

PAPER • OPEN ACCESS

Dynamics of oscillator populations with disorder in the coupling phase shifts

To cite this article: Arkady Pikovsky and Franco Bagnoli 2024 *New J. Phys.* **26** 023054

View the [article online](#) for updates and enhancements.

You may also like

- [Study on neural entrainment to continuous speech using dynamic source connectivity analysis](#)
Kai Yang, Shuang Wu, Di Zhou et al.
- [HYDRODYNAMIC SIMULATIONS OF H ENTRAINMENT AT THE TOP OF He-SHELL FLASH CONVECTION](#)
Paul R. Woodward, Falk Herwig and Pei-Hung Lin
- [Simultaneous iterative learning control of mode entrainment and error field](#)
W. Choi and F.A. Volpe

**PAPER**

Dynamics of oscillator populations with disorder in the coupling phase shifts

OPEN ACCESS**RECEIVED**

30 December 2023

REVISED

11 February 2024

ACCEPTED FOR PUBLICATION

19 February 2024

PUBLISHED

28 February 2024

Original Content from
this work may be used
under the terms of the
[Creative Commons
Attribution 4.0 licence](#).

Any further distribution
of this work must
maintain attribution to
the author(s) and the title
of the work, journal
citation and DOI.

Arkady Pikovsky^{1,*}  and Franco Bagnoli^{2,3} ¹ Institute of Physics and Astronomy, University of Potsdam, Karl-Liebknecht-Str. 24/25, 14476 Potsdam-Golm, Germany² Department of Physics and Astronomy and CSDC, University of Florence, via G. Sansone 1, Sesto Fiorentino 50019, Italy³ INFN, section of Florence, via G. Sansone 1, Sesto Fiorentino 50019, Italy

* Author to whom any correspondence should be addressed.

E-mail: pikovsky@uni-potsdam.de and franco.bagnoli@unifi.it**Keywords:** synchronization, frequency entrainment, disorder, Kuramoto model**Abstract**

We study populations of oscillators, all-to-all coupled by means of quenched disordered phase shifts. While there is no traditional synchronization transition with a nonvanishing Kuramoto order parameter, the system demonstrates a specific order as the coupling strength increases. This order is characterized by partial phase locking, which is put into evidence by the introduced novel correlation order parameter, which is shown to grow monotonically with the coupling strength, and via frequency entrainment by following concentration of the oscillators frequencies. Simulations with phase oscillators, Stuart–Landau oscillators, and chaotic Roessler oscillators demonstrate similar scaling of the correlation order parameter with the coupling and the system size and also similar behavior of the frequencies with maximal entrainment (at which the standard deviation of the frequencies is reduced by a factor close to four) at some finite coupling.

1. Introduction

Synchronization phenomena in populations of oscillators have been widely studied since the pioneering works of Winfree [1] and Kuramoto [2]. The Kuramoto model and its modifications, based on the phase dynamics, serve as paradigmatic models of synchronization (with a possibility, in many situations, of a fully analytical treatment). The relevance of the findings within phase dynamics is confirmed by studies of realistic models of limit-cycle oscillators, where the amplitude dynamics is also included. If limit cycles are close to a Hopf bifurcation point, a universal normal form description can be adopted, which is expressed as an ensemble of Stuart–Landau oscillators, which also exhibits the synchronization transition [3, 4]. Moreover, Kuramoto-type synchronization can be also observed in populations of chaotic oscillators [5].

Effects of disorder on the synchronization properties of large populations of coupled oscillators are a subject of intensive current research. Randomness can be restricted to properties of individual units; e.g. already in the original formulation of the Kuramoto model [2], a distribution of natural frequencies was assumed. Also, the interactions between the oscillators can hold randomness, e.g. in the form of a random (weighted or unweighted) network [6–8]. In many cases, in the thermodynamic limit of an infinite number of units, a theory of synchronization transition can be constructed where the randomness is smeared out, and the averaged coupling, if attractive, leads to a synchronization transition [9]. In some cases, disorder in the system may even play a constructive role, facilitating synchrony [10].

A special class of problems appears if the random coupling is maximally frustrating, i.e. if it disappears after a naive averaging. To the best of our knowledge, the first model of such type has been suggested by Daido [11], who considered a population of oscillators with mutual couplings sampled from a Gaussian distribution with zero mean. Although such a model does not exhibit synchronization as a transition to a state with a non-vanishing order parameter, it demonstrates some entrainment properties. By exploiting an analogy with spin glasses [12], he found regimes with a slow relaxation to equilibrium. Furthermore, Daido found a volcano transition, where a distribution of random coupling inputs on oscillators changes its form

from one with a maximum at zero to one with a maximum at a finite distance from zero. The Daido model and its modifications have been since intensively studied [13–20].

This paper focuses on another type of randomness in a population of oscillators, where couplings are characterized by random phase shifts [21]. If the distribution of the random shifts is nonuniform, one obtains certain renormalized averaged interaction in the thermodynamic limit, which begets a standard synchronization transition [9]. However, no standard synchronization transition occurs for a maximally frustrating disorder with a uniform distribution of the phase shifts in $[0, 2\pi)$. As we will show in this paper, a certain ordered state nevertheless appears, and the present work aims to characterize it. We will see that the effects under disorder in the phase shifts are quite different from those reported in [13–20] for disorder in coupling strengths; in particular we do not observe a volcano transition and a glassy state in our setup.

The paper is organized as follows. In section 2, we will formulate the three basic models that will be explored numerically: phase oscillators, Stuart–Landau limit cycle oscillators, and Roessler chaotic oscillators. In section 3, we introduce the tools by which the regimes will be characterized. In particular, we introduce a novel correlation order parameter and discuss its relation to the standard Kuramoto order parameter. Detailed numerical studies of populations of phase oscillators are presented in section 4. Similar simulations for Stuart–Landau and Roessler oscillations reported in section 5 show, on the one hand, that the effects found for phase oscillators are reproduced for these more realistic models, while, on the other hand, novel features related to the amplitude dynamics appear. We conclude in section 6 with a comparison to Daido’s disorder model.

2. Basic models

In this section, we take three popular models of global coupling (Kuramoto–Sakaguchi model [22]; system of Stuart–Landau oscillators [3]; system of chaotic Roessler oscillators [5]), widely used for exploration of the collective dynamics and of the synchronization transition, and extend them to the case of disorder in the phase shifts.

2.1. Phase oscillators

The standard Kuramoto–Sakaguchi model of N oscillators with all-to-all couplings reads

$$\dot{\varphi}_k = \omega_k + f \sum_{j \neq k} \sin(\varphi_j - \varphi_k - \alpha), \quad k = 1, \dots, N. \quad (1)$$

Here φ_k are the phases of the oscillating units, and ω_k are their natural frequencies. Parameter f is the coupling strength (for future compatibility, we do not normalize it by N). Here, all mutual coupling terms are assumed to include the same phase shift α . We introduce a disorder in these phase shifts by assuming that phase shifts α_{jk} are independent, identically distributed random numbers (cf [21]):

$$\dot{\varphi}_k = \omega_k + f \sum_{j \neq k} \sin(\varphi_j - \varphi_k - \alpha_{jk}). \quad (2)$$

Furthermore, in this paper we assume the maximally possible disorder, where α_{jk} are uniformly distributed in $[0, 2\pi)$ (for non-uniform distributions, see theory in [9]). Furthermore, we mostly concentrate on the asymmetric case $\alpha_{jk} \neq \alpha_{kj}$.

The main physical motivation behind the disordered phase shifts is that such shifts can be associated with time delays in coupling. In particular, if the frequencies of all oscillators are close to a central frequency ω_0 , then the shifts can be expressed via the coupling time delays τ_{jk} as $\alpha_{jk} = \omega_0 \tau_{jk}$, where ω_0 is the central frequency. The basic assumption is that the time delays are smaller than the characteristic time of slow variations of the phases; see [23] for a precise formulation. Uniformity (or almost uniformity) of the distribution of the phase shifts corresponds to a broad distribution of the time delays (a situation where delays in coupling obey a gamma distribution has been treated in [24]). We mention here that in addition to time delays in coupling, other factors may influence the phase shifts. For example, transmitting a signal from the driving to the driven unit may include filters with random characteristics.

We mention here that following the pioneering work of Daido [11], one often considers an implementation of disorder in the coupling, where the randomness is not in the phase shifts but in the coupling constants, see, e.g. [13–19]. We discuss a possible relation to these models in section 6; here we only mention that the phenomenologies in the cases of the coupling strength disorder and the phase shift disorders are different.

Below, we consider a Gaussian distribution of natural frequencies ω_k . Since one can rescale time in equation (2), we will set the variance of this distribution to one. Furthermore, to avoid potential finite-size effects due to realizations of the frequencies, we assign ω_k not via random sampling, but via deterministic sampling using quantiles of the Gaussian distribution (see discussion on the differences between random and deterministic samplings in [25]). Below, we will also consider a version of equation (2) where all the oscillators have the same natural frequency (which is set to zero, $\omega_k = 0$), but are subject to independent Gaussian white noise forces:

$$\begin{aligned}\dot{\varphi}_k &= \eta_k(t) + f \sum_{j \neq k} \sin(\varphi_j - \varphi_k - \alpha_{jk}) , \\ \langle \eta_k \rangle &= 0, \langle \eta_k(t) \eta_l(t') \rangle = 2\delta_{kl} \delta(t - t') .\end{aligned}\quad (3)$$

Here again, the intensity of noise is fixed via the time rescaling.

In the case of deterministic identical oscillators, one can rescale the coupling parameter to one, and the problem is reduced to

$$\dot{\varphi}_k = \sum_{j \neq k} \sin(\varphi_j - \varphi_k - \alpha_{jk}) , \quad (4)$$

where the only parameter is the size of the ensemble N .

2.2. Stuart-Landau oscillators (SLO)

A SLO is a prototypic model of a self-sustained oscillator, which includes both the phase and the amplitude dynamics. We consider in this paper a network of N SLOs with random phase shifts in coupling

$$\dot{a}_k = i\omega_k a_k + \mu a_k (1 - |a_k|^2) + f \sum_{j \neq k} e^{-i\alpha_{jk}} a_j . \quad (5)$$

Here a_k are complex amplitudes of the units, and parameter μ defines the relaxation rate of the amplitudes (μ^{-1} is the characteristic relaxation time of the amplitude dynamics). In the limiting case of strongly stable amplitude, $\mu \rightarrow \infty$, the dynamics of the phases $\varphi_k = \arg(a_k)$ follows equation (2); otherwise, the dynamics of the amplitudes influences that of the phases.

2.3. Chaotic Roessler-type oscillators

A chaotic Roessler oscillator [26] (RO) possesses a well-defined phase variable; and an ensemble of such systems with a Kuramoto–Sakaguchi-type coupling equation (1) demonstrates a phase synchronization transition similar to the Kuramoto transition [5]. Below we explore the effect of randomness in the phase shifts in the coupling. For this purpose, we slightly modify the original Roessler system to obtain equations close to the dynamics of the SLO:

$$\begin{aligned}\dot{a} &= \omega (ia + da - z) , \\ \dot{z} &= \omega (b + z(\operatorname{Re}(a) - c)) .\end{aligned}\quad (6)$$

Here, a is the complex ‘amplitude’ of the oscillations (related to ‘usual’ variables of the Roessler system as $a = x + iy$), and z is a variable, the dynamics of which ensures chaotic ‘saturation’ of the unstable oscillations of a . We introduce a parameter ω responsible for the natural frequency as a common factor to avoid difficulties related to possible bifurcations ‘order-chaos’ in dependence on such a parameter (cf a similar approach in [27]). The phase variable for the Roessler-type oscillator can be defined as $\varphi_k = \arg(a_k)$. The effective dynamics of the phase is diffusive due to chaotic variations of the amplitudes.

We now introduce a global coupling in a population with different natural frequencies ω_k like in (5)

$$\begin{aligned}\dot{a}_k &= \omega_k (i a_k + d a_k - z_k) + f \sum_{j \neq k} e^{-i\alpha_{jk}} a_j , \\ \dot{z}_k &= \omega_k (b + z_k (\operatorname{Re}(a_k) - c)) .\end{aligned}\quad (7)$$

3. Characterizations of the collective dynamics

There are two ways to characterize synchronization in populations of oscillators: To follow the phase locking or the frequency entrainment. In many situations, these aspects have significant overlap, but there are also

cases where phase locking is accompanied by frequency anti-entrainment [28, 29]. We will see that the two characterizations deliver complementary characterizations for the abovementioned ensembles.

3.1. Correlation order parameter for phase locking

Phase locking in a population of oscillators without disorder is typically characterized by the complex Kuramoto order parameter (complex KOP)

$$Z(t) = \frac{1}{N} \sum_{k=1}^N e^{i\varphi_k(t)} .$$

At the synchronization transition, the absolute value of Z starts to grow, and large values of $|Z|$ indicate that the phases are concentrated in a small range (compared to 2π). As we will see below, this KOP is nearly vanishing (up to finite-size fluctuations) in the presence of disorder. Therefore, we suggest to characterize relations between the phases with a complex phase correlation order parameter (complex COP) defined for the systems equations (2)–(4) as (cf [30, 31])

$$\begin{aligned} Q(t) &= \frac{1}{N} \sum_k Q_k(t) , \\ Q_k(t) &= \frac{1}{N-1} \sum_{j \neq k} e^{i(\varphi_j - \varphi_k - \alpha_{jk})} . \end{aligned} \quad (8)$$

We notice that the imaginary part of Q_k is, up to a factor, the force acting on the oscillator k . The real part of Q_k will be more important. This quantity characterizes the closeness (on average) of the phase φ_k to the value $\varphi_j + \alpha_{jk}$, which would be the locked state if only one coupling $j \rightarrow k$ would be present (but this value is, of course, frustrated in the presence of all couplings).

Let us first demonstrate that complex COP and complex KOP are directly related in the case without disorder $\alpha_{jk} = 0$, i.e. for the standard Kuramoto setup:

$$|Z|^2 = \frac{1}{N} \sum_{k=1}^N e^{-i\varphi_k(t)} \frac{1}{N} \sum_{j=1}^N e^{i\varphi_j(t)} = \frac{1}{N^2} \sum_{j,k} e^{i(\varphi_j - \varphi_k)} = \frac{(N-1)Q + 1}{N} . \quad (9)$$

For large N , one thus has $|Z|^2 \approx Q$ (note that in the ordered case, Q is always real).

Below, in the applications of the complex COP to the disordered populations, we will perform additional averaging of Q over time (for a particular sample of phase shifts), and over realizations of disorder, and separate real and imaginary parts of the resulting complex quantity as

$$\langle Q \rangle = C + iD . \quad (10)$$

We will refer to C as the real correlation order parameter (real COP). Correspondingly, in a slight discrepancy with a usual definition of the (real) Kuramoto order parameter, we will define it as $|Z|^2$, averaged over time and over the realizations of disorder:

$$R = \langle |Z|^2 \rangle . \quad (11)$$

The COP Q can be directly applied to populations of SLOs (5) and ROs (7), provided the phases $\exp[i\varphi_k] = a_k/|a_k|$ are used as the observables. One can also define the corresponding correlation order parameters that include full amplitudes

$$\begin{aligned} Q^{(a)}(t) &= \frac{1}{N} \sum_k Q_k^{(a)}(t) , \\ Q_k^{(a)}(t) &= \frac{1}{N-1} \sum_{j \neq k} a_j a_k^* e^{-i\alpha_{jk}} , \end{aligned} \quad (12)$$

where a_j are complex amplitudes either of SLOs or of ROs. Similar to (10), $\langle Q^{(a)} \rangle = C^{(a)} + iD^{(a)}$.

3.2. Frequency entrainment

The definition of the average observed frequency for each oscillator in the population is straightforward (of course, the average includes only time average, not averaging over the realizations of disorder):

$$\Omega_k = \langle \dot{\varphi}_k \rangle . \quad (13)$$

This definition will be applied to all types of oscillators (phase oscillators, SLOs, and ROs). In the Kuramoto system without disorder, the spread of frequencies decreases beyond the synchronization transition. Moreover, in the deterministic case, a portion (growing with the coupling strength) of oscillators is perfectly entrained so that their frequencies coincide.

We notice here that the difference between the observed frequency Ω_k (13) and the natural frequency ω_k in (2) can be expressed as the average of the imaginary part of the local complex COP ($\Omega_k - \omega_k$) $\propto \langle \text{Im}(Q_k) \rangle$. We also note that the entrainment of frequencies is a relevant characterization if the intrinsic frequencies ω_k are different; for the noisy identical oscillators without the disorder, the frequencies are always equal.

4. Phase oscillators

In this section, we apply the introduced above characterizations of synchrony in a population of oscillators to the phase oscillators models equations (2)–(4).

4.1. Phase oscillators with a Gaussian distribution of frequencies

In exploring model (2) for different system sizes N , we select the frequencies according to a Gaussian distribution with unit variance deterministically as corresponding quantiles. This is done to avoid additional diversity due to the random sampling of frequencies. The only parameters of the model are the coupling strength f and the system size N . The phase shifts in each run are chosen randomly according to a uniform distribution $0 \leq \alpha_{jk} < 2\pi$, an averaging is performed over time and over different samples of phase shifts.

Figure 1 presents the characterization in terms of the KOP and of the COP. The main finding is that while the value of the KOP R defined in equation (11) exhibits finite size fluctuations $\sim N^{-1}$, which are practically independent on the coupling strength f (panel (c)), the value of the real COP C defined in equation (10) exhibits a monotonic growth with a saturation at large f (panel (a)). At the same time, the imaginary COP D fluctuates close to zero (panel (c)).

In figure 1(b) we also demonstrate, that the COP C fulfills a scaling relation

$$C(f, N) = N^{-1/2} c \left(f N^{1/2} \right) . \quad (14)$$

Remarkably, the COP demonstrates a continuous transition without any particular threshold value of the coupling.

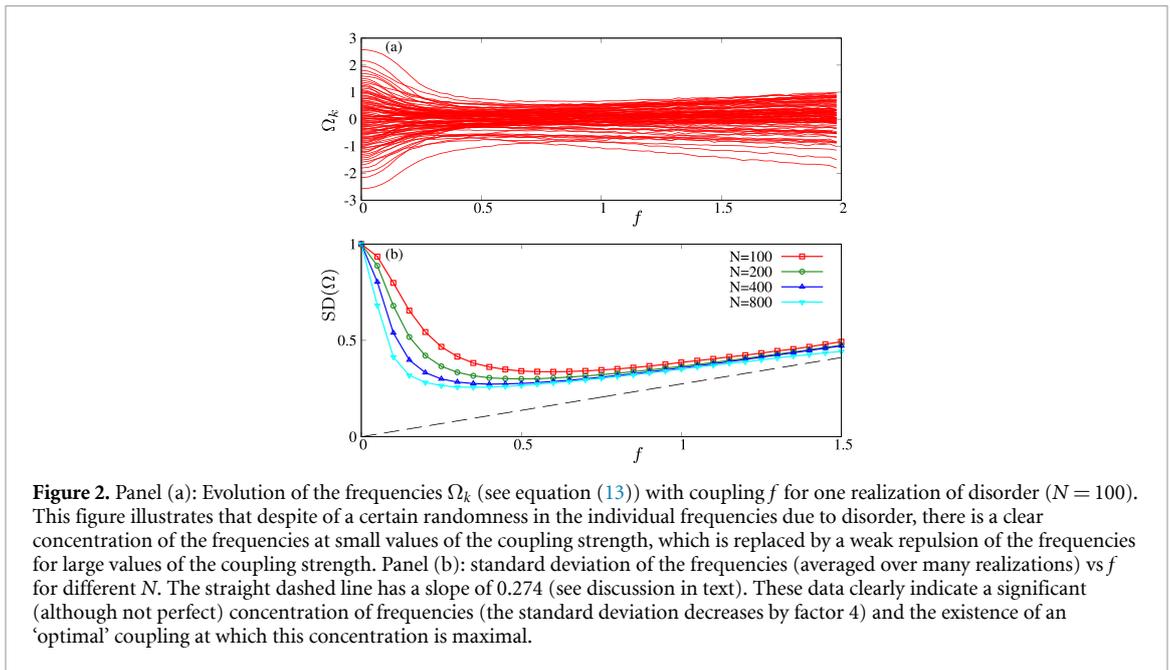
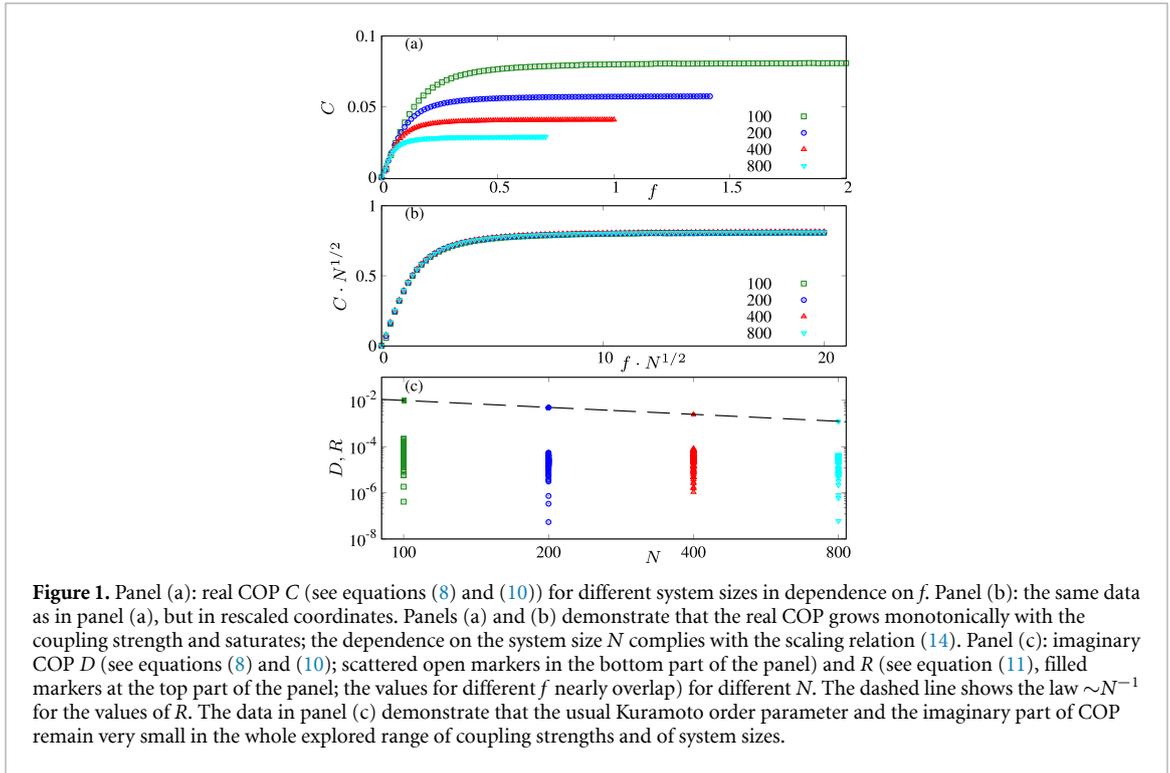
Next, in figure 2, we explore the properties of the frequency entrainment. We show in panel (a) the dependencies of frequencies Ω_k according to equation (13) for one realization of disorder in dependence on the coupling strength f , and in panel (b) the standard deviations of the frequencies from the mean value, averaged for several realizations of disorder. One can see that at small coupling strengths f , the dispersion of frequencies decreases (from the value 1, defined according to the distribution of natural frequencies), but after reaching a minimum, it starts to increase again, nearly linear in f . We will discuss this feature and explain the value of the slope of the dashed line in more detail in section 4.3 below.

4.2. Noisy phase oscillators

Here, we apply the same analysis as in section 4.1 above to noisy oscillators described by equation (3). First, we show in figure 3 the results for the order parameters characterizing phase coherence. One can see that there are no qualitative differences with the case of deterministic oscillators (figure 1). However, there are differences in the behavior of the frequencies. Since the oscillators are identical, the observed frequencies Ω_k in the absence of coupling coincide. However, with the onset of the coupling, such differences do appear, and, as figure 4 shows, the standard deviation grows linearly with f . This amplification corresponds to the increase of the dispersion of the frequencies for non-identical oscillators equation (2) at large coupling strengths f (figure 2(b)).

4.3. Identical phase oscillators

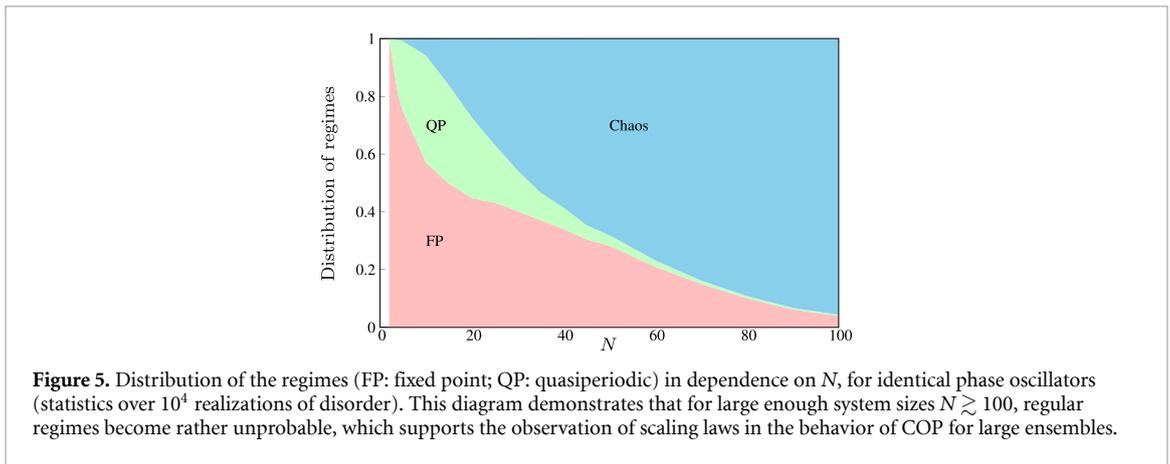
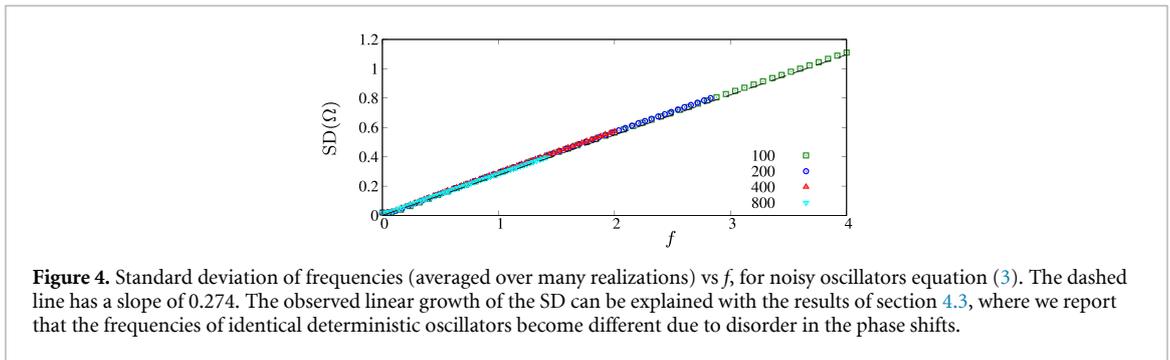
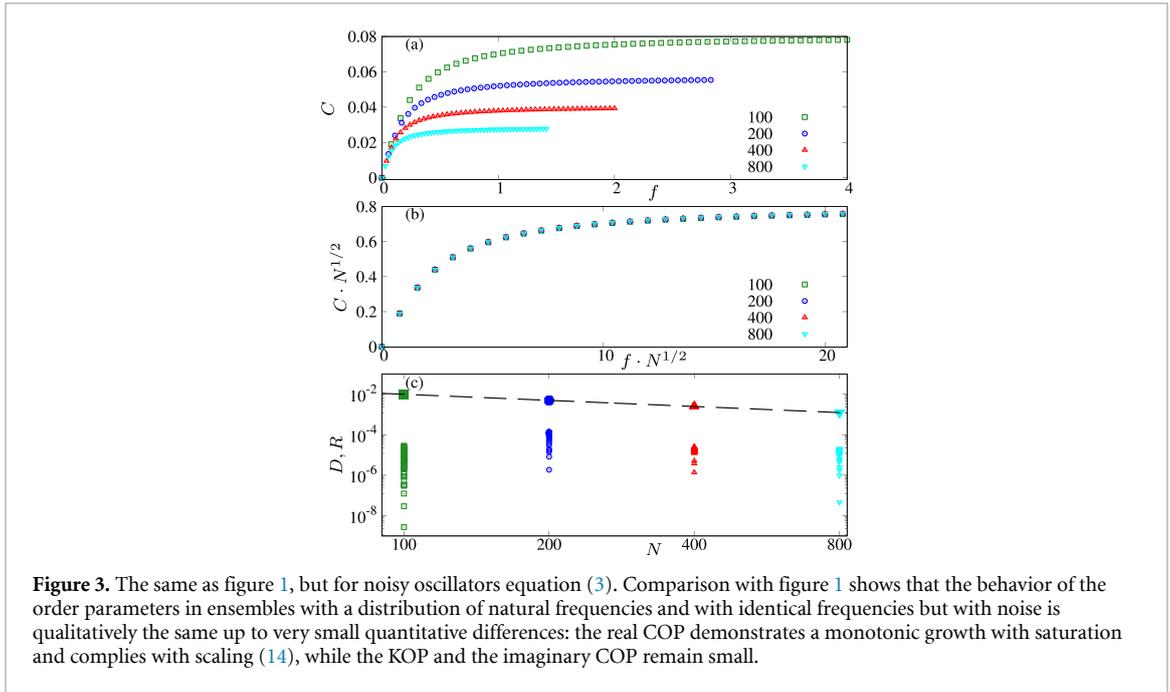
The results of sections 4.1 and 4.2 above show that at large coupling strengths f , a universal regime appears, at which the real COP saturates at a value $C_{sat} \approx 0.8N^{-1/2}$, and the dispersion of the frequencies grows



proportional to f . In this state, one can neglect differences between the oscillators (either due to their different natural frequencies or due to independent noise terms) and study a population of identical deterministic oscillators as defined in equation (4). Because now the time is rescaled according to the coupling strength, a constant dispersion of the observed frequencies will explain linear growth of this dispersion with f in setups where the coupling strength is not rescaled, i.e. in figures 2(b) and 4.

The only parameter of equation (4) is the population size N , what makes a statistical evaluation of different characteristics easier.

First, we would like to characterize the dynamical regimes. Computation of the maximal Lyapunov exponent for the deterministic system (4) allows us to distinguish chaos and regular states. Because the system (4) is invariant with respect to the shift of all phases, one Lyapunov exponent is always zero. To distinguish whether a state with vanishing maximal Lyapunov exponent is a synchronous one where all the frequencies coincide or a quasiperiodic state where the frequencies are different, we calculated in each run



the spread of the observed frequencies (13) as $\Delta\Omega = \Omega_{\max} - \Omega_{\min}$. Regimes where this spread nearly vanishes (we used a threshold 10^{-4} as a practical criterion) were identified as steady states (in some reference frame rotating with a common frequency of the oscillators). Otherwise, the states with nearly vanishing maximal Lyapunov exponent and $\Delta\Omega > 10^{-4}$ were identified as quasiperiodic regimes. The probabilities of observing three possible regimes depend on the ensemble size as depicted in figure 5. For $N = 2$, for all phase shifts, the oscillators synchronize (as the theory of coupled oscillators predicts, because there is no frustration for two oscillators). At the same time, for larger N , quasiperiodic and chaotic regimes become possible. Finally, for large $N \gtrsim 100$, chaotic regimes dominate.

Next, we explored the statistics of the observed frequencies equation (13) at large N , and found that the distribution of these frequencies (when averaged over many realizations of disorder) is Gaussian with high accuracy. Furthermore, the variance is practically the same in a large range of N . This observation explains the linear behavior of the standard deviations of the frequencies in figures 2(b) and 4, and also the fact that these standard deviations for different N nearly coincide at a value ≈ 0.274 . This gives a prediction $\text{Std}(\Omega) \sim 0.274f$ for large f in the presence of a distribution of natural frequencies or in the presence of noise; this prediction is shown with dashed lines in figures 2(b) and 4.

Finally, we studied the statistics of the complex COP Q_k defined in equation (8). We have found that the values $\text{Re}(Q_k)$ and $\text{Im}(Q_k)$ are practically uncorrelated. The average of $\text{Im}(Q_k)$ is nearly zero, while the average of $\text{Re}(Q_k)$ gives the averaged COP C , as discussed above. The variances of $\text{Re}(Q_k)$ and $\text{Im}(Q_k)$ are found to be $\sim N^{-1}$. Remarkably, the distribution of $\text{Re}(Q_k)$ is rather close to a Gaussian one, with small skewness values and excess kurtosis values. The distribution of $\text{Im}(Q_k)$ is nearly symmetric with a small skewness, but the excess kurtosis values of ≈ 3 indicate deviations from a Gaussian distribution.

5. Stuart-Landau and Roessler-type oscillators

This section aims to show that the essential effects described above for the phase oscillators are also observed for SLOs and ROs. We do not repeat all the numerical setups but focus on ensembles with a distribution of natural frequencies without noise.

5.1. SLOs

SLOs have a parameter μ that governs the stability of the limit cycle. For $\mu \rightarrow \infty$, a reduction to the phase dynamics equation (2) becomes exact. Thus, we focus mainly on additional effects appearing at small μ .

We start with presenting in figure 6 the evolution of the frequencies for an ensemble of SLOs (5), where natural frequencies are sampled according to a Gaussian distribution with variance 1. One can see that the evolution of standard deviation is very similar to that for phase oscillators (cf figure 2): entrainment at small values of coupling and growth of the standard deviation for strong coupling. The dependence of these effects on the parameter μ is only quantitative: for small values of μ , the minimal standard deviation is achieved at smaller values of f (see in figure 6 the curves with an open square marker, corresponding to $\mu = 0.1$).

As discussed above, for the COP, we now have two variants: one is based on the transformation to the phases and thus is mostly close to the COP for the phase oscillators (equation (8)); another one is based on the cross-correlations of the amplitudes (equation (12)). Both these COPs are depicted in figure 7, and they demonstrate very different behaviors. While the phase-COP is practically μ -independent, the amplitude-COP grows strongly with the coupling if μ is small. We explain this by noticing that the linear coupling of the SLOs affects the stability of the equilibrium $a_k = 0$ in a population of SLOs. Indeed, equation (5) for small amplitudes a_k can be written as a linear system

$$\dot{a}_k = i\omega_k a_k + \mu a_k + f \sum_{j \neq k} e^{-i\alpha_{jk}} a_j, \quad (15)$$

and stability is affected by the random coupling matrix $e^{-i\alpha_{jk}}$. According to the circular law [32], for large N , the eigenvalues of this matrix are uniformly distributed in a disk with radius \sqrt{N} around the origin. Thus, the largest real part of the eigenvalues in equation (15) can be estimated as $\tilde{\mu} \approx \mu + f\sqrt{N}$. If we assume that this instability saturates at amplitude $|a|^2 \approx \tilde{\mu}/\mu$, we conclude that the average amplitude of oscillations, for large couplings, grows $\sim f\sqrt{N}/\mu$. This linear growth corresponds to what is observed in figure 7, and explains why the amplitude-COP is not a proper quantity to characterize phase coherence in the ensemble for small values of μ .

5.2. Roessler oscillators

Here, we report on the effects of coupling with disorder in the phase shifts on the Roessler system (7). The natural frequencies ω_k have been sampled from Gaussian distribution with mean value 1 (because the RO is not rotationally invariant, one cannot set the mean frequency to zero by a transformation to a rotating reference frame, as it is possible for the phase oscillators and the SLOs) and standard deviation 0.1. Other parameters were set to $d = 0.09$, $b = 0.4$, $c = 8.5$. We also note that because the disordered coupling enhances instability of the oscillations, for large enough values of f , no finite regimes in system (7) have been observed

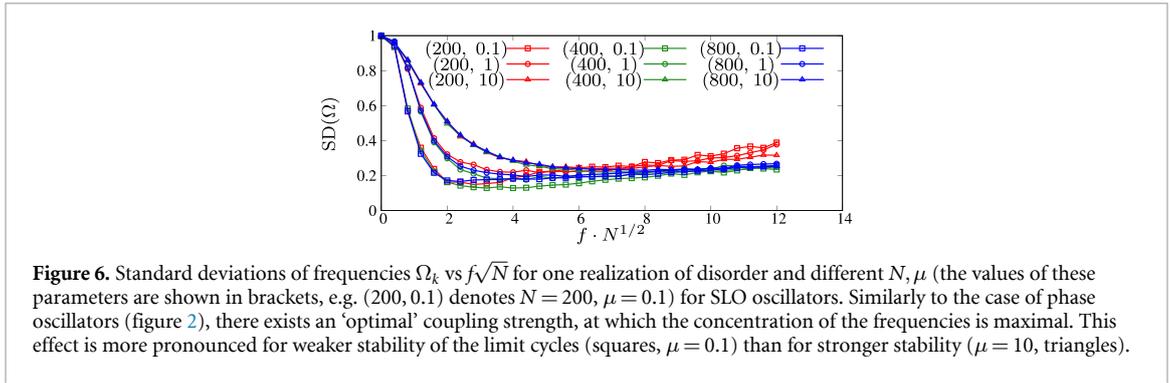


Figure 6. Standard deviations of frequencies Ω_k vs $f\sqrt{N}$ for one realization of disorder and different N, μ (the values of these parameters are shown in brackets, e.g. (200, 0.1) denotes $N = 200, \mu = 0.1$) for SLO oscillators. Similarly to the case of phase oscillators (figure 2), there exists an ‘optimal’ coupling strength, at which the concentration of the frequencies is maximal. This effect is more pronounced for weaker stability of the limit cycles (squares, $\mu = 0.1$) than for stronger stability ($\mu = 10$, triangles).

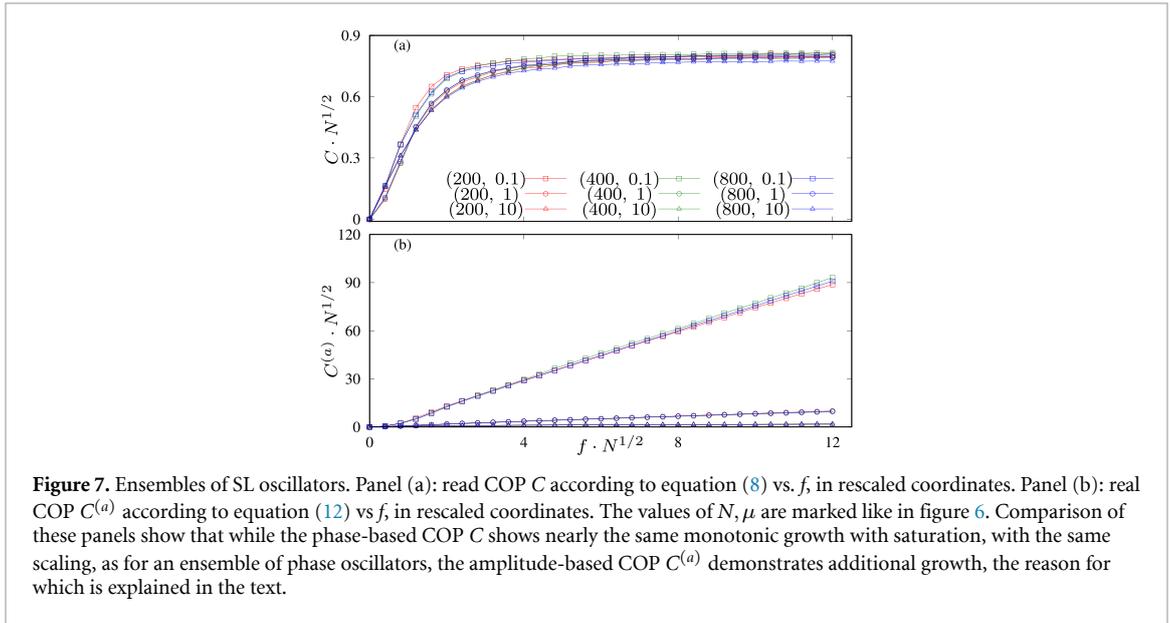


Figure 7. Ensembles of SL oscillators. Panel (a): read COP C according to equation (8) vs. f , in rescaled coordinates. Panel (b): real COP $C^{(a)}$ according to equation (12) vs f , in rescaled coordinates. The values of N, μ are marked like in figure 6. Comparison of these panels show that while the phase-based COP C shows nearly the same monotonic growth with saturation, with the same scaling, as for an ensemble of phase oscillators, the amplitude-based COP $C^{(a)}$ demonstrates additional growth, the reason for which is explained in the text.

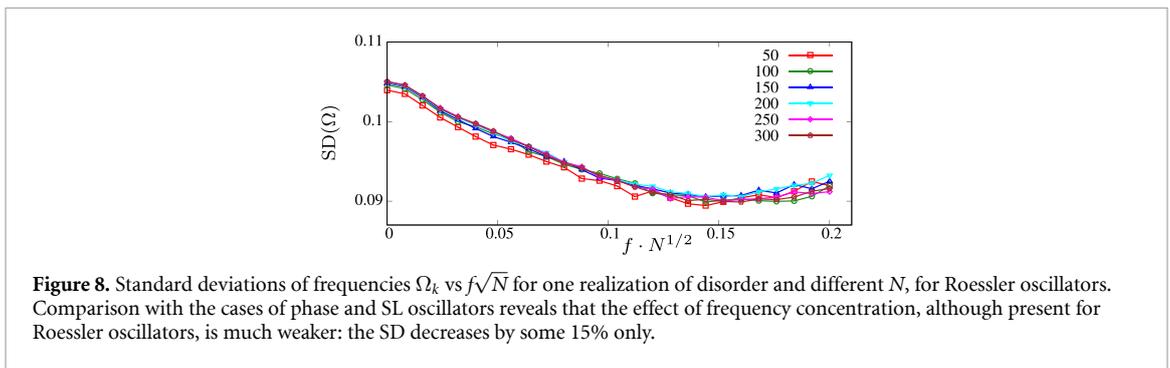


Figure 8. Standard deviations of frequencies Ω_k vs $f\sqrt{N}$ for one realization of disorder and different N , for Roessler oscillators. Comparison with the cases of phase and SL oscillators reveals that the effect of frequency concentration, although present for Roessler oscillators, is much weaker: the SD decreases by some 15% only.

because a trajectory escaped to infinity. Thus, we report on the regimes for relatively small values of coupling strength f only.

The evolution of the standard deviations of the observed frequencies with f is reported in figure 8. One can see that in the explored range of system sizes, this standard deviation is a universal function of rescaled coupling $f\sqrt{N}$. Frequency entrainment is maximal at $f\sqrt{N} \approx 0.14$; however, the overall decrease of the standard deviation is less pronounced than for the phase oscillators and the SLOs (cf figures 2(b) and 6).

Calculations of the COP, depicted in figure 9, revealed the same scaling in dependence on the f, N as in relation (14). The main difference between the corresponding results for the phase oscillators and the SLOs is that the imaginary part D also deviates from zero. We attribute this to the fact that the RO equation (7), due to the existence of the variable z , is non-invariant toward rotations of the phase $a \rightarrow ae^{i\theta}$.

The difference in the amplitude of the values of Q and $Q^{(a)}$ has the same explanation as in the case of the SLOs above: disordered coupling contributes to the instability of oscillations of the variable a , thus increasing the amplitude of the chaotic attractor, which contributes to the growth of $Q^{(a)}$.

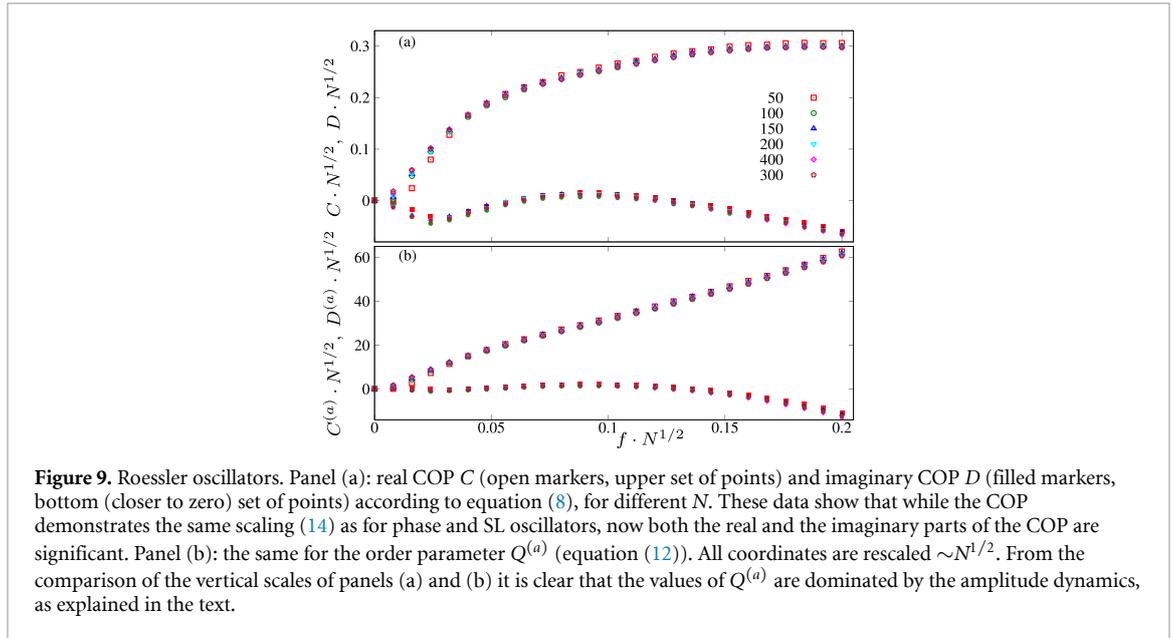


Figure 9. Roessler oscillators. Panel (a): real COP C (open markers, upper set of points) and imaginary COP D (filled markers, bottom (closer to zero) set of points) according to equation (8), for different N . These data show that while the COP demonstrates the same scaling (14) as for phase and SL oscillators, now both the real and the imaginary parts of the COP are significant. Panel (b): the same for the order parameter $Q^{(a)}$ (equation (12)). All coordinates are rescaled $\sim N^{1/2}$. From the comparison of the vertical scales of panels (a) and (b) it is clear that the values of $Q^{(a)}$ are dominated by the amplitude dynamics, as explained in the text.

6. Conclusion

We first summarize the results of our study. We have demonstrated that globally coupled oscillators with a maximally frustrating disorder in the coupling phase shifts show a transition to a partially ordered state. This state cannot be identified by the standard Kuramoto order parameter, which remains close to zero. We introduced a correlation order parameter, the calculation of which requires knowledge of the phase shifts, and demonstrated that it continuously grows with coupling achieving a saturation level $\sim N^{-1/2}$ beyond a characteristic coupling strength $\sim N^{-1/2}$ (expression (14)).

The transition to order can also be characterized via frequency entrainment; here, one does not need to know the phase shifts, and this characterization can be easily implemented in experimental setups. Coupling leads to a concentration of frequencies, which can be straightforwardly quantified by a reduced standard deviation of their distribution. Remarkably, there is an optimal coupling strength at which the concentration is maximal; beyond this strength, the frequencies start to spread again. The disorder-induced spread of the frequencies in a population of identical oscillators explains the latter effect.

We have demonstrated that the ordering behavior described above holds for oscillators with a distribution of natural frequencies and noisy identical oscillators. Furthermore, we showed the effect not only for the phase oscillators but also for the SLOs and for the coupled chaotic Roessler systems. We focused on the case where the phase shifts are not symmetric (i.e. $\alpha_{jk} \neq \alpha_{kj}$). Selected simulations with symmetric random phase shifts (i.e. $\alpha_{jk} = \alpha_{kj}$) demonstrated only minimal quantitative difference to the presented features.

We have explored populations of oscillators with sizes up to several hundred; the scaling properties are well-established in this range. However, probably not all features can be downscaled to small ensembles. The reason is that for small populations, regular quasiperiodic (and periodic for identical deterministic units) become dominant, as the calculations of the largest Lyapunov exponent show. However, for $N \gtrsim 100$, chaotic regimes dominate.

Next, we compare our findings with other results on the disordered populations of coupled oscillators from the recent literature. We have found only paper [21] where a Kuramoto-type model with random phase shifts has been explored. Adopting the replica method for noisy phase oscillators with a Gaussian distribution of natural frequencies, the authors predicted for a maximally frustrating disordered case (phase shifts uniformly distributed in $[0, 2\pi)$) the appearance of a glass synchronization state if $fN^{1/2}$ exceeds a certain level, which nontrivially depends on the noise intensity. It is unclear if this ‘glass synchronization’ state corresponds to a saturated ordered regime observed in our study at large couplings (no numerics is presented in [21]); such a comparison deserves further investigation.

In this respect, we also mention that we have not found any indication of the extremely slow dynamics, typically associated with a ‘glassy’ state. Such a slowing has been observed in the Kuramoto model with another type of disorder, not in the phase shifts, but in the coupling constants (which are sampled from a Gaussian distribution with zero mean), see [11, 13–20]. In the latter model also a sharp ‘volcano transition’

has been observed; we have not found this transition in the random phase shifts model. This comparison shows that different types of disorder apparently lead to different types of dynamics. Thus, it is appealing to look at combinational models, where both types of the disorder (in the phase shifts and the coupling constants) are present; this work is in progress.

Our approach in this paper was purely numerical. Definitely, theoretical progress in understanding the described effects would be highly desirable. Still, there are several numerical setups that we did not explore—in particular, for the phase oscillators, we assumed the coupling in the simplest Kuramoto form, while generally, one has a many-harmonic Daido–Kuramoto coupling. We concentrated on the popular ‘standard’ models of coupled oscillators. One of the possible extensions could be exploring other systems with a disorder in the coupling, which in the ordered case demonstrate a transition to synchrony, such as active rotators [33–35], phase oscillators ‘with inertia’ described by the second-order equations [36–38], higher-dimensional generalizations of the Kuramoto model [20, 39–42]. While disorder in the phase shifts is quite natural for phase oscillators (and other oscillators with a well-defined phase), a corresponding formulation for pulse-coupled units like neurons [43–45] is another challenging task for future studies.

Data availability statement

All data that support the findings of this study are included within the article (and any supplementary files).

Acknowledgments

We thank L Smirnov, D Pazo, I Kiss, O Omelchenko, and M Rosenblum for valuable discussions. A P thanks the University of Florence for warm hospitality.

ORCID iDs

Arkady Pikovsky  <https://orcid.org/0000-0001-9682-7122>

Franco Bagnoli  <https://orcid.org/0000-0002-6293-0305>

References

- [1] Winfree A T 1967 Biological rhythms and the behavior of populations of coupled oscillators *J. Theor. Biol.* **16** 15
- [2] Kuramoto Y 1975 Self-entrainment of a population of coupled nonlinear oscillators *Int. Symp. on Mathematical Problems in Theoretical Physics (Springer Lecture Notes Physics vol 39)* ed H Araki (Springer) p 420
- [3] Hakim V and Rappel W-J 1992 Dynamics of the globally coupled complex Ginzburg–Landau equation *Phys. Rev. A* **46** R7347
- [4] Nakagawa N and Kuramoto Y 1994 From collective oscillations to collective chaos in a globally coupled oscillator system *Physica D* **75** 74
- [5] Pikovsky A, Rosenblum M and Kurths J 1996 Synchronization in a population of globally coupled chaotic oscillators *Europhys. Lett.* **34** 165
- [6] Kalloniatis A C 2010 From incoherence to synchronicity in the network Kuramoto model *Phys. Rev. E* **82** 066202
- [7] Chiba H, Medvedev G S and Mizuhara M S 2018 Bifurcations in the Kuramoto model on graphs *Chaos* **28** 073109
- [8] Juhász R, Kelling J and Ódor G 2019 Critical dynamics of the Kuramoto model on sparse random networks *J. Stat. Mech.* **2019** 053403
- [9] Smirnov L A and Pikovsky A 2023 Dynamics of oscillator populations globally coupled with distributed phase shifts (arXiv:2307.12563 [nlin.AO])
- [10] Zhang Y, Ocampo-Espindola J L, Kiss I Z and Motter A E 2021 Random heterogeneity outperforms design in network synchronization *Proc. Natl Acad. Sci.* **118** e2024299118
- [11] Daido H 1992 Quasientrainment and slow relaxation in a population of oscillators with random and frustrated interactions *Phys. Rev. Lett.* **68** 1073
- [12] Fischer K H and Hertz J A 1993 *Spin Glasses* (Cambridge University Press)
- [13] Stiller J C and Radons G 1998 Dynamics of nonlinear oscillators with random interactions *Phys. Rev. E* **58** 1789
- [14] Daido H 2000 Algebraic relaxation of an order parameter in randomly coupled limit-cycle oscillators *Phys. Rev. E* **61** 2145
- [15] Stiller J C and Radons G 2000 Self-averaging of an order parameter in randomly coupled limit-cycle oscillators *Phys. Rev. E* **61** 2148
- [16] Daido H 2018 Superslow relaxation in identical phase oscillators with random and frustrated interactions *Chaos* **28** 045102
- [17] Ottino-Löffler B and Strogatz S H 2018 Volcano transition in a solvable model of frustrated oscillators *Phys. Rev. Lett.* **120** 264102
- [18] Pazó D and Gallego R 2023 Volcano transition in populations of phase oscillators with random nonreciprocal interactions *Phys. Rev. E* **108** 014202
- [19] Prüser A, Rosmej S and Engel A 2023 Nature of the volcano transition in the fully disordered Kuramoto model (arXiv:2310.09079 [cond-mat.stat-mech])
- [20] Lee S, Jeong Y, Son S-W and Krischer K 2024 Volcano transition in a system of generalized Kuramoto oscillators with random frustrated interactions *J. Phys. A: Math. Theor.* **57** 085702
- [21] Park K, Rhee S W and Choi M Y 1998 Glass synchronization in the network of oscillators with random phase shifts *Phys. Rev. E* **57** 5030
- [22] Sakaguchi H and Kuramoto Y 1986 A soluble active rotator model showing phase transition via mutual entrainment *Prog. Theor. Phys.* **76** 576
- [23] Izhikevich E M 1998 Phase models with explicit time delays *Phys. Rev. E* **58** 905

- [24] Lee W S, Ott E and Antonsen T M 2009 Large coupled oscillator systems with heterogeneous interaction delays *Phys. Rev. Lett.* **103** 044101
- [25] Peter F and Pikovsky A 2018 Transition to collective oscillations in finite Kuramoto ensembles *Phys. Rev. E* **97** 032310
- [26] Rössler O E 1976 An equation for continuous chaos *Phys. Lett. A* **57** 397
- [27] Rosa E, Ott E and Hess M H 1998 Transition to phase synchronization of chaos *Phys. Rev. Lett.* **80** 1642
- [28] Pimenova A V, Goldobin D S, Rosenblum M and Pikovsky A 2016 Interplay of coupling and common noise at the transition to synchrony in oscillator populations *Sci. Rep.* **6** 38518
- [29] Goldobin D S, Pimenova A V, Rosenblum M and Pikovsky A 2017 Competing influence of common noise and desynchronizing coupling on synchronization in the Kuramoto-Sakaguchi ensemble *Eur. Phys. J. Spec. Top.* **226** 1921
- [30] Gómez-Gardenes J, Moreno Y and Arenas A 2007 Paths to synchronization on complex networks *Phys. Rev. Lett.* **98** 034101
- [31] Schröder M, Timme M and Witthaut D 2017 A universal order parameter for synchrony in networks of limit cycle oscillators *Chaos* **27** 073119
- [32] Tao T and Vu V 2010 Random matrices: universality of ESDs and the circular law *Ann. Probab.* **38** 2023
- [33] Sakaguchi H, Shinomoto S and Kuramoto Y 1988 Phase transitions and their bifurcation analysis in a large population of active rotators with mean-field coupling *Prog. Theor. Phys.* **79** 600
- [34] Sonnenschein B, Zaks M, Neiman A and Schimansky-Geier L 2013 Excitable elements controlled by noise and network structure *Eur. Phys. J. Spec. Top.* **222** 2517
- [35] Klinshov V V, Kirillov S Y, Nekorkin V I and Wolfrum M 2021 Noise-induced dynamical regimes in a system of globally coupled excitable units *Chaos* **31** 083103
- [36] Tanaka H, Lichtenberg A and Oishi S 1997 First order phase transition resulting from finite inertia in coupled oscillator systems *Phys. Rev. Lett.* **78** 2104
- [37] Gao J and Efstathiou K 2018 Self-consistent method and steady states of second-order oscillators *Phys. Rev. E* **98** 042201
- [38] Munyayev V O, Smirnov L A, Kostin V A, Osipov G V and Pikovsky A 2020 Analytical approach to synchronous states of globally coupled noisy rotators *New J. Phys.* **22** 023036
- [39] Tanaka T 2014 Solvable model of the collective motion of heterogeneous particles interacting on a sphere *New J. Phys.* **16** 023016
- [40] Chandra S, Girvan M and Ott E 2019 Continuous versus discontinuous transitions in the D-dimensional generalized Kuramoto model: odd D is different *Phys. Rev. X* **9** 011002
- [41] Dai X, Li X, Guo H, Jia D, Perc M, Manshour P, Wang Z and Boccaletti S 2020 Discontinuous transitions and rhythmic states in the D-dimensional Kuramoto model induced by a positive feedback with the global order parameter *Phys. Rev. Lett.* **125** 194101
- [42] Zheng C, Toenjes R and Pikovsky A 2021 Transition to synchrony in a three-dimensional swarming model with helical trajectories *Phys. Rev. E* **104** 014216
- [43] Goel P and Ermentrout B 2002 Synchrony, stability and firing patterns in pulse-coupled oscillators *Physica D* **163** 191
- [44] Montbrió E, Pazó D and Roxin A 2015 Macroscopic description for networks of spiking neurons *Phys. Rev. X* **5** 021028
- [45] Daffertshofer A and Pietras B 2020 Phase synchronization in neural systems *Synergetics* vol 221 (Springer)