, 0<k<1, x_{0}=0, \sigma=-1$ also satisfies equations (8) with reduction (5c), since in this case the condition is also satisfied

$$
f^{*}(-x) \equiv f(x)
$$



Figure 3. The real part of the function $p(x, t)(13)$ when $\sigma=1, a=3, k=0,8$.

## 6. Solutions in elementary functions

However, there are two special cases where equation (11) has a solution in elementary functions.
In the first of them the constant $c=0$ and equation (11) is reduced to the simplest differential equation of the first order

$$
\frac{d f}{d x}=f \sqrt{b+\sigma f^{2}} .
$$

Dividing the variables and integrating, we have

$$
x=\int \frac{d f}{f \sqrt{b+\sigma f^{2}}} .
$$

Calculating the integral standing on the right side and expressing the function $f$, we obtain, depending on the sign of $b$, three different cases.

1. If $b=0$, then

$$
f(x)=\frac{1}{\left(x-x_{0}\right) \sqrt{\sigma}}
$$

and

$$
\begin{equation*}
p=\frac{1}{\left(x-x_{0}\right) \sqrt{\sigma}}, q=\frac{\sqrt{\sigma}}{\left(x-x_{0}\right)} . \tag{14}
\end{equation*}
$$

2. If $b=-a^{2}<0$, then

$$
f(x)=\frac{a}{\sqrt{\sigma} \sin \left[a\left(x-x_{0}\right)\right]}
$$

and

$$
\begin{equation*}
p=\frac{a e^{-i a^{2} t}}{\sqrt{\sigma} \sin \left[a\left(x-x_{0}\right)\right]}, q=\frac{\sqrt{\sigma} a e^{i a^{2} t}}{\sin \left[a\left(x-x_{0}\right)\right]} . \tag{15}
\end{equation*}
$$

3. If $b=a^{2}>0$, then

$$
f(x)=\frac{i a}{\sqrt{\sigma} \cosh \left[a\left(x-x_{0}\right)\right]}
$$

and

$$
\begin{equation*}
p=\frac{i a e^{i a^{2} t}}{\sqrt{\sigma} \cosh \left[a\left(x-x_{0}\right)\right]}, q=\frac{i \sqrt{\sigma} a e^{-i a^{2} t}}{\cosh \left[a\left(x-x_{0}\right)\right]} . \tag{16}
\end{equation*}
$$

Figure 4 shows the real part of the function $p(x, t)(16)$ when $\sigma=1, a=3$


Figure 4. The real part of the function $p(x, t)(16)$ when $\sigma=1, a=3$.

It is easy to see that for $a \in \mathrm{R}, x_{0}=0, \sigma=-1$, the solution (16) also satisfies equations (8) with reduction (5c).

In the second special case

$$
c=\frac{b^{2}}{4 \sigma} \neq 0
$$

and equation (10) is also reduced to the first order differential equation

$$
\frac{d f}{d x}=\sqrt{\sigma} f^{2}+\frac{b}{2 \sqrt{\sigma}} \text { or } \frac{d(\sqrt{\sigma} f)}{d x}=(\sqrt{\sigma} f)^{2}+\frac{b}{2} .
$$

By introducing a new variable $\tilde{f}=\sqrt{\sigma} f$, separating the variables and integrating, we have

$$
x=\int \frac{d \tilde{f}}{\left(\tilde{f}^{2}+\frac{b}{2}\right)}
$$

Calculating the integral standing on the right side and expressing the function $h$, we obtain, depending on the sign of $b$ two different cases

1. If $b=2 a^{2}>0$, then

$$
f(x)=\frac{a}{\sqrt{\sigma}} \operatorname{ctg}\left[a\left(x-x_{0}\right)\right]
$$

and

$$
\begin{equation*}
p=\frac{a e^{2 i a^{2} t}}{\sqrt{\sigma}} \operatorname{ctg}\left[a\left(x-x_{0}\right)\right], q=\sqrt{\sigma} a e^{-2 i a^{2} t} \operatorname{ctg}\left[a\left(x-x_{0}\right)\right] . \tag{17}
\end{equation*}
$$

2. If $b=-2 a^{2}>0$, then

$$
f(x)=\frac{a}{\sqrt{\sigma}} \tanh \left[a\left(x-x_{0}\right)\right]
$$

and

$$
\begin{equation*}
p=\frac{a e^{-2 i a^{2} t}}{\sqrt{\sigma}} \tanh \left[a\left(x-x_{0}\right)\right], q=\sqrt{\sigma} a e^{2 i a^{2} t} \tanh \left[a\left(x-x_{0}\right)\right] . \tag{18}
\end{equation*}
$$

Figure 5 shows the real part of the function $p(x, t)(18)$ when $\sigma=1, a=3$


Figure 5. The real part of the function $p(x, t)(18)$ when $\sigma=1, a=3$.

It is easy to check that when $a \in \mathrm{R}, x_{0}=0, \sigma=-1$, the solution (18) also satisfies equations (8) with reduction (5c).

## 7. Conclusion

In conclusion, we note that the solutions found above satisfy the subsequent odd equations of the AKNS hierarchy, but with a different coefficient before the variable $t$. In particular, substituting (6) into equation

$$
i p_{l_{3}}+H_{4}(p, q)=0
$$

and simplifying, we obtain (for $b=b_{3}$ ).

$$
\begin{equation*}
-b_{3} f-6 \sigma^{2} f^{5}+10 \sigma f f_{x}^{2}+10 \sigma f^{2} f_{x x}-f_{x x x x}=0 . \tag{19}
\end{equation*}
$$

Suppose that the function $f(x)$ satisfies equations (9) and (10).
Then

$$
f_{x x x}=6 \sigma f^{2} f_{x}+b f_{x}
$$

and

$$
f_{x x x}=12 \sigma f f_{x}^{2}+6 \sigma f^{2} f_{x x}+b f_{x x}
$$

or

$$
\begin{equation*}
f_{x x x}=24 \sigma^{2} f^{5}+20 \sigma b f^{\beta}+12 \sigma c f+b^{2} f . \tag{20}
\end{equation*}
$$

Substituting (20), (9) and (10) into (19) and simplifying, we obtain

$$
b_{3}=-b^{2}-2 \sigma c .
$$

Naturally, the Darboux transformations of the solutions found will also satisfy all odd equations of the AKNS hierarchy with corresponding reductions and their mixed forms.

## References

[1] Ablowitz M J, Feng B F, Luo X D and Musslimani Z H 2018 Reverse space-time nonlocal Sine-Gordon/Sinh-Gordon equations with nonzero boundary conditions Stud. Appl. Math.
[2] Ablowitz M J and Musslimani Z H 2013 Integrable nonlocal nonlinear Schrodinger equation Phys. Rev. L. 110064105
[3] Ablowitz M J and Musslimani Z H 2014 Integrable discrete PT symmetric model Phys. Rev. E 90032912
[4] Ablowitz M J and Musslimani Z H 2016 Inverse scattering transform for the integrable nonlocal nonlinear Schr" odinger equation Nonlinearity 29 915-46
[5] Ablowitz M J and Musslimani Z H 2017 Integrable nonlocal nonlinear equations Stud. Appl. Math. 139 7-59
[6] Cao Y, Malomed B and He J 2018 Two (2+1)-dimensional integrable nonlocal nonlinear Schrodinger equations: Breather, rational and semi-rational solutions Chaos Solitons and Fractals 114 99-107
[7] Cao Y, Rao J, Mihalache D and He J 2018 Semi-rational solutions for the (2+1)- dimensional nonlocal Fokas system Appl. Math. L 80 27-34
[8] Chen K, Deng X, Lou S and Zhang D J 2018 Solutions of Nonlocal Equations Reduced from the AKNS Hierarchy Stud. Appl. Math. 141 113-41
[9] Gurses M and Pekcan A 2019 Nonlocal modified KdV equations and their soliton solutions by Hirota method Commun Nonlinear Sci. Numer. Simulat. 67 427-48
[10] Horikis T P and Ablowitz M J 2017 Rogue waves in nonlocal media Phys. Rev. E 95042211
[11] Liu D Y and Sun W R 2018 Rational solutions for the nonlocal sixth-order nonlinear Schrodinger equation Appl. Math. L 84 63-9
[12] Liu W and Li X 2018 General soliton solutions to a (2+1)-dimensional nonlocal nonlinear Schrodinger equation with zero and nonzero boundary conditions Nonlinear Dynamics 93 721-31
[13] Manikandan K, Vishnu Priya N, Senthilvelan M and Sankaranarayanan R 2018 Deformation of dark solitons in a PT-invariant variable coefficients nonlocal nonlinear Schrodinger equation Chaos 28083103
[14] Matveev V B and Smirnov A O 2016 Solutions of the Ablowitz-Kaup-Newell-Segur hierarchy equations of the "rogue wave" type: a unified approach Theor. Math. Phys. 186 156-82
[15] Matveev V B and Smirnov A O 2018 AKNS hierarchy, MRW solutions, Pn breathers, and beyond J. Math. Phys. 59091419
[16] Qian C, Rao J, Mihalache D and He J 2018 Rational and semi-rational solutions of the y-nonlocal Davey-Stewartson I equation Computers and Mathematics with Applications 75 3317-30
[17] Rao J, Zhang Y, Fokas A and He J 2018 Rogue waves of the nonlocal Davey-Stewartson I equation Nonlinearity 31 4090-107
[18] Tang X Y, Liang Z F and Hao X Z 2018 Nonlinear waves of a nonlocal modified KdV equation in the atmospheric and oceanic dynamical system Commun Nonlinear Sci. Numer. Simulat. 60 62-71
[19] Vinayagam P, Radha R, Al Khawaja U and Ling L 2018 New classes of solutions in the coupled PT symmetric nonlocal nonlinear Schrodinger equations with four wave mixing Commun Nonlinear Sci. Numer. Simulat. 59 387-95
[20] Yang B and Chen Y 2018 Reductions of Darboux transformations for the PT-symmetric nonlocal Davey-Stewartson equations Appl. Math. L 82 43-9
[21] Yang Y, Suzuki T and Cheng X 2020 Darboux transformations and exact solutions for the integrable nonlocal Lakshmanan-Porsezian-Daniel equation Appl. Math. L 99105998
[22] Yang Z J, Zhang S M, Li XL and Pang Z G 2081 Variable sinh-Gaussian solitons in nonlocal nonlinear Schrodinger equation Appl. Math. L 82 64-70
[23] Zhang H Q and Gao M 2018 Rational soliton solutions in the parity-time-symmetric nonlocal coupled nonlinear Schrodinger equations Comm. Nonlin. Sci. and Num. Sim. 63 253-60
[24] Zhang Q, Zhang Y and Ye R 2019 Exact solutions of nonlocal Fokas-Lenells equation Appl. Math. L 98 336-43
[25] Zhou Z X 2018 Darboux transformations and global solutions for a nonlocal derivative nonlinear Schrodinger equation Commun Nonlinear Sci. Numer. Simulat. 62 480-8

