

## AN EMPIRICAL CALIBRATION TO ESTIMATE COOL DWARF FUNDAMENTAL PARAMETERS FROM *H*-BAND SPECTRA

ELISABETH R. NEWTON<sup>1</sup>, DAVID CHARBONNEAU<sup>1</sup>, JONATHAN IRWIN<sup>1</sup>, AND ANDREW W. MANN<sup>2</sup>

<sup>1</sup> Harvard-Smithsonian Center for Astrophysics, 60 Garden Street, Cambridge, MA 02138, USA; [enewton@cfa.harvard.edu](mailto:enewton@cfa.harvard.edu)

<sup>2</sup> Astronomy Department, University of Texas at Austin, Austin, TX 78712, USA

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### ABSTRACT

Interferometric radius measurements provide a direct probe of the fundamental parameters of M dwarfs. However, interferometry is within reach for only a limited sample of nearby, bright stars. We use interferometrically measured radii, bolometric luminosities, and effective temperatures to develop new empirical calibrations based on low-resolution, near-infrared spectra. We find that *H*-band Mg and Al spectral features are good tracers of stellar properties, and derive functions that relate effective temperature, radius, and log luminosity to these features. The standard deviations in the residuals of our best fits are, respectively, 73 K,  $0.027 R_{\odot}$ , and 0.049 dex (an 11% error on luminosity). Our calibrations are valid from mid K to mid M dwarf stars, roughly corresponding to temperatures between 3100 and 4800 K. We apply our *H*-band relationships to M dwarfs targeted by the M<sub>Earth</sub> transiting planet survey and to the cool Kepler Objects of Interest (KOIs). We present spectral measurements and estimated stellar parameters for these stars. Parallaxes are also available for many of the M<sub>Earth</sub> targets, allowing us to independently validate our calibrations by demonstrating a clear relationship between our inferred parameters and the stars' absolute *K* magnitudes. We identify objects with magnitudes that are too bright for their inferred luminosities as candidate multiple systems. We also use our estimated luminosities to address the applicability of near-infrared metallicity calibrations to mid and late M dwarfs. The temperatures we infer for the KOIs agree remarkably well with those from the literature; however, our stellar radii are systematically larger than those presented in previous works that derive radii from model isochrones. This results in a mean planet radius that is 15% larger than one would infer using the stellar properties from recent catalogs. Our results confirm the derived parameters from previous in-depth studies of KOIs 961 (*Kepler*-42), 254 (*Kepler*-45), and 571 (*Kepler*-186), the latter of which hosts a rocky planet orbiting in its star's habitable zone.

**Key words:** stars: atmospheres – binaries: general – stars: fundamental parameters – stars: low-mass

**Supporting material:** machine-readable tables

### 1. INTRODUCTION

The characterization of planets is often limited by our knowledge of the stars they orbit. Stellar characterization is not yet as reliable for cooler dwarfs as it is for F, G, and K dwarfs, for which fundamental stellar parameters can be determined with reasonable precision and accuracy (e.g., Valenti & Fischer 2005). The interpretation of observables for cooler dwarfs is complicated by uncertain sources of opacity, the appearance of complex molecules in their atmospheres, and by discrepancies between their observed and theoretical properties. Thus, empirical calibrations remain an important component for our understanding of M dwarfs.

Empirically derived relations between basic stellar parameters provide the basis for determining the physical properties of field M dwarfs. For stars with parallaxes, masses can be estimated using the mass to absolute *K*-band magnitude relation of Delfosse et al. (2000), and radii can then be calculated using a mass–radius relation (e.g., Bayless & Orosz 2006; Boyajian et al. 2012). The mass–magnitude and mass–radius relations are determined from double-lined eclipsing binaries and have precisions of about 10%. Johnson et al. (2012) circumvented the lack of a distance for Kepler Object of Interest (KOI) 254 by constraining the mass and radius of this star using four separate empirical relationships between photometric quantities, mass, radius, and [Fe/H]. Another technique that can be used is the infrared flux method (IRFM; Blackwell & Shallis 1977; Blackwell et al. 1979, 1980), which was extended to M dwarfs

by Casagrande et al. (2008). The IRFM uses the ratio of infrared to bolometric flux and a grid of stellar models to determine the effective temperature ( $T_{\text{eff}}$ ) and bolometric flux ( $F_{\text{bol}}$ ) of a star.

Planet surveys have driven many authors to investigate methods to estimate the stellar parameters of cool dwarfs without parallaxes, which are predominantly based on fitting observations to grids of stellar models. This has proved particularly fruitful for M dwarfs targeted by the *Kepler* survey, some of which host confirmed or candidate planets, including several that are likely to be Earth-sized and within their star's habitable zone. Dressing & Charbonneau (2013) matched observed colors to Dartmouth stellar isochrones (Dotter et al. 2008; Feiden et al. 2011). Mann et al. (2012, 2013b) matched optical spectra to synthetic BT-Settl and PHOENIX spectra, respectively. Muirhead et al. (2012a, 2014) used moderate resolution *K*-band spectra to determine  $T_{\text{eff}}$  and metallicity for 103 cool KOIs, which they then interpolated onto Dartmouth isochrones to infer the stars' radii.

The methods described above rely on relationships derived from binary stars or on matching observed properties to stellar models. While there is reasonable agreement between predicted and observed masses and radii (or luminosities), the observed radii of M dwarfs at a given  $T_{\text{eff}}$  are larger than expected from the models. This was first noted for binary stars (e.g., Popper 1997; Torres & Ribas 2002; Berger et al. 2006; Lopez Morales 2007; Chabrier et al. 2007), but has been demonstrated in stars with interferometric radii measurements (e.g., Berger et al. 2006; Boyajian et al. 2012). Additionally, if magnetic

activity is responsible for the inflated radii, as discussed in Lopez Morales & Ribas (2005), Ribas (2006), Lopez Morales (2007), and Chabrier et al. (2007), a mass–radius–activity relationship would be required to accurately determine the radii of field dwarfs. On the theoretical side, synthetic spectra—which have been improved significantly during the past two decades—suffer from incomplete line lists for the molecules that blanket the spectra, particularly for TiO, VO, metal hydrides, and water vapor (e.g., Valenti et al. 1998; Allard et al. 2000; Leggett et al. 2000; Bean et al. 2006; Önehag et al. 2012; Rajpurohit et al. 2013).

The use of proxies—stars with directly measured parameters that are similar to the star of interest—allows the parameters of a field dwarf to be inferred with limited reliance on stellar models. Muirhead et al. (2012b) used Barnard’s star to anchor stellar models to direct measurements and infer the properties of a similar M dwarf, KOI 961/*Kepler*-42. Building on the method used by Muirhead et al. (2012b), Ballard et al. (2013) used four stars with interferometric radii as proxies to infer the stellar properties of *Kepler*-61. As discussed by Ballard et al. (2013), the use of proxies fills a particularly crucial niche: while the temperature-sensitive index used by Muirhead et al. (2012a, 2012b, 2014) saturates  $T_{\text{eff}} > 3800$  K, the *Kepler* sample is rich in planets orbiting early M and late K dwarfs. Ballard et al. (2013) did not use models in their derivation of stellar parameters, and instead directly used the interferometric radii. However, their method assumes that the properties of *Kepler*-61 match those of stars with the same spectral type, and M dwarf spectral types represent a coarse binning in stellar properties. It has also been shown that late-type dwarfs may have different spectral types in the optical and infrared (e.g., Rojas-Ayala et al. 2012; Pecaut & Mamajek 2013; Newton et al. 2014).

The idea of proxies naturally extends to the identification of empirical tracers of stellar parameters, and the development of such tracers is the motivation for this work. We develop methods to estimate stellar temperatures, radii, and luminosities that do not require parallaxes or stellar models; that are precise, accurate, and extensible; and that can be applied to the large body of moderate-resolution near-infrared (NIR) spectra that is presently available. This idea was also used by Mann et al. (2013b), who used the sample of stars with interferometric measurements to derive index-based calibrations for  $T_{\text{eff}}$  in the visible, *J*, *H*, and *K* bands (although they used model fitting for the stellar parameters reported therein). Our approach is also based on the spectra of the interferometric sample, but it differs from their work in several ways. We use the NIR, which is not as strongly blanketed as the optical by molecular bands that are often sensitive to metallicity. Instead of spectral indices, we use equivalent widths (EWs), which do not require flux-calibrated spectra and are less sensitive to instrument characterization and atmospheric dispersion. Due to ongoing discussions in the literature on measurements of  $T_{\text{eff}}$ , we opt to directly calibrate relations for  $T_{\text{eff}}$ , radius, and luminosity, rather than inferring one property and using additional relations to determine the others. Finally, we have the benefit of three new interferometric measurements from von Braun et al. (2014), and include Gl 725B, which was excluded from the analysis in Mann et al. (2013b), in our radius calibration.

In this paper, we present purely empirical relations between the EWs of NIR features and a cool dwarf’s  $T_{\text{eff}}$ , radius, and luminosity. These calibrations are based on what we refer to as the interferometric sample: the set of stars with interferometrically measured angular diameters for which bolometric fluxes

**Table 1**  
*H*-band Spectral Features

Feature	Feature Window ( $\mu\text{m}$ )		Blue Continuum ( $\mu\text{m}$ )		Red Continuum ( $\mu\text{m}$ )	
Mg (1.48 $\mu\text{m}$ )	1.4872	1.4892	1.4790	1.4850	1.4900	1.4950
Mg (1.50 $\mu\text{m}$ )	1.5020	1.5060	1.4965	1.5000	1.5070	1.5120
K (1.51 $\mu\text{m}$ )	1.5160	1.5180	1.5105	1.5135	1.5185	1.5215
Mg (1.57 $\mu\text{m}$ )	1.5740	1.5780	1.5640	1.5680	1.5785	1.5815
Si (1.58 $\mu\text{m}$ )	1.5880	1.5925	1.5845	1.5875	1.5925	1.5955
CO (1.61 $\mu\text{m}$ )	1.6190	1.6220	1.6120	1.6150	1.6265	1.6295
CO (1.62 $\mu\text{m}$ )	1.6245	1.6265	1.6120	1.6150	1.6265	1.6295
Al-a (1.67 $\mu\text{m}$ )	1.6715	1.6735	1.6550	1.6650	1.6780	1.6820
Al-b (1.67 $\mu\text{m}$ )	1.6745	1.6775	1.6550	1.6650	1.6780	1.6820
Mg (1.71 $\mu\text{m}$ )	1.7100	1.7125	1.7025	1.7055	1.7130	1.7160

**Note.** All wavelengths are given in vacuum.

and distances have been measured, thereby allowing their effective temperatures and luminosities to be determined. We discuss our observations and measurements in Section 2. In Section 3, we discuss the behavior of *H*-band spectral features and present our new calibrations, including a discussion of systematic errors. We then apply our calibrations to M<sub>Earth</sub> M dwarfs in Section 4 and consider the behavior of inferred stellar parameters with absolute *K* magnitude. We also address the applicability of NIR metallicity calibrations from Mann et al. (2013a, 2014) and Newton et al. (2014) to mid and late M dwarfs. We consider the cool KOIs in Section 5, and compare our stellar and planetary parameters to those derived using other techniques. In Section 6, we summarize our findings.

## 2. OBSERVATIONS AND MEASUREMENTS

Our observations and data reduction were carried out as discussed in Newton et al. (2014), which we summarize briefly here. We used the SpeX instrument on IRTF with the  $0''.3 \times 15''$  slit, yielding spectra from 0.8–2.4  $\mu\text{m}$  with  $R = 2000$ . Four observations were acquired for each object, with two observations at each of the two nod positions. Telluric standards were observed directly before or after observations of the interferometry stars used in this work. Flats and wavelength calibrations were taken throughout the night (using an internal quartz lamp or internal thorium–argon lamp, respectively), at roughly 1 hr intervals or after large slews.

We reduced our data using Spextool (Cushing et al. 2004) and used xtellcor to perform telluric corrections (Vacca et al. 2003). We determined absolute radial velocities using the method described in Newton et al. (2014), using telluric features to provide an absolute wavelength calibration. We then shifted the spectra to rest wavelengths.

To measure EWs we defined a wavelength range for the feature and nearby continuum regions on either side. We oversampled the data and numerically integrated the flux within the feature. We estimated errors on the radial velocities and EWs by creating 50 realizations of our data, taking into account correlated noise for the high signal-to-noise ratio observations of the interferometry and M<sub>Earth</sub> samples. We measured the EWs of 26 spectral features and 13 spectral indices in the NIR, which were primarily the same lines as examined by Newton et al. (2014). Our *H*-band features are included in Table 1. We measure the temperature-sensitive indices used in Mann et al. (2013b) and the H<sub>2</sub>O–K2 index (Rojas-Ayala et al. 2012). Wavelengths are given in vacuum.

**Table 2**  
Updated Parameters for New Interferometry Stars

Object	$F_{\text{bol}}$ ( $\text{erg s}^{-1} \text{ cm}^{-2} \times 10^{-8}$ )	$T_{\text{eff}}$ (K)	$L_{\text{bol}}$ ( $L_{\odot}$ )	$R$ ( $R_{\odot}$ )	$\chi_{\text{red}}^2$ <sup>a</sup>	[Fe/H] <sup>b</sup> (dex)
Gl 176	$1.2544 \pm 0.0099$	$3701 \pm 90$	$0.03352 \pm 0.00027$	$0.4525 \pm 0.0221$	0.36	+0.17
Gl 649	$1.3171 \pm 0.0071$	$3604 \pm 46$	$0.04379 \pm 0.00023$	$0.5387 \pm 0.0157$	1.94	+0.05
Gl 876	$1.8863 \pm 0.0115$	$3176 \pm 20$	$0.01290 \pm 0.00008$	$0.3761 \pm 0.0059$	0.78	+0.35

**Notes.**

<sup>a</sup>  $\chi_{\text{red}}^2$  describes the goodness of fit as described in the text and in Mann et al. (2013b).

<sup>b</sup> [Fe/H] estimated using Equation (16) from Mann et al. (2013a).

### 2.1. Stars with Interferometric Measurements

Our calibration sample comprises 25 stars with interferometrically measured radii. We preferentially use the interferometric measurements from Boyajian et al. (2012) and incorporate the measurements they collected from the literature. Where more than one such literature measurement is available, we use the weighted average. The literature sources for measurements are Segransan et al. (2003), Boyajian et al. (2008, 2012), Demory et al. (2009), and von Braun et al. (2012). We also include newer measurements from von Braun et al. (2014).

We observed 20 M dwarfs with interferometrically measured radii from the sample described above. We supplement this sample with the five remaining spectra—Gl 581 and Gl 892 are available in the IRTF spectral library (Cushing et al. 2005; Rayner et al. 2009)<sup>3</sup> and Gl 876, Gl 649, and Gl 176 were observed by Mann et al. (2013b). We cross-correlated these spectra with the radial velocity (RV) standard we use in the rest of our work to ensure agreement between our wavelength calibrations.

The bolometric fluxes and luminosities ( $F_{\text{bol}}$  and  $L_{\text{bol}}$ ) for the stars from Boyajian et al. (2012) were re-measured using spectra and photometry by Mann et al. (2013b). To measure  $F_{\text{bol}}$ , Boyajian et al. (2012) used multicolor photometry to select the best-fitting template spectra from the Pickles (1998) catalog, which they extrapolated to the NIR using photometry. Mann et al. (2013b) demonstrated that fluxes extrapolated beyond  $1.1 \mu\text{m}$  do not match their observed spectra, so they instead used optical and infrared spectra that they obtained for each star. They used models to cover gaps in the spectra and assume Wein’s and Rayleigh–Jeans’ laws at shorter and longer wavelengths than those covered by their spectra. They then adjusted the overall flux level using a correction factor calculated by comparing photometric fluxes to the fluxes measured from the spectra. Mann et al. (2013b) measured  $T_{\text{eff}}$  systematically higher than Boyajian et al. (2012) by 1%.

We use the updated parameters from Mann et al. (2013b) in this work, but note that the radii measured are insensitive to these changes in temperature. We apply the method from Mann et al. (2013b), described above, to the three new objects from von Braun et al. (2014) that we include in this paper. Updated parameters for Gl 176, Gl 649, and Gl 876 are given in Table 2. We also use the [Fe/H] calibration from Mann et al. (2013a) to estimate the iron abundances of these stars (see Section 4.1 for discussion of NIR metallicity calibrations).

The reduced  $\chi^2$  ( $\chi_{\text{red}}^2$ ) of the photometric corrections applied to each spectrum are also included in the table. Large  $\chi_{\text{red}}^2$  indicate that errors in the photometry or spectrum are underestimated, or that there are systematic offsets between the two.

The  $\chi_{\text{red}}^2$  for the three stars presented here are typical of those reported by Mann et al. (2013b).

Mann et al. (2013b) excluded Gl 725B (GJ 725B in their work) due to its atypically large  $\chi_{\text{red}}^2$  (7.9). This indicates a large disagreement between this star’s photometry and spectrum, which affects its  $F_{\text{bol}}$  in an unknown manner. Therefore, we also exclude Gl 725B from our analysis of  $T_{\text{eff}}$  and  $L_{\text{bol}}$ , both of which are subject to this uncertainty. However, we include Gl 725B in our analysis of radius, because while changing a star’s temperature affects the interferometrically measured radius through the adopted limb darkening model, the effect is small (see Boyajian et al. 2012). Finally, we note that while Gl 725B has an unusually cool temperature given its radius, Gl 876 has similar properties; Gl 876 has  $\chi_{\text{red}}^2 = 0.78$ , which is typical of the values in Mann et al. (2013b) and is not excluded from any part of our analysis.

### 2.2. MEarth M Dwarfs

Newton et al. (2014) obtained spectra of 447 mid to late M dwarfs in the solar neighborhood that are targets of the MEarth transiting planet survey (Berta et al. 2012; Irwin et al. 2014). Half of these M dwarfs had previously published parallaxes, which were compiled in Newton et al. (2014). Dittmann et al. (2014) measured parallaxes for most of these stars from MEarth astrometry; with the parallaxes from Dittmann et al. (2014), we have distances to 388 of these M dwarfs.

We note that LSPM J0035+5241S is improperly matched in Newton et al. (2014), who identified it as LSPM J0035+5241N. It is identified properly in this work as LSPM J0035+5241S.

### 2.3. Cool KOIs

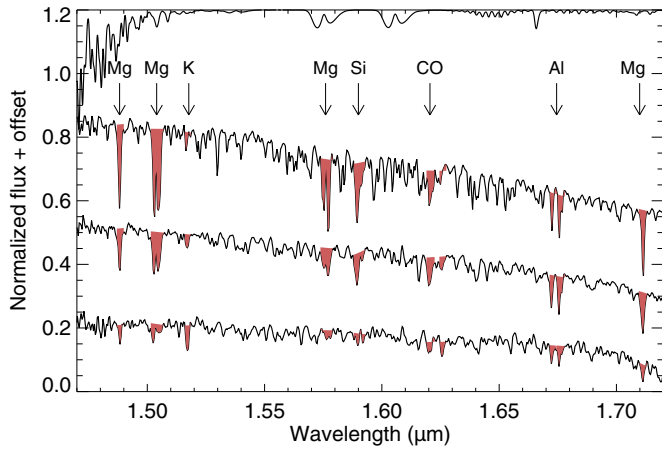
Muirhead et al. (2012a, 2014) presented *H*- and *K*-band spectra for 103 KOIs with  $r - J > 2$ , which implies that they have temperatures below 4000 K. They obtained spectra using the TripleSpec instrument on the Palomar 200 inch (5.1 m) Hale Telescope (Herter et al. 2008), which simultaneously obtains  $R = 2700$  spectra in *J*, *H*, and *K* bands. We convolved the spectra with a Gaussian of fixed width to degrade the resolution to that of IRTF. We then cross-correlated each spectrum with that of our RV standard to place them on the same wavelength calibration as our observations.

Measurements of EWs should not depend on the resolution of the spectrograph. However, most lines in the NIR spectra of a cool dwarf are not free of contaminating features at moderate resolution. We found systematic differences of  $0.1 \text{ \AA}$  between the EWs of some features measured at  $R = 2000$  and  $R = 2700$ , so we stress the importance of accounting for even moderate differences in spectral resolution.

While the subtraction of sky emission lines in IRTF spectra is very robust, it is more difficult for TripleSpec spectra—because

<sup>3</sup> [http://irtfweb.ifa.hawaii.edu/~spex/IRTF\\_Spectral\\_Library/](http://irtfweb.ifa.hawaii.edu/~spex/IRTF_Spectral_Library/)





**Figure 1.** Representative *H*-band spectra of cool dwarfs spanning the temperature range of our calibration. From the top down: a K4V from the IRTF spectral library, an M2V composite spectrum, and an M5V composite spectrum. The composite spectra are from Newton et al. (2014). We indicate the EW measurements with red shading and the elements that dominate the major absorption features. Along the top, we show a typical atmospheric transmission spectrum for Mauna Kea, from Lord (1992).

of the tilt of the slit, the removal of the sky background requires illumination to be very well characterized. Sky emission features persist in some of the KOI spectra and contaminate many spectral features of interest, necessitating the removal of the affected spectra from our analysis. We first identified contaminated spectra by eye, and then developed the following quantitative method to remove spectra based on scatter between the EWs of the components of doublets.

We compared the EWs of the two components of the Mg doublet at  $1.50\ \mu\text{m}$ , of the Mg I doublet at  $1.57\ \mu\text{m}$ , and of the Al I doublet at  $1.67\ \mu\text{m}$ . One component in each of the  $1.50\ \mu\text{m}$  Mg doublet and the Al doublet is contaminated by sky emission. Neither component is contaminated in the  $1.57\ \mu\text{m}$  Mg doublet. We found that in the spectra of the MEarth and interferometry stars, for which sky emission is not prevalent, the EWs of the doublet components are linearly correlated with little scatter. For the  $1.57\ \mu\text{m}$  Mg doublet, which is free of sky emission, there is also little scatter for the KOIs. However, a fraction of the KOIs show inconsistencies in the EWs of the components of the contaminated doublets. The stars that show scatter tend to be those that we determined by eye to contain sky emission.

For each of the two doublets contaminated by sky emission, we used the interferometric sample to determine a linear correlation between the doublet components. Objects deviating from the best-fitting line by more than  $0.75\ \text{\AA}$  for Mg at  $1.50\ \mu\text{m}$  or  $0.5\ \text{\AA}$  for Al at  $1.67\ \mu\text{m}$  were discarded. These limits replicate the contamination we determined by eye and the amount of scatter expected based on observations of the MEarth and interferometry stars. Our final sample of KOIs includes 66 stars. The number of spectra excluded by our cuts does not depend sensitively on the limits we adopt.

### 3. INFERRING STELLAR PARAMETERS

The interferometric sample, which we use to calibrate our empirical relationships, consists of 25 cool stars with radii directly measured using interferometry. We discussed observations of these stars in Section 2.1. Typical *H*-band spectra, with our EW measurements indicated, are shown in Figure 1.

#### 3.1. Behavior of *H*-band Spectral Features as a Function of Effective Temperature

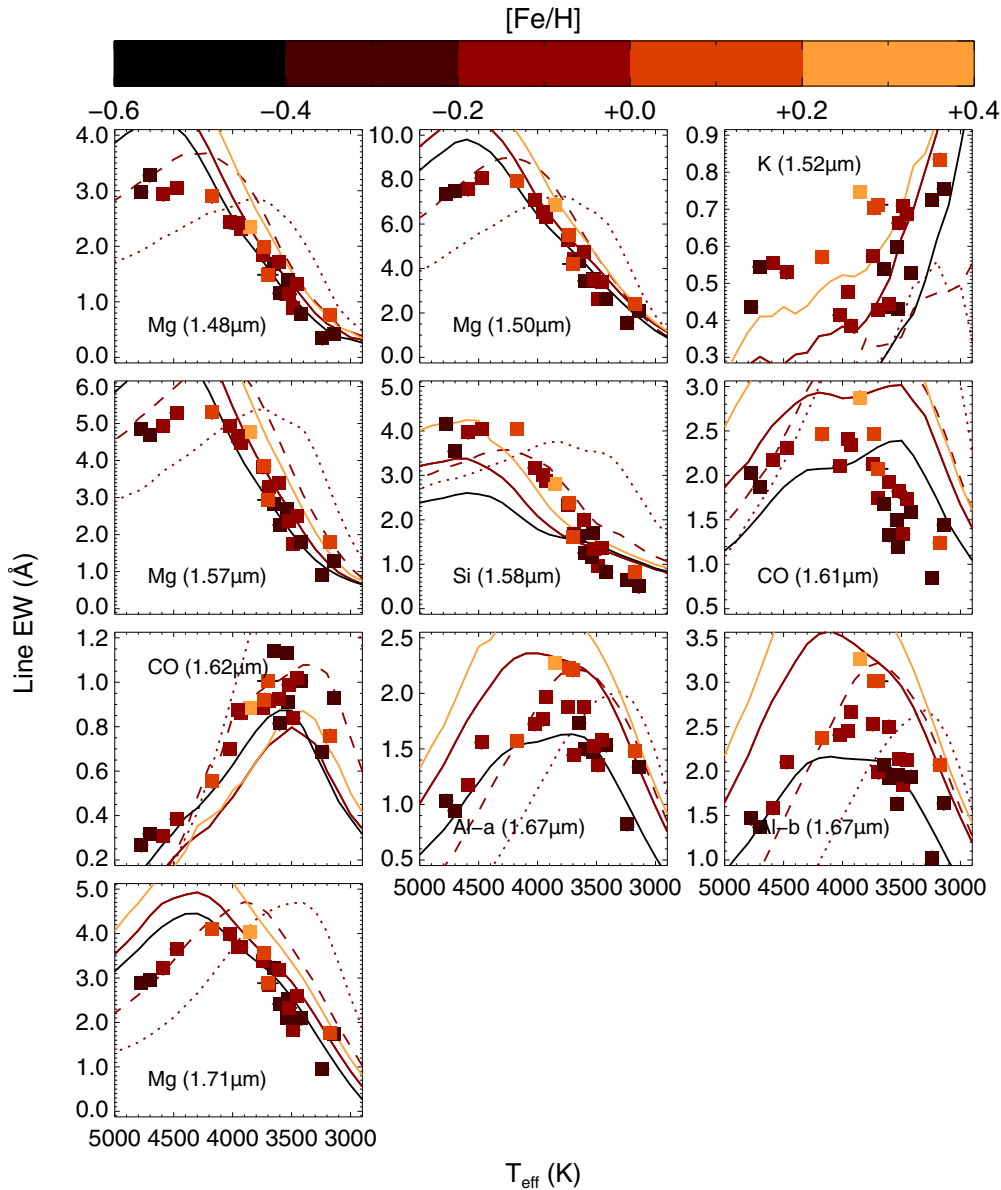
In Figures 2 and 3, we plot selected EWs, EW ratios, and spectral indices for the interferometric calibration sample against the measured  $T_{\text{eff}}$ . Included in these plots are all *H*-band EW measurements, the two EW ratios used in our final analysis (each of which is Mg I / Al I, involving different atomic transitions), and one EW ratio that we found to have a strong temperature dependence for mid M dwarfs. This last EW ratio, K I / Si I, may be useful for studies focusing on later objects. The data points in these figures are colored by their [Fe/H], using the iron abundance calibration from Mann et al. (2013a).

We measured the same spectral features in BT-Settl model spectra (Allard 2014, available online)<sup>4</sup> as we do in our observed spectra. We used models based on the Asplund et al. (2009) solar abundances and degraded their resolution to that of SpeX by convolving the spectra with a fixed-width Gaussian. We then measured EWs numerically using the same technique that we used with our observed spectra. Figures 2 and 3 include the EWs, EW ratios, and spectral features that we measured for a suite of synthetic spectra. We show temperature tracks for three surface gravities, using solar metallicity:  $\log g = 5, 4.5$ , and  $4$ . Late dwarfs have long main sequence lifetimes, so age does not strongly affect  $\log g$ . For  $\log g = 5$ , we also show temperature tracks for two non-solar metallicities, [Fe/H] =  $-0.5$  and [Fe/H] =  $+0.3$ , which roughly span the metallicity range of the interferometric sample. The synthetic spectra with [Fe/H] =  $-0.5$  have an alpha enhancement of  $+0.2$ ; the remainder have no alpha enhancement. The EW measurements of M dwarfs should most closely match the theoretical  $\log g = 5$  temperature track. The surface gravities of the late K dwarfs ( $T_{\text{eff}} > 4000$ ) decrease with increasing temperature, reaching  $\log g = 4.5$  for the hottest star in our sample. Observed EWs for the late K dwarfs are therefore expected to most closely match the temperature track with that surface gravity.

Spectral features that show the smallest metallicity dependence show the most agreement between the EWs measured in the synthetic spectra and in our observed spectra. The Mg I EWs show the best agreement (although the slight metallicity dependence in the models is not evident in our data). For M dwarfs, the EWs closely follow the temperature tracks for stars with  $\log g = 5$ , as expected. For the four K dwarfs, the EWs drop due to the lower surface gravities of these objects. For the Mg I features at  $1.50\ \mu\text{m}$  and  $1.71\ \mu\text{m}$ , the EWs of K dwarfs closely follow the temperature tracks for  $\log g = 4.5$ ; however for those at  $1.48\ \mu\text{m}$  and  $1.57\ \mu\text{m}$ , the EWs are smaller than expected for stars with their temperatures and surface gravities. The EWs of the Si I feature at  $1.58\ \mu\text{m}$  and the CO feature at  $1.62\ \mu\text{m}$  both show tight trends with  $T_{\text{eff}}$ , albeit ones that deviate from those expected from synthetic spectra.

The EWs we observe deviate most strongly from the EWs we calculate from synthetic spectra for features where the models indicate a strong metallicity dependence in the EWs: K I at  $1.52\ \mu\text{m}$ , Si I at  $1.58\ \mu\text{m}$ , CO at  $1.61\ \mu\text{m}$ , and the components of the Al I doublet at  $1.67\ \mu\text{m}$ . The EWs of our observed spectra match neither the strength of the features in the synthetic spectra, nor the amplitude of their metallicity dependence. This result is similar to what was found by Rojas-Ayala et al. (2012), who noted disagreement between EWs measured from synthetic and observed spectra for two metallicity-sensitive *K*-band features: Na I at  $2.21\ \mu\text{m}$  and Ca I at  $2.26\ \mu\text{m}$ . This suggests that there

<sup>4</sup> <http://phoenix.ens-lyon.fr/Grids/BT-Settl/>



**Figure 2.** EWs of  $H$ -band spectral features plotted against measured  $T_{\text{eff}}$  for stars in our interferometric sample. We use the updated temperatures from Mann et al. (2013b) and adopt iron abundances using the Mann et al. (2013a) calibration. The EWs measured from BT-Settl model spectra (Allard 2014) are also plotted. Solid lines represent  $\log g = 5$ , dashed lines  $\log g = 4.5$ , and dotted lines  $\log g = 4$ , all with solar metallicities. For  $\log g = 5$ , we show three metallicities:  $-0.5$  dex (black),  $+0.0$  dex (red), and  $+0.3$  dex (orange).

may be an issue with the treatment of metal abundances in M dwarf atmosphere models.

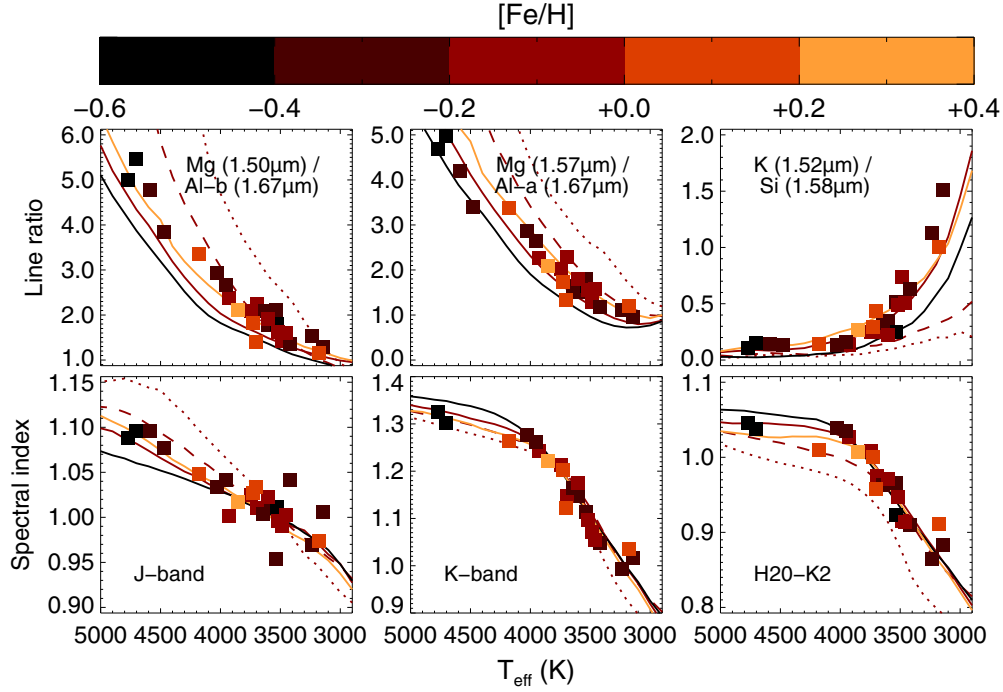
Last, we note that the  $K$ -band spectral indices we measure in our observed spectra are broadly in agreement with those in the synthetic spectra.

### 3.2. Empirical Calibrations for Stellar Parameters Using Spectral Features

We first conducted a principal component analysis (PCA), looking for the strongest correlations between the EW measurements of our calibration sample. We then fit for  $T_{\text{eff}}$ , radius, and luminosity using a linear combination of between one and five principal components. These fits were not better than the simple functions we tried later, and we required line ratios to best fit the full temperature range. The PCA was also hampered by the limited metallicity range of the interferometry stars, but may be a more useful tool in the future.

We investigate simple parameterizations of two or three EWs or EW ratios and simple functions of one EW. The multi-line functions we test (using  $x$ ,  $y$ , and  $z$  to represent an EW measurement or ratio of measurements) are  $ax+by$ ,  $ax+by+cz$ , and  $ax+bx^2+cy$ . The single-line functions we test are  $ax+b/x$ ,  $ax+b\sqrt{x}$ , and  $ax+bx^2$ . We restrict our fits to a single NIR band, and perform a comprehensive search across the possible combinations of features. We use least-squares regression to determine the best-fitting parameters for each combination and the Bayesian information criterion (BIC) to quantitatively compare each fit. After removing fits that qualitatively show systematics in the residuals, we select the fit with the lowest BIC. Our best fits are shown in Figure 4.

We found no viable  $K$ -band relationship and the only viable  $J$ -band relationship has significantly higher scatter than our adopted  $H$ -band relation. Combining features from different regions of the spectrum did not result in significant improvement either. We excluded the Si I, K I, and CO ( $1.62\mu\text{m}$ ) features



**Figure 3.** Same as Figure 2 but plotting select EW ratios and spectral indices against  $T_{\text{eff}}$  using the updated temperatures from Mann et al. (2013b).

from our fits because the behavior of these features as observed in our larger samples of MEarth and *Kepler* M dwarfs was not well represented by the behavior of these features in the interferometric sample, which may be due to metallicity. While empirically these lines can be used to fit stellar parameters as well as those we do use, such fits would have more limited applicability.

In the following equations, we use the element responsible for the spectral feature to indicate an EW measurement, for example, Al-a(1.67  $\mu\text{m}$ ) represents the EW measurement of the blue component of the Al doublet at 1.67  $\mu\text{m}$ . All EWs are in Ångströms. After each equation, we indicate the standard deviation of the residuals ( $\sigma$ ) and the mean absolute deviation (MAD). Our best-fitting relationships are

$$\begin{aligned}
 T_{\text{eff}}/\text{K} &= +271.4 \times \text{Al-a}(1.67 \mu\text{m}) \\
 &+ 392.7 \times \text{Mg}(1.50 \mu\text{m})/\text{Al-b}(1.67 \mu\text{m}) \\
 &+ 2427, \\
 \sigma_{T_{\text{eff}}/\text{K}} &= 73, \\
 \text{MAD}/\text{K} &= 55,
 \end{aligned} \tag{1}$$

$$\begin{aligned}
 R/R_{\odot} &= -0.0489 \times \text{Mg}(1.57 \mu\text{m}) \\
 &+ 0.275 \times \text{Al-a}(1.67 \mu\text{m}) \\
 &+ 0.201 \times \text{Mg}(1.57 \mu\text{m})/\text{Al-a}(1.67 \mu\text{m}) \\
 &- 0.216, \\
 \sigma_{R/R_{\odot}} &= 0.027, \\
 \text{MAD}/R_{\odot} &= 0.020,
 \end{aligned} \tag{2}$$

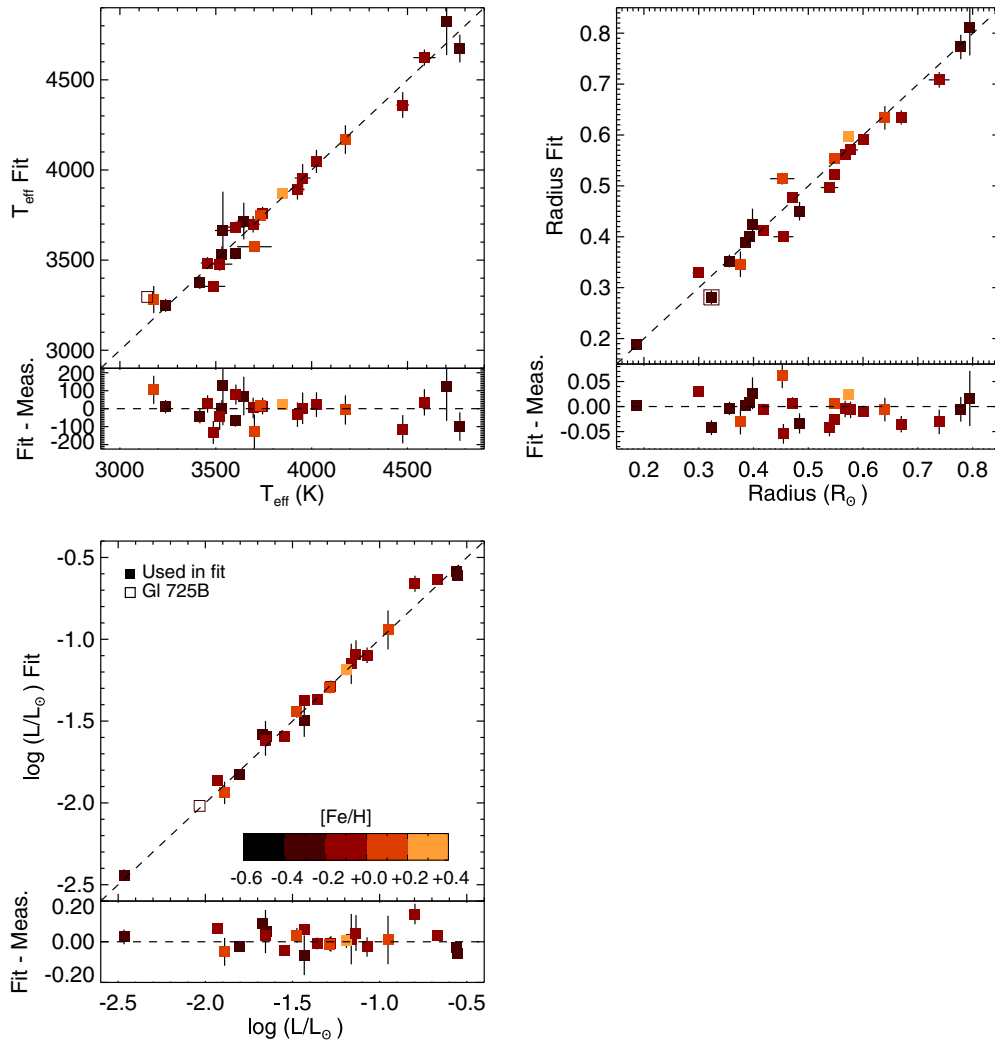
$$\begin{aligned}
 \log L/L_{\odot} &= +0.832 \times \text{Mg}(1.71 \mu\text{m}) \\
 &- 0.176 \times [\text{Mg}(1.71 \mu\text{m})]^2 \\
 &+ 0.266 \times \text{Mg}(1.50 \mu\text{m}) \\
 &- 3.491, \\
 \sigma_{\log L/L_{\odot}} &= 0.049, \\
 \text{MAD}/\text{dex} &= 0.039.
 \end{aligned} \tag{3}$$

We show our best fits in Figure 4. We present the stellar parameters we estimate for the interferometric calibration sample in Table 3; our EW measurements are included in Table 4. These fits are calibrated from our sample of 24 calibrators (25 for our radius fit, for which Gl 725B is included, see Section 2.1). They are valid for stars with  $3200 < T_{\text{eff}} < 4800$  K,  $0.18 < R < 0.8 R_{\odot}$ , and  $-2.5 < \log L/L_{\odot} < -0.5$ .

We note that these calibrations depend on some of the same spectral features. More fundamentally, spectral features are determined by a star's  $T_{\text{eff}}$ ,  $\log g$ , and  $[\text{Fe}/\text{H}]$ —not by its radius or  $L_{\text{bol}}$ . It could be argued that it is only appropriate to derive  $T_{\text{eff}}$ ,  $\log g$ , and  $[\text{Fe}/\text{H}]$  from spectral features, which could then be used to determine radius and  $L_{\text{bol}}$ . We opt to present relations calibrated directly to radius and  $L_{\text{bol}}$  for several reasons: we do not have masses for these stars or other constraints on  $\log g$ , efforts are currently underway that will revise the temperatures and luminosities of our calibration sample, and disagreement continues over  $T_{\text{eff}}$  scales (e.g., Casagrande et al. 2014).

### 3.3. Systematic Uncertainties in the Calibrations

To assess systematic uncertainties in our calibrations, we performed two bootstrap analyses with 1000 samples each. In each case, we randomly created a new realization of our calibration sample and re-fit our relations using the new realization. In the first bootstrap analysis, we randomly drew stars from our calibration sample, allowing calibrators to be sampled repeatedly. Despite the limited number of calibrators, our fits are well constrained at the limits of our calibration. In the second bootstrap analysis, we randomly permuted the residuals from our best fit, assigning each residual to a new data point. This probes the effect of deviations from our best fit that are not captured in the errors on the measured stellar parameters. We find no systematic deviation in  $T_{\text{eff}}$ . For radius and  $L_{\text{bol}}$ , small systematic errors are evident in the difference between the stellar parameters we infer from our bootstrapped fits and those we infer from our fiducial fit. For radius, the difference is  $0.006 R_{\odot}$ . For luminosity, the difference is  $-0.014$  dex for the coolest stars



**Figure 4.** Our best-fitting relationships and the residuals for  $T_{\text{eff}}$  (top left), radius (top right), and  $L_{\text{bol}}$  (bottom left). The horizontal axes show the directly measured stellar parameter using the updated values from Mann et al. (2013b). In the top plot of each panel, the vertical axis shows the stellar parameter we infer from our best fits; in the bottom plot, the vertical axis shows the residuals between our best-fitting values and the directly measured values. The  $[\text{Fe}/\text{H}]$  is indicated by the color of each data point. Gl 725B—used only in our radius fit—is indicated by an open square.

( $\log L/L_{\odot} = -2.5$  dex) and reaches  $-0.021$  dex for the hottest ( $\log L/L_{\odot} > -1.3$  dex). We adopt systematic uncertainties of  $0.006 R_{\odot}$  for radius and  $0.02$  dex for  $\log L/L_{\odot}$ .

We additionally demonstrate the robustness of our calibrations for small stars by removing Gl 699 (Barnard’s star), which anchors the radius and luminosity fits, and re-fit our relations. The effect on the parameters we infer for other stars is negligible, and the relations also provide a good fit if we extrapolate them to Gl 699. Extrapolation gives a radius of  $0.21 R_{\odot}$  (the measured value is  $0.19 R_{\odot}$ ) and a  $\log$  luminosity of  $-2.36$  dex (measured:  $-2.47$  dex) for Gl 699.

$T_{\text{eff}}$ , radius, and  $L_{\text{bol}}$  are not independent parameters and our calibrations should produce consistent values. Looking at the difference between  $\log L/L_{\odot}$  and  $2 \log R/R_{\odot} + 4 \log T/5777 \text{ K}$ , the scatter is  $0.1$  dex with a systematic offset of  $-0.06$  dex. This corresponds to a  $7\%$  difference in radius if the entire systematic offset is assigned to the radius calibration, or a  $3.5\%$  difference in  $T_{\text{eff}}$  if it is assigned to the temperature calibration.

#### 4. APPLICATION TO M DWARFS FROM MEarth

We applied our calibrations for  $T_{\text{eff}}$ , radius, and luminosity to the sample of MEarth M dwarfs observed by Newton et al. (2014).

We present measurements of the EWs used in our best fits, the  $K$ -band temperature index from Mann et al. (2013b), and our inferred stellar parameters in Table 5. The quoted errors on our stellar parameters are random errors (propagated from EW errors) added in quadrature to the standard deviation of the residuals in our best fits. The errors do not include our adopted systematic uncertainties, which are  $0.006 R_{\odot}$  for radius and  $0.02$  dex for  $\log L/L_{\odot}$ . We found no evidence of systematic uncertainties for  $T_{\text{eff}}$ .

The latest spectral type represented in our interferometry calibration sample is M4V (Gl 699, Barnard’s star). The MEarth targets predominantly have temperatures and radii at the extreme end of the calibration range, while close to half of sample have EWs that indicate they are cooler than Gl 699. While our calibrations are not valid for these cool stars, they remain well behaved for late M dwarfs and are useful diagnostics of stellar properties. We therefore report estimated stellar properties for stars beyond the limits of our calibration, but caution that these values may only be used to assess the properties of stars relative to one another.

We limit the application of our temperature calibration to objects with small uncertainties in the EWs. When EW uncertainties are a large fraction of the measurement, their

**Table 3**  
Inferred Parameters for the Interferometric Sample

Star	Measured Parameters			Inferred Parameters		
	$T_{\text{eff}}^{\text{a}}$ (K)	Radius <sup>b</sup> ( $R_{\odot}$ )	$L_{\text{bol}}^{\text{a}}$ ( $\log L/L_{\odot}$ )	$T_{\text{eff}}^{\text{c}}$ (K)	Radius <sup>d</sup> ( $R_{\odot}$ )	$L_{\text{bol}}^{\text{e}}$ ( $\log L/L_{\odot}$ )
Gl 725B <sup>f</sup>	3142 ± 29	0.3232 ± 0.0061	−2.0329 ± 0.0052	(3295 ± 81)	0.281 ± 0.030	(−2.02 ± 0.06)
Gl 876 <sup>g</sup>	3176 ± 20	0.3761 ± 0.0059	−1.8893 ± 0.0027	3281 ± 105	0.346 ± 0.037	−1.94 ± 0.08
Gl 699	3238 ± 11	0.1869 ± 0.0012	−2.4660 ± 0.0038	3248 ± 81	0.188 ± 0.029	−2.44 ± 0.06
Gl 725A	3417 ± 17	0.3561 ± 0.0039	−1.8033 ± 0.0052	3375 ± 81	0.352 ± 0.030	−1.82 ± 0.05
Gl 687	3457 ± 35	0.4183 ± 0.0070	−1.6521 ± 0.0086	3483 ± 80	0.413 ± 0.028	−1.62 ± 0.05
Gl 581	3487 ± 62	0.2990 ± 0.0100	−1.9278 ± 0.0077	3354 ± 74	0.329 ± 0.027	−1.86 ± 0.05
Gl 436	3520 ± 66	0.4546 ± 0.0182	−1.5476 ± 0.0110	3477 ± 81	0.400 ± 0.028	−1.59 ± 0.05
Gl 411	3532 ± 17	0.3924 ± 0.0033	−1.6708 ± 0.0061	3532 ± 85	0.401 ± 0.029	−1.58 ± 0.05
Gl 412A	3537 ± 41	0.3982 ± 0.0091	−1.6548 ± 0.0047	3664 ± 227	0.425 ± 0.041	−1.61 ± 0.12
Gl 15A	3602 ± 13	0.3863 ± 0.0021	−1.6467 ± 0.0052	3534 ± 79	0.388 ± 0.028	−1.60 ± 0.05
Gl 649 <sup>g</sup>	3604 ± 46	0.5387 ± 0.0157	−1.3586 ± 0.0023	3683 ± 79	0.497 ± 0.028	−1.37 ± 0.06
Gl 526	3646 ± 34	0.4840 ± 0.0084	−1.4325 ± 0.0060	3716 ± 125	0.450 ± 0.033	−1.50 ± 0.11
Gl 887	3695 ± 35	0.4712 ± 0.0086	−1.4325 ± 0.0088	3698 ± 86	0.478 ± 0.028	−1.37 ± 0.05
Gl 176 <sup>g</sup>	3701 ± 90	0.4525 ± 0.0221	−1.4746 ± 0.0034	3574 ± 78	0.514 ± 0.029	−1.44 ± 0.06
Gl 880	3731 ± 16	0.5477 ± 0.0048	−1.2856 ± 0.0049	3749 ± 76	0.555 ± 0.028	−1.30 ± 0.06
Gl 809	3744 ± 27	0.5472 ± 0.0067	−1.2798 ± 0.0057	3758 ± 82	0.522 ± 0.028	−1.29 ± 0.06
Gl 205	3850 ± 22	0.5735 ± 0.0044	−1.1905 ± 0.0094	3872 ± 75	0.597 ± 0.027	−1.19 ± 0.06
Gl 338B	3926 ± 37	0.5673 ± 0.0137	−1.1627 ± 0.0145	3892 ± 92	0.562 ± 0.028	−1.15 ± 0.13
Gl 338A	3953 ± 41	0.5773 ± 0.0131	−1.1357 ± 0.0164	3955 ± 106	0.571 ± 0.029	−1.09 ± 0.10
Gl 820B	4025 ± 24	0.6010 ± 0.0072	−1.0722 ± 0.0064	4047 ± 97	0.591 ± 0.028	−1.10 ± 0.07
Gl 380	4176 ± 19	0.6398 ± 0.0046	−0.9518 ± 0.0065	4168 ± 107	0.634 ± 0.036	−0.94 ± 0.13
Gl 702B	4475 ± 33	0.6697 ± 0.0089	−0.7972 ± 0.0108	4360 ± 102	0.635 ± 0.030	−0.66 ± 0.07
Gl 570A	4588 ± 58	0.7390 ± 0.0190	−0.6654 ± 0.0064	4623 ± 85	0.708 ± 0.031	−0.63 ± 0.05
Gl 105A	4704 ± 21	0.7949 ± 0.0062	−0.5589 ± 0.0056	4825 ± 202	0.811 ± 0.061	−0.59 ± 0.06
Gl 892	4773 ± 20	0.7784 ± 0.0053	−0.5521 ± 0.0060	4673 ± 106	0.773 ± 0.036	−0.61 ± 0.05

**Notes.**

<sup>a</sup> Calculated from the interferometric radius and bolometric flux as describe in Section 2.1. Values are from Mann et al. (2013b) Table 1 unless otherwise noted.

<sup>b</sup> Measured from interferometry; see Boyajian et al. (2012) Table 6 for references unless otherwise noted.

<sup>c</sup> Inferred from EWs using Equation (1); fit shown in Figure 4 (top left).

<sup>d</sup> Inferred from EWs using Equation (2); fit shown in Figure 4 (top right).

<sup>e</sup> Inferred from EWs using Equation (3); fit shown in Figure 4 (bottom left).

<sup>f</sup> Gl 725B was excluded from our  $T_{\text{eff}}$  and  $L_{\text{bol}}$  calibrations, as discussed in Section 2.1. The  $T_{\text{eff}}$  and  $L_{\text{bol}}$  we infer from Equations (1) and (3) for this star are nevertheless included in this table.

<sup>g</sup> Values for interferometric radii are from von Braun et al. (2014); updated  $T_{\text{eff}}$  and  $L_{\text{bol}}$  from this work.

ratios have asymmetric error distributions, and we find that the temperatures of cool stars with large errors are systematically hotter than those of similar stars with small errors. Therefore, we do not report estimated temperatures for stars for which the contribution from EW uncertainties (random error) to the total error exceeds 100 K. We found no evidence for similar effects in our radius and luminosity calibrations.

#### 4.1. Using Luminosities to Revisit the Metallicities of the MEarth Sample

Newton et al. (2014) estimated the metallicities of the M dwarfs in our sample from the EW of the NaI feature at  $2.2\mu\text{m}$  in the  $K$  band of IRTF spectra (see also Rojas-Ayala et al. 2010, 2012; Terrien et al. 2012). Their empirical relation was based on M dwarfs with NIR spectral types M5V and earlier. Mann et al. (2013a) constructed a similar calibration for K5–M5 dwarfs using a combination of the Na feature and other NIR spectral features. Mann et al. (2014) used early–late M dwarf pairs to bootstrap a calibration valid for M dwarfs with spectral types M4V–M9V, and showed that Newton et al. (2014) and other previous works either over- or underestimated the metallicities of M7–M9 dwarfs. In the following, we make

use of our estimated stellar luminosities to investigate the best metallicity calibration to use on our sample of stars.

Newton et al. (2014) looked for systematic trends in the metallicities of stars with spectral type, which they determined by matching each NIR spectrum to that of a spectral standard. They found that the M5–M7 dwarfs in their sample appeared metal-rich by 0.1 dex relative to earlier stars (Figure 13 in their work). Because M5 dwarfs comprised a quarter of their metallicity calibration sample, they assumed the metallicities of the MEarth M5 dwarfs—and by extension the M6 and M7 dwarfs—were estimated accurately. However, M spectral types are a coarse binning in stellar parameters. In Figure 5, we show that while the M4V metallicity calibrators used by Newton et al. (2014) are typical of the M4 dwarfs in the MEarth sample, 70% of the M5 dwarfs have luminosities lower than the median luminosity of the M5V metallicity calibrators. The calibration from Newton et al. (2014) may not be valid for all M dwarfs assigned an M5V spectral type.

In Figure 6, we compare  $[\text{Fe}/\text{H}]$  values measured for the MEarth M dwarfs using the calibrations from Newton et al. (2014) and Mann et al. (2013a, 2014). We only show M dwarfs whose metallicities, as estimated from the Newton et al. (2014) calibration, are less than 0.25 dex, because the



**Table 4**  
Measured Quantities for the Interferometric Sample

Star	Mg (1.48 $\mu\text{m}$ ) ( $\text{\AA}$ )	Mg (1.50 $\mu\text{m}$ ) ( $\text{\AA}$ )	K (1.51 $\mu\text{m}$ ) ( $\text{\AA}$ )	Mg (1.58 $\mu\text{m}$ ) ( $\text{\AA}$ )	Si (1.58 $\mu\text{m}$ ) ( $\text{\AA}$ )	CO (1.61 $\mu\text{m}$ ) ( $\text{\AA}$ )	CO (1.62 $\mu\text{m}$ ) ( $\text{\AA}$ )	Al-a (1.67 $\mu\text{m}$ ) ( $\text{\AA}$ )	Al-b (1.67 $\mu\text{m}$ ) ( $\text{\AA}$ )	Mg (1.71 $\mu\text{m}$ ) ( $\text{\AA}$ )	<i>J</i> Band <sup>a</sup>	<i>H</i> Band <sup>a</sup>	<i>K</i> Band <sup>a</sup>	H <sub>2</sub> O–K2 <sup>b</sup>
Gl 725B	0.42 $\pm$ 0.07	2.11 $\pm$ 0.08	0.75 $\pm$ 0.06	1.28 $\pm$ 0.11	0.50 $\pm$ 0.08	1.45 $\pm$ 0.08	0.93 $\pm$ 0.06	1.34 $\pm$ 0.06	1.64 $\pm$ 0.08	1.74 $\pm$ 0.10	1.01	0.80	1.02	0.88
Gl 876	0.76 $\pm$ 0.18	2.38 $\pm$ 0.20	0.83 $\pm$ 0.11	1.79 $\pm$ 0.19	0.83 $\pm$ 0.19	1.24 $\pm$ 0.22	0.76 $\pm$ 0.13	1.48 $\pm$ 0.17	2.07 $\pm$ 0.21	1.77 $\pm$ 0.20	0.97	0.86	1.04	0.91
Gl 699	0.34 $\pm$ 0.03	1.57 $\pm$ 0.05	0.73 $\pm$ 0.03	0.91 $\pm$ 0.05	0.64 $\pm$ 0.04	0.85 $\pm$ 0.04	0.69 $\pm$ 0.04	0.82 $\pm$ 0.03	1.03 $\pm$ 0.05	0.95 $\pm$ 0.07	0.97	0.73	0.99	0.86
Gl 725A	0.78 $\pm$ 0.08	2.63 $\pm$ 0.09	0.53 $\pm$ 0.06	1.79 $\pm$ 0.09	0.84 $\pm$ 0.09	1.59 $\pm$ 0.06	1.01 $\pm$ 0.07	1.53 $\pm$ 0.09	1.94 $\pm$ 0.07	2.09 $\pm$ 0.10	1.04	0.85	1.05	0.91
Gl 687	1.32 $\pm$ 0.07	3.39 $\pm$ 0.09	0.69 $\pm$ 0.05	2.49 $\pm$ 0.08	1.36 $\pm$ 0.08	1.74 $\pm$ 0.05	1.02 $\pm$ 0.05	1.58 $\pm$ 0.06	2.13 $\pm$ 0.08	2.59 $\pm$ 0.09	1.00	0.79	1.06	0.91
Gl 581	0.89 $\pm$ 0.02	2.63 $\pm$ 0.03	0.71 $\pm$ 0.02	1.74 $\pm$ 0.04	0.96 $\pm$ 0.03	1.34 $\pm$ 0.03	0.84 $\pm$ 0.02	1.35 $\pm$ 0.03	1.84 $\pm$ 0.03	1.83 $\pm$ 0.05	0.99	0.83	1.07	0.92
Gl 436	1.14 $\pm$ 0.06	3.46 $\pm$ 0.08	0.66 $\pm$ 0.06	2.38 $\pm$ 0.07	1.35 $\pm$ 0.08	1.82 $\pm$ 0.08	0.99 $\pm$ 0.05	1.53 $\pm$ 0.06	2.13 $\pm$ 0.10	2.33 $\pm$ 0.10	1.00	0.87	1.10	0.95
Gl 411	1.40 $\pm$ 0.09	3.52 $\pm$ 0.08	0.43 $\pm$ 0.05	2.43 $\pm$ 0.11	1.70 $\pm$ 0.10	1.20 $\pm$ 0.09	0.91 $\pm$ 0.08	1.47 $\pm$ 0.07	1.96 $\pm$ 0.10	2.53 $\pm$ 0.15	1.01	0.80	1.11	0.92
Gl 412A	1.14 $\pm$ 0.29	3.46 $\pm$ 0.39	0.60 $\pm$ 0.19	2.69 $\pm$ 0.34	1.15 $\pm$ 0.41	1.49 $\pm$ 0.29	1.13 $\pm$ 0.27	1.50 $\pm$ 0.25	1.63 $\pm$ 0.36	2.09 $\pm$ 0.29	0.95	0.84	1.11	0.97
Gl 15A	1.15 $\pm$ 0.04	3.44 $\pm$ 0.07	0.44 $\pm$ 0.05	2.26 $\pm$ 0.06	1.25 $\pm$ 0.06	1.33 $\pm$ 0.06	0.81 $\pm$ 0.05	1.50 $\pm$ 0.05	1.93 $\pm$ 0.06	2.42 $\pm$ 0.05	1.02	0.90	1.15	0.97
Gl 649	1.72 $\pm$ 0.05	4.76 $\pm$ 0.07	0.44 $\pm$ 0.04	3.39 $\pm$ 0.11	2.01 $\pm$ 0.08	1.93 $\pm$ 0.06	0.93 $\pm$ 0.06	1.88 $\pm$ 0.06	2.50 $\pm$ 0.08	3.18 $\pm$ 0.08	1.02	0.87	1.18	0.97
Gl 526	1.67 $\pm$ 0.17	4.32 $\pm$ 0.23	0.54 $\pm$ 0.10	2.84 $\pm$ 0.21	1.63 $\pm$ 0.18	1.68 $\pm$ 0.12	1.14 $\pm$ 0.13	1.73 $\pm$ 0.14	2.07 $\pm$ 0.21	3.23 $\pm$ 0.25	1.00	0.87	1.17	0.96
Gl 887	1.50 $\pm$ 0.07	4.45 $\pm$ 0.07	0.43 $\pm$ 0.05	3.30 $\pm$ 0.09	1.69 $\pm$ 0.08	1.74 $\pm$ 0.07	0.91 $\pm$ 0.06	1.44 $\pm$ 0.06	1.98 $\pm$ 0.09	2.85 $\pm$ 0.08	1.01	0.85	1.15	0.98
Gl 176	1.49 $\pm$ 0.10	4.20 $\pm$ 0.11	0.71 $\pm$ 0.06	2.94 $\pm$ 0.10	1.62 $\pm$ 0.12	2.07 $\pm$ 0.09	1.01 $\pm$ 0.08	2.21 $\pm$ 0.07	3.01 $\pm$ 0.10	2.88 $\pm$ 0.12	1.03	0.91	1.12	0.96
Gl 880	1.99 $\pm$ 0.04	5.52 $\pm$ 0.08	0.70 $\pm$ 0.04	3.85 $\pm$ 0.06	2.38 $\pm$ 0.07	2.47 $\pm$ 0.06	0.92 $\pm$ 0.05	2.22 $\pm$ 0.07	3.02 $\pm$ 0.05	3.55 $\pm$ 0.06	1.03	0.91	1.20	1.00
Gl 809	1.83 $\pm$ 0.05	5.29 $\pm$ 0.09	0.57 $\pm$ 0.05	3.82 $\pm$ 0.09	2.34 $\pm$ 0.06	2.12 $\pm$ 0.06	0.88 $\pm$ 0.05	1.88 $\pm$ 0.07	2.53 $\pm$ 0.09	3.38 $\pm$ 0.08	1.02	0.89	1.21	1.01
Gl 205	2.35 $\pm$ 0.05	6.88 $\pm$ 0.05	0.75 $\pm$ 0.04	4.77 $\pm$ 0.06	2.81 $\pm$ 0.06	2.88 $\pm$ 0.05	0.89 $\pm$ 0.06	2.27 $\pm$ 0.04	3.26 $\pm$ 0.06	4.05 $\pm$ 0.05	1.02	0.89	1.22	1.01
Gl 338B	2.31 $\pm$ 0.11	6.32 $\pm$ 0.14	0.39 $\pm$ 0.08	4.48 $\pm$ 0.14	2.88 $\pm$ 0.10	2.35 $\pm$ 0.13	0.86 $\pm$ 0.09	1.97 $\pm$ 0.13	2.66 $\pm$ 0.12	3.70 $\pm$ 0.25	1.00	0.87	1.24	1.03
Gl 338A	2.42 $\pm$ 0.11	6.55 $\pm$ 0.13	0.48 $\pm$ 0.10	4.65 $\pm$ 0.15	3.01 $\pm$ 0.13	2.41 $\pm$ 0.12	0.87 $\pm$ 0.10	1.77 $\pm$ 0.10	2.45 $\pm$ 0.16	3.71 $\pm$ 0.17	1.04	0.88	1.26	1.03
Gl 820B	2.43 $\pm$ 0.05	7.07 $\pm$ 0.08	0.41 $\pm$ 0.05	4.94 $\pm$ 0.10	3.17 $\pm$ 0.08	2.11 $\pm$ 0.08	0.70 $\pm$ 0.08	1.72 $\pm$ 0.09	2.41 $\pm$ 0.12	3.98 $\pm$ 0.07	1.03	0.90	1.28	1.04
Gl 380	2.90 $\pm$ 0.10	7.95 $\pm$ 0.17	0.57 $\pm$ 0.08	5.30 $\pm$ 0.14	4.05 $\pm$ 0.09	2.46 $\pm$ 0.08	0.55 $\pm$ 0.10	1.57 $\pm$ 0.13	2.37 $\pm$ 0.12	4.11 $\pm$ 0.18	1.05	0.87	1.26	1.01
Gl 702B <sup>c</sup>	3.05 $\pm$ 0.06	8.07 $\pm$ 0.07	0.53 $\pm$ 0.05	5.28 $\pm$ 0.09	4.05 $\pm$ 0.07	2.31 $\pm$ 0.09	0.38 $\pm$ 0.09	1.56 $\pm$ 0.07	2.10 $\pm$ 0.09	3.65 $\pm$ 0.10	1.08	...	...	...
Gl 570A <sup>c</sup>	2.93 $\pm$ 0.03	7.56 $\pm$ 0.03	0.55 $\pm$ 0.02	4.94 $\pm$ 0.04	3.97 $\pm$ 0.03	2.17 $\pm$ 0.04	0.31 $\pm$ 0.04	1.18 $\pm$ 0.03	1.58 $\pm$ 0.04	3.22 $\pm$ 0.03	1.10	...	...	...
Gl 105A	3.28 $\pm$ 0.06	7.48 $\pm$ 0.09	0.54 $\pm$ 0.04	4.69 $\pm$ 0.08	3.54 $\pm$ 0.08	1.87 $\pm$ 0.07	0.32 $\pm$ 0.06	0.94 $\pm$ 0.07	1.37 $\pm$ 0.12	2.95 $\pm$ 0.08	1.10	0.86	1.30	1.04
Gl 892	2.98 $\pm$ 0.03	7.34 $\pm$ 0.05	0.44 $\pm$ 0.02	4.86 $\pm$ 0.04	4.16 $\pm$ 0.04	2.03 $\pm$ 0.04	0.27 $\pm$ 0.04	1.03 $\pm$ 0.04	1.47 $\pm$ 0.06	2.90 $\pm$ 0.06	1.09	0.96	1.32	1.05

**Notes.** EW definitions given in Table 1.

<sup>a</sup> Temperature-sensitive indices from Mann et al. (2013b).

<sup>b</sup> Temperature-sensitive index from Covey et al. (2010) and Rojas-Ayala et al. (2012).

<sup>c</sup> Spectral windows for the *K*-band indices fall outside of calibrated wavelength range.

**Table 5**  
Measured Quantities and Inferred Stellar Parameters for M Dwarfs in the MEarth Sample

LSPM Name	Comp	Mg (1.48 $\mu\text{m}$ ) ( $\text{\AA}$ )	Mg (1.50 $\mu\text{m}$ ) ( $\text{\AA}$ )	Mg (1.58 $\mu\text{m}$ ) ( $\text{\AA}$ )	Al-a (1.67 $\mu\text{m}$ ) ( $\text{\AA}$ )	Al-b (1.67 $\mu\text{m}$ ) ( $\text{\AA}$ )	Mg (1.71 $\mu\text{m}$ ) ( $\text{\AA}$ )	<i>K</i> Band <sup>a</sup>	$T_{\text{eff}}$ <sup>b</sup> (K)	Radius <sup>c</sup> ( $R_{\odot}$ )	$\log L/L_{\odot}$ <sup>d</sup> (dex)	Parallax (mas)	Ref. <sup>e</sup>	[Fe/H] <sub>M13</sub> <sup>f</sup> (dex)	[Fe/H] <sub>M14</sub> <sup>g</sup> (dex)
LSPM J0001+0659		0.80 $\pm$ 0.05	0.59 $\pm$ 0.11	−0.10 $\pm$ 0.16	0.48 $\pm$ 0.08	0.86 $\pm$ 0.14	0.40 $\pm$ 0.13	0.94	2822 $\pm$ 100	−0.122 $\pm$ 0.072	−3.03 $\pm$ 0.11	57.1 $\pm$ 2.6	D14	−0.05	−0.08
LSPM J0007+0800		0.73 $\pm$ 0.10	2.10 $\pm$ 0.12	1.56 $\pm$ 0.12	1.05 $\pm$ 0.05	1.10 $\pm$ 0.09	1.67 $\pm$ 0.12	1.03	3459 $\pm$ 105	0.295 $\pm$ 0.032	−2.04 $\pm$ 0.06	44.0 $\pm$ 6.3	N14	−0.71	−0.72
LSPM J0008+2050		0.19 $\pm$ 0.05	1.36 $\pm$ 0.05	0.66 $\pm$ 0.09	0.84 $\pm$ 0.04	1.09 $\pm$ 0.07	1.03 $\pm$ 0.10	0.95	3147 $\pm$ 82	0.140 $\pm$ 0.032	−2.46 $\pm$ 0.07	67.5 $\pm$ 2.7	D14	+0.02	+0.12
LSPM J0015+4344		0.97 $\pm$ 0.09	2.05 $\pm$ 0.13	1.27 $\pm$ 0.15	1.07 $\pm$ 0.07	1.32 $\pm$ 0.08	1.33 $\pm$ 0.13	0.95	3328 $\pm$ 93	0.255 $\pm$ 0.034	−2.15 $\pm$ 0.08	35.1 $\pm$ 1.5	D14	+0.22	+0.36
LSPM J0015+1333		0.59 $\pm$ 0.05	1.58 $\pm$ 0.07	0.92 $\pm$ 0.11	0.94 $\pm$ 0.06	1.19 $\pm$ 0.07	1.34 $\pm$ 0.09	1.02	3201 $\pm$ 83	0.194 $\pm$ 0.033	−2.27 $\pm$ 0.06	99.8 $\pm$ 4.0	D14	−0.22	−0.24

**Notes.** For binaries that were not resolved in the LSPM catalogs but which we resolve, we use the appropriate LSPM name in the first column and identify the components in the second column.

<sup>a</sup> Temperature-sensitive *K*-band index from Mann et al. (2013b).

<sup>b</sup> Inferred from EWs using Equation (1); fit shown in Figure 4 (top left).

<sup>c</sup> Inferred from EWs using Equation (2); fit shown in Figure 4 (top right).

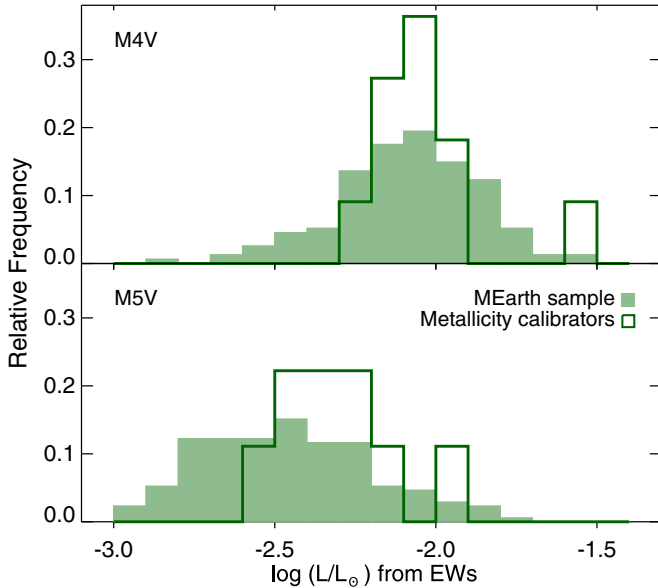
<sup>d</sup> Inferred from EWs using Equation (3); fit shown in Figure 4 (bottom left).

<sup>e</sup> D14: parallax from Dittmann et al. (2014); N14: parallax compiled from the literature by Newton et al. (2014). For parallaxes compiled from the literature, we refer the reader to Newton et al. (2014) for the reference.

<sup>f</sup> [Fe/H] using the Mann et al. (2013a) calibration, appropriate for early M dwarfs ( $\log L/L_{\odot} > -2.25$ ).

<sup>g</sup> [Fe/H] using the Mann et al. (2014) calibration, appropriate for late M dwarfs ( $\log L/L_{\odot} < -2.25$ ).

(This table is available in its entirety in machine-readable form.)

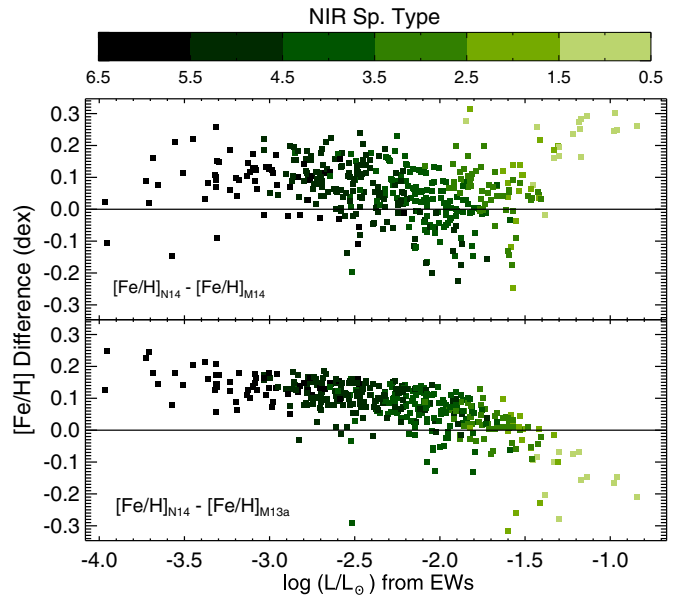


**Figure 5.** Distribution of inferred luminosities for M4 dwarfs (top panel) and for M5 dwarfs (bottom panel). M dwarfs in the MEarth sample are represented by the light green shaded histogram. The M dwarfs in the Newton et al. (2014) metallicity calibration sample are represented by the dark green histogram.

Newton et al. (2014) calibration saturates for more metal-rich stars. This demonstrates that while the Newton et al. (2014) calibration is generally valid for M2–M4V stars, it over- and underestimates the metallicities of later and earlier stars, respectively.

We compare the Newton et al. (2014) calibration to the Mann et al. (2014) calibration, which is applicable to late M dwarfs, in the top panel of Figure 6. This demonstrates that Newton et al. (2014) overestimates the metallicities of M dwarfs with  $\log L/L_\odot < -2.25$ , which includes the majority of stars M5V and later. The two calibrations agree for stars with  $-2.25 < \log L/L_\odot < -1.75$  (approximately M4V), with a median offset of 0.03 dex. For hotter stars, the Mann et al. (2014) calibration is not applicable, and the estimated metallicities deviate significantly. In the bottom panel of Figure 6, we find an almost linear trend with luminosity when we compare the Newton et al. (2014) to the Mann et al. (2013a) calibration, which is applicable to early M dwarfs. The two calibrations agree most closely for  $-2.0 < \log L/L_\odot < -1.5$ , (approximately M2V–M3V). For earlier stars the Newton et al. (2014) calibration underestimates metallicities, which is a finding anticipated in that work.

We update the metallicities of the MEarth M dwarfs by stitching together the Mann et al. (2013a) and Mann et al. (2014) relations. For stars with  $\log L/L_\odot < -1.75$ , we use the Mann et al. (2014) relation, which was calibrated specifically for late-type dwarfs. For stars with  $\log L/L_\odot > -1.75$ , we use the Mann et al. (2013a) relation, for which the calibration was dominated by early M dwarfs. We use the code provided by A. Mann<sup>5</sup> to calculate metallicities. We note that the method we use to measure EWs (oversampling and using the trapezoidal rule to numerically integrate the flux beneath the pseudo-continuum), while more robust than the IDL routines TOTAL or SUM without oversampling, produces different (typically larger) EW measurements.



**Figure 6.** Difference between literature [Fe/H] values vs. stellar luminosity. We use relations from three works: Mann et al. (2013a, M13a), Newton et al. (2014, N14), and Mann et al. (2014, M14). The Mann et al. (2013a) relation is calibrated for M4V stars and earlier, while the Mann et al. (2014) relation is calibrated for stars M4V and later. Points are colored by the NIR spectral type assigned by eye by Newton et al. (2014).

#### 4.2. Comparison to the Mann et al. Temperature and Radius Calibrations

Mann et al. (2013b) developed empirical relations for cool dwarf temperatures that are based on spectral indices in visual and infrared bands. As we have done in this work, they used stars with interferometric measurements to calibrate the relationships. The temperature-sensitive indices are the ratios between the median flux in three wavelength windows (Equation (13) in their paper), and quantify the curvature of the spectrum in each band. Mann et al. (2013b) chose the continuum windows to minimize the scatter in the resulting temperature calibration.

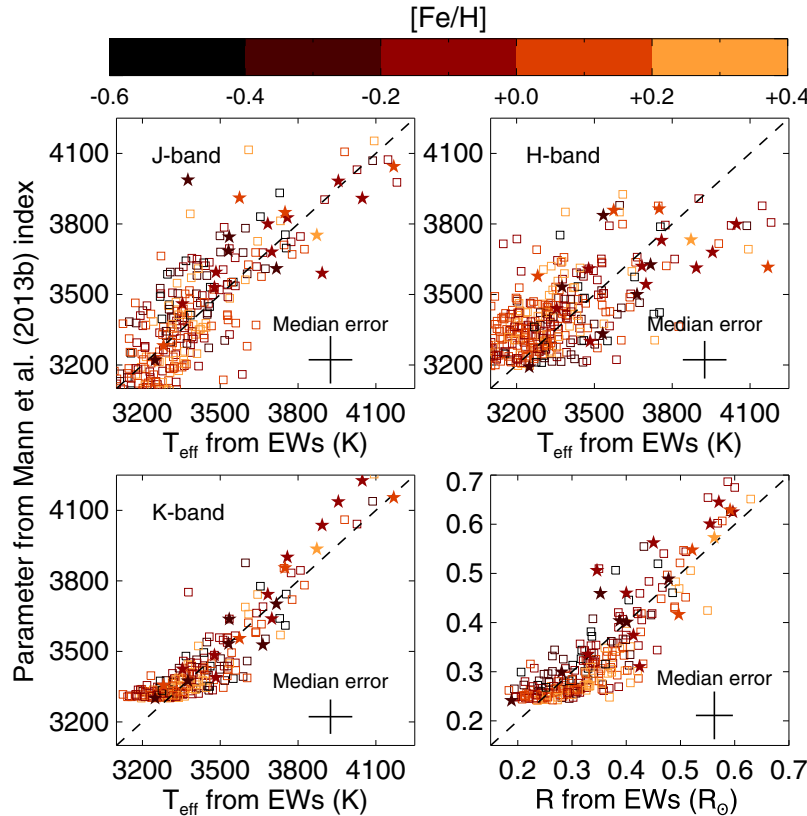
We compare the stellar parameters we estimated for the MEarth M dwarfs to those we calculated using the Mann et al. (2013b) *J*-, *H*-, and *K*-band  $T_{\text{eff}}$  calibrations in Figure 7. We then applied their temperature–radius polynomial (Equation (6) in their work) to convert the *K*-band temperatures to radii. Accurate transformation requires additional significant figures, included here for completeness:

$$\begin{aligned} R/R_\odot = & -16.883175 + 1.1835396 \times 10^{-2} \times (T_{\text{eff}}/\text{K}) \\ & - 2.7087196 \times 10^{-6} \times (T_{\text{eff}}/\text{K})^2 \\ & + 2.1050281 \times 10^{-10} \times (T_{\text{eff}}/\text{K})^3. \end{aligned} \quad (4)$$

We limit our comparison to those stars within the limits of the calibrations in question; only these objects are included in Figure 7. We note that the Mann et al. (2013b) *K*-band calibration saturates for stars with  $T_{\text{eff}} < 3300$  K, so there is a minimum value for the temperatures derived from this relation.

The Mann et al. (2013b) *J*- and *H*-band indices are not good predictors of  $T_{\text{eff}}$ . The differences between the  $T_{\text{eff}}$  inferred from our two calibrations have standard deviations ( $\sigma_{\Delta T}$ ) that are larger than expected from the errors in the calibrations:  $\sigma_{\Delta T} = 140$  K for the *J*-band calibration and 170 K for the *H*-band calibration. This is because the wavelength windows for these indices fall in regions with strong telluric absorption. Temperatures from the *K*-band index and our EW-based calibration

<sup>5</sup> <https://github.com/awmann/metal>



**Figure 7.** Comparison of the calibrations from Mann et al. (2013b) and from this work, showing the  $T_{\text{eff}}$  and radii estimated for the M dwarfs from Newton et al. (2014). On the horizontal axis we show the parameters we infer using the EW-based methods developed in this work. On the vertical axis we show the parameters we infer using the index-based methods from Mann et al. (2013b). We use the polynomial presented in that work to convert temperatures estimated from the  $K$ -band index to radii. The color of the points indicates the stars’ metallicities.

agree within the errors of the calibrations, with  $\sigma_{\Delta T} = 90$  K, although for  $3500 < T_{\text{eff}} < 3800$  the median temperature we infer from EWs is 40 K hotter.

The radii we infer directly from EWs and the radii calculated from the  $K$  band using the temperature–radius relation have  $\sigma_{\Delta R} = 0.05 R_{\odot}$ . However, there is evidence for a systematic offset for stars with  $0.3 < R/R_{\odot} < 0.4$ , where the median offset is  $0.045 R_{\odot}$  (12%), with radii calculated from EWs being larger. We believe this is due to systematics in the Mann et al. (2013b) relations: the  $K$ -band temperature relation predicts temperatures that are too cool by 50–100 K for the three interferometry stars with temperatures around 3500 K (Figure 11 in their work). Because the slope of the temperature–radius relation is steep for stars of this size, temperatures that are 50–100 K too cool result in radii that are 0.03–0.06  $R_{\odot}$  (15%) too small. This is consistent with the differences we find and supports the temperatures and radii we infer from EWs. Additionally, the temperature–radius relation has larger scatter for stars of this size relative to hotter and larger stars (Figure 4 in their work).

Intriguingly, we find that stars with  $[\text{Fe}/\text{H}] < -0.2$  dex are assigned larger radii by the Mann et al. (2013b) relations than stars with  $[\text{Fe}/\text{H}] > +0.2$  dex. The mean difference is  $0.05 R_{\odot}$ . The effect persists when applying the temperature–radius relation to temperatures inferred using our EW relations, so the difference must either be a result of the Mann et al. (2013b) temperature–radius relation or of our radius calibration. We suggest that the root cause is the temperature–radius relation for two reasons: (1) we expect a metallicity dependence in the temperature–radius relation from theory (one has yet to present itself in observations of the interferometry stars, but the

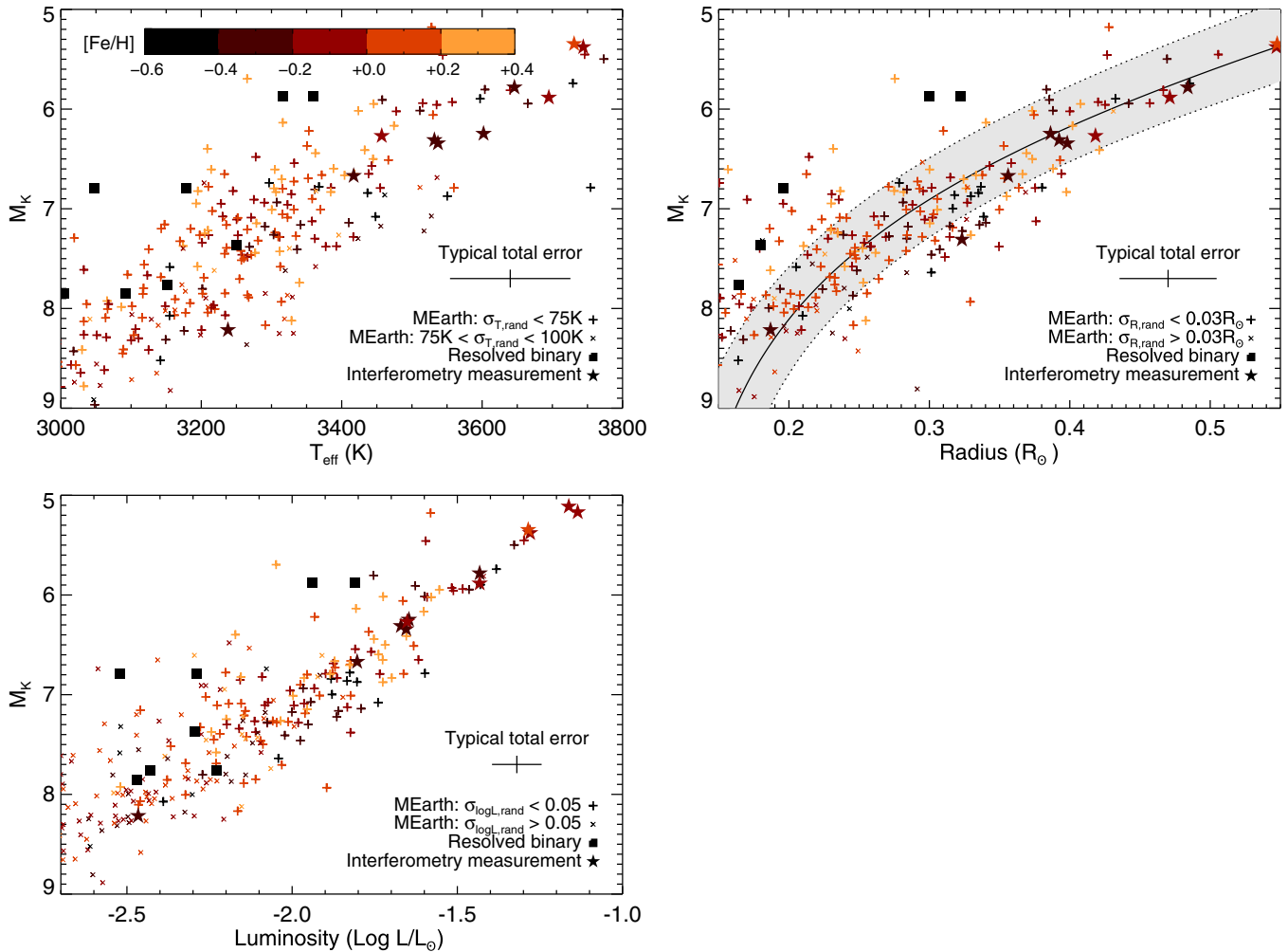
metallicity range spanned by these data is narrow), and (2) we do not see a metallicity dependence when we consider inferred radius as a function of  $M_K$  (Section 4.3).

#### 4.3. Trends Between Absolute $M_K$ and Inferred Stellar Parameters

We consider the relationship between  $M_K$  and the  $T_{\text{eff}}$ , radius, or  $L_{\text{bol}}$  estimated from EWs in Figure 8. NIR photometry is from Two Micron All Sky Survey (2MASS; Skrutskie et al. 2006). We exclude 53 objects with  $\sigma_{M_K} > 0.2$  and 9 objects lacking high quality magnitudes from 2MASS (for which the `qual_flag` is anything other than AAA). For  $T_{\text{eff}}$ , we also exclude those objects that have random errors on  $T_{\text{eff}} > 100$  K as discussed at the beginning of Section 4. We fit a quadratic for  $M_K$  as a function of stellar temperature, radius, or log luminosity; the standard deviation in  $M_K$  is about 0.5 mag for each fit. Because  $M_K$  is an independent indicator of stellar parameters—for example, it was tied to stellar mass by Delfosse et al. (2000)—the existence of a clear relationship between  $M_K$  and the  $T_{\text{eff}}$ , radii, and  $L_{\text{bol}}$  we infer provides additional validation of our method.

Visually, our plot of  $M_K$  versus  $T_{\text{eff}}$  shows the largest scatter. Large variations in  $T_{\text{eff}}$  and  $M_K$  are also seen in the Dartmouth models (Dotter et al. 2008; Feiden et al. 2011), which predict that a star with  $[\text{Fe}/\text{H}] = -0.5$  dex and  $T_{\text{eff}} = 3400$  will be 1 mag fainter than a star with the same temperature but with  $[\text{Fe}/\text{H}] = +0.3$  dex. While neither our  $T_{\text{eff}}$  estimates nor the interferometric measurements indicate such a metallicity dependence, the models still demonstrate the strong influence atmospheric opacities can have on  $T_{\text{eff}}$  (they





**Figure 8.** Absolute  $K$  magnitude vs. inferred  $T_{\text{eff}}$ , radius, or  $L_{\text{bol}}$  for MEarth M dwarfs from Newton et al. (2014). Included are stars with  $\sigma_{M_K} < 0.2$  mag; the symbol indicates the size of the error on stellar parameter (plus symbols for smaller errors or crosses for larger errors). We overplot the properties of the interferometry stars (filled stars), and we show the measured radii, temperatures, and luminosities calculated in Mann et al. (2013b) and in this work. The  $[\text{Fe}/\text{H}]$  of each data point is indicated by its color. We also indicate binaries that are unresolved in 2MASS but resolved in our IRTF observations (filled black squares). In the radius plot (upper right), we also include an  $M_K$ –radius relation that was calculated from the Delfosse et al. (2000)  $M_K$ –mass relation and the Boyajian et al. (2012) mass–radius relation. The shaded region indicates  $\pm 15\%$ .

predict less sensitivity for  $M_K$  versus radius or  $L_{\text{bol}}$ ). Unresolved binaries also contribute significantly to the scatter, which we discuss in Section 4.4.

#### 4.4. Identifying Overluminous Objects

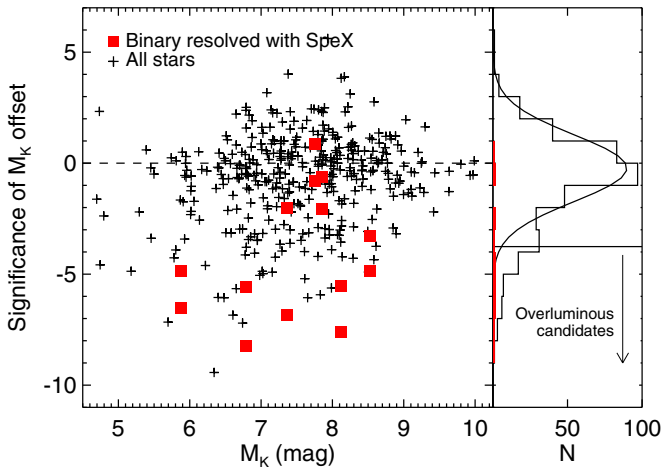
An object that is an unresolved multiple will have an  $M_K$  smaller than that of a single star with the same value of  $L_{\text{bol}}$  predicted from our Equation (3). The EWs of spectral features are largely unchanged by the object’s multiplicity (although for very tight binaries features might appear broadened), but the object appears brighter. Our sample contains binaries that are unresolved in 2MASS but for which independent spectra of the components were obtained—when two stars could be identified in the SpeX guider, Newton et al. (2014) aligned the slit along both components and extracted each object separately. As expected, many of these objects are brighter than single stars with similar spectrally inferred stellar parameters. The diagram of  $M_K$  versus  $L_{\text{bol}}$  shows the excess scatter at brighter magnitudes most clearly.

To quantify the likelihood of an object being a multiple, we calculate the significance of the magnitude offset ( $S$ ) between the object and a quadratic that describes our empirical

$M_K$ –luminosity sequence. We fit  $M_K$  as a function of  $\log L/L_{\odot}$  for stars with  $\log L/L_{\odot} > -2.5$ , which is the limit of our calibration, using  $\sigma_{M_K}$  as our measurement error. We then exclude those objects whose magnitudes are brighter than the best-fitting magnitude by more than the 0.52 mag, which is the standard deviation in the residuals, and re-fit the data. We extrapolate the fit to stars with fainter luminosities.  $S$  is then given by

$$S = \frac{M_K - M_{K,\text{fit}}}{\sqrt{\sigma_{M_K}^2 + \sigma_{M_{K,\text{fit}}}^2}}. \quad (5)$$

We show the distribution of  $S$  values in Figure 9, which is well modeled by a Gaussian with a width of  $\sigma = 1.5$ . An obvious feature is an overabundance of stars that appear too bright for the  $L_{\text{bol}}$  we infer (objects with  $S < 0$ ), and we identify objects with very negative  $S$  values as potential multiples. We note that the width of the Gaussian indicates that our estimates of the errors do not account for all scatter in the relation. One possible contributor is binaries with more extreme mass ratios, as the secondary would increase the combined brightness by a lesser amount.



**Figure 9.** Significance  $S$  vs. absolute  $K$  magnitude for the MEarth M dwarfs.  $S$  is the offset in magnitudes from our empirical  $M_K$ –log luminosity main sequence divided by the error, which includes the error in absolute magnitude and inferred luminosity. On the right side of the figure, we include a histogram of the significances for our entire sample (black histogram) and the best-fitting Gaussian, which has a width of 1.5 (solid curve). We also show the histogram of significances for the IRTF-resolved binaries (solid red histogram). In the histogram panel, we indicate the  $2.5\sigma$  cut we use to select candidate overluminous objects, which are listed in Table 6.

We select those stars with  $S < -3.76$  ( $-2.5\sigma$ ) as candidate multiples. We include the list of candidates in Table 6, which we note does not include the systems identified as visual binaries by Newton et al. (2014), which can be found in that work. Eighteen of these stars are in fact known binaries that were not resolved in our IRTF observations, as indicated in the table. Others may be previously unidentified multiples. Whereas youth is another possibility for explaining their brightnesses, a candidate overluminous object would need to be younger than several hundred Myr to show an enhanced  $K$ -band magnitude, which is unlikely for a field M dwarf.

## 5. APPLICATIONS TO KEPLER OBJECTS OF INTEREST

We apply our calibrations to the high fidelity sample of 66 KOIs selected as described in Section 2.3. The KOIs are primarily early M to mid K dwarfs, with a small number of mid M dwarfs. Muirhead et al. (2014) identified three KOIs as visual binaries (two of which meet our sky emission cuts) and KOI 256 as a white dwarf–M dwarf eclipsing binary, which we exclude from this analysis. Unlike the MEarth objects, the EWs for the KOIs do not approach zero, so it is not necessary to use the limits on the errors we adopted for the MEarth objects. However, for ease of comparison, we restrict the discussion in this section to those KOIs with random errors less than 150 K on  $T_{\text{eff}}$ . This leaves 51 non-visual binaries in our KOI sample.

The cool KOI sample has been the subject of several recent works, and as part of this work we compare the stellar parameters we infer to those presented in Dressing & Charbonneau (2013), Mann et al. (2013b), and Muirhead et al. (2014). Dressing & Charbonneau (2013) matched observed colors to Dartmouth models of different ages, metallicities, and masses. From the best-fitting model, they revised the temperatures and radii for all the M dwarfs in *Kepler* with formal errors typically between 50–100 K for  $T_{\text{eff}}$  and  $0.06 R_{\odot}$  for radius. Mann et al. (2013b) fit flux-calibrated optical spectra to PHOENIX models to determine the  $T_{\text{eff}}$  of *Kepler* M dwarfs with an estimated error of 57 K, and used relations derived from stars with interferometric

**Table 6**  
Candidate Overluminous Objects

Object	$ S ^a$	Binary Ref. <sup>b</sup>	Binary Type <sup>c</sup>
LSPM J0008+2050	4.8	B04	VB
LSPM J0028+5022	4.1	D07	VB
LSPM J0105+2829	4.1	...	...
LSPM J0111+1526	4.9	B04	VB
LSPM J0159+0331E	4.6	S10	SB2
LSPM J0259+3855	3.9	L08	VB
LSPM J0409+0546	4.1	L08	VB
LSPM J0438+2813	4.0	B04	VB
LSPM J0528+1231	6.0	...	...
LSPM J0540+2448	5.4	D76	VB
LSPM J0711+4329	4.2	M06	VB
LSPM J0736+0704	5.1	H97	VB
LSPM J0810+0109	3.9	...	...
LSPM J0835+1408	6.0	...	...
LSPM J0918+6037W	4.1	...	...
LSPM J1000+3155	4.7	...	...
LSPM J1233+0901	5.7	WDS <sup>d</sup>	VB
LSPM J1331+2916	7.1	G02	SB2
		B04	VB
LSPM J1332+3059	4.2	...	...
LSPM J1419+0254	3.7	...	...
LSPM J1547+2241	5.6	...	...
LSPM J1555+3512	4.2	M01	VB
LSPM J1604+2331	5.0	...	...
LSPM J1616+5839	4.8	...	...
LSPM J1707+0722	3.8	P05	VB
LSPM J1733+1655	5.9	...	...
LSPM J1841+2447S	6.8	G96	SB2
LSPM J2010+0632	4.6	S10	SB2
LSPM J2040+1954	3.9	WDS <sup>e</sup>	VB
LSPM J2117+6402	7.1	...	...
LSPM J2223+3227	4.8	W60	VB

### Notes.

<sup>a</sup> Absolute value of the significance of the magnitude offset from our empirical main sequence as defined by Equation (5). All objects are overluminous given their inferred luminosities.

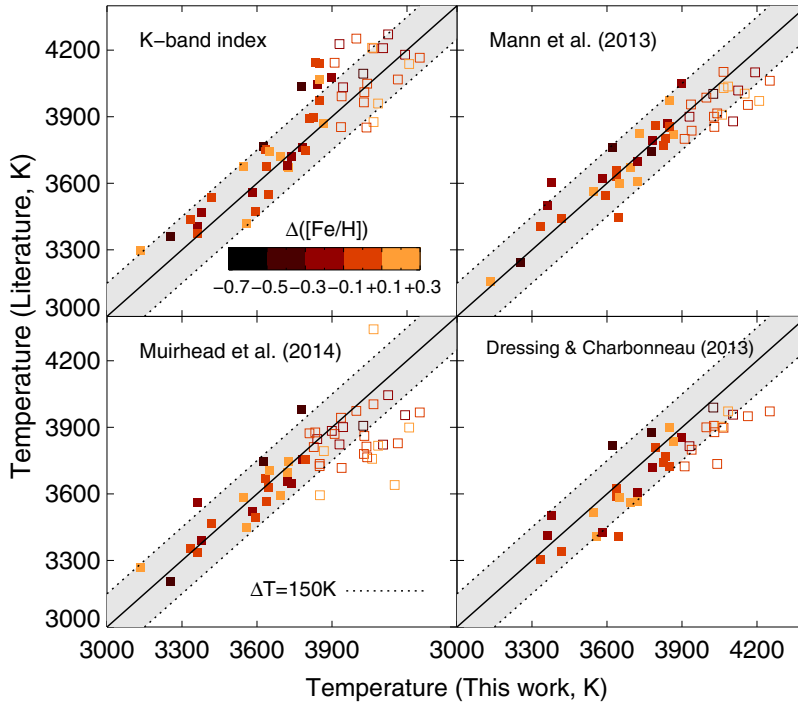
<sup>b</sup> Reference for previous discovery of the object as binary. If no reference is listed, our literature search did not identify that the object is known to be binary. References—W60: Worley 1960; D76: Dahn et al. 1976; G96: Gizis & Reid 1996; H97: Henry et al. 1997; M01: McCarthy et al. 2001; G02: Gizis et al. 2002; B04: Beuzit et al. 2004; P05: Pravdo et al. 2005; M06: Montagnier et al. 2006; D07: Daemgen et al. 2007; L08: Law et al. 2008; S10: Shkolnik et al. 2010.

<sup>c</sup> Type of binary: VB—visual binary; SB2—double-lined spectroscopic binary.

<sup>d</sup> This object is the well-known binary Wolf 424 (Gl 473AB), which is identified in the Washington Double Star catalog. We confirmed that its original discovery was by Reuyl (1938).

<sup>e</sup> This object is listed in the Washington Double Star Catalog as a visual binary, discovered by Riddle et al. (2015) using Robo-AO, who identify it as a possible chance alignment.

measurements to calculate radius and luminosity from  $T_{\text{eff}}$ . Their radius and luminosity errors are 7% and 13%, respectively. They refined and tested their method using the interferometric sample that is also the basis for our work. We note that when fitting spectra without continuum removal, the spectral shape—and therefore the color—is an important determinant. Muirhead et al. (2014) followed the method developed in Muirhead et al. (2012a), using moderate resolution  $K$ -band spectra to determine temperatures and metallicities for 103 cool KOIs. Their temperatures are based on the  $\text{H}_2\text{O}$ – $\text{K}_2$  index (Covey et al. 2010), which was calibrated by Rojas-Ayala et al. (2012)



**Figure 10.** Comparison between the temperatures we infer for the KOIs and those inferred from previous works. Our estimates are on the horizontal axis, and those from the literature are on the vertical axis. The literature works are indicated on each panel. We show the one-to-one line and indicate  $\pm 150$  K as the shaded region. Stars with  $T_{\text{eff}} > 3900$  K ( $T_{\text{eff}} > 3800$  K for Muirhead et al. 2014) are shown as open squares; due to selection effects, we limit our discussion to cooler stars.

as a tracer of  $T_{\text{eff}}$  using measurements of BT-Settl model spectra. Typical formal errors for  $T_{\text{eff}}$  are 65 K. The authors then inferred the stars’ radii by interpolating the  $T_{\text{eff}}$  and metallicities onto Dartmouth isochrones, achieving a median formal error on stellar radius of  $0.06 R_{\odot}$ . We note that Muirhead et al. (2014) used new stellar models, and provided updated parameters for those objects that were first presented in Muirhead et al. (2012b).

### 5.1. Comparison to Previous Work

We compare our inferred temperatures and radii to those from previous works in Figures 10 and 11. The literature sources we query are Mann et al. (2013b), Muirhead et al. (2014), and Dressing & Charbonneau (2013). We also apply the K-band calibration from Mann et al. (2013b) to the KOI spectra to estimate  $T_{\text{eff}}$ , after which we use their temperature–radius relation to estimate radii (see Section 4.2). We adopt the metallicities from Muirhead et al. (2014).

We cannot make statements about temperature differences above 3900 K. The literature sources we query have their own target selection criteria, each of which places an upper limit of roughly 4000 K on the temperatures of the stars in their sample. This is particularly important for the sample from Dressing & Charbonneau (2013), which is strictly limited to stars with updated temperatures of less than 4000 K. Mann et al. (2013b) and Muirhead et al. (2014) use color cuts to select red objects, which could also bias the resulting temperatures because their methods are related to spectral colors. Finally, the  $\text{H}_2\text{O}$ –K2 index used by Muirhead et al. (2014) saturates for  $T_{\text{eff}} > 3800$  K.

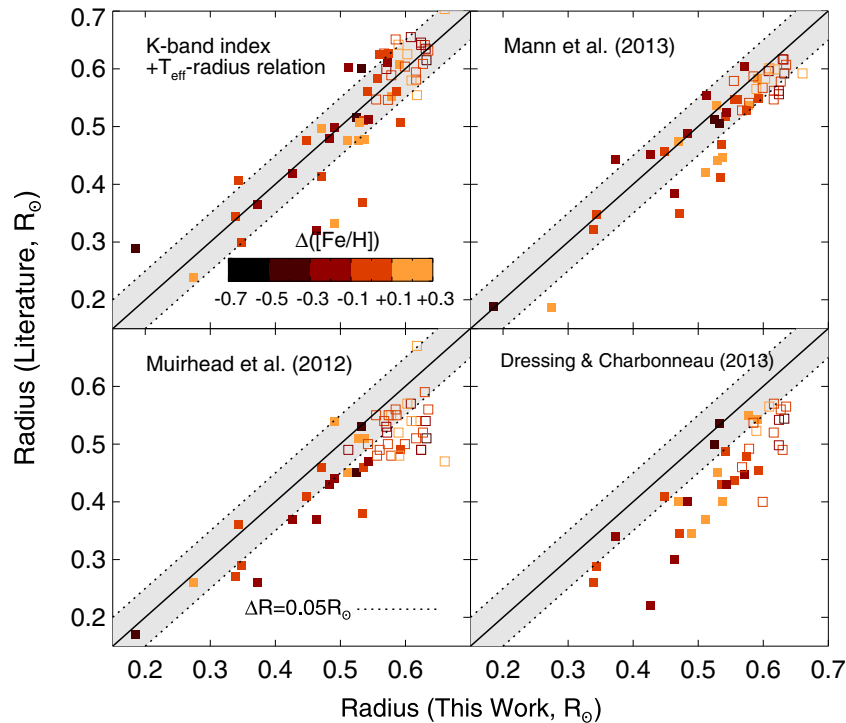
The temperatures we find generally agree well with those from other works. They are cooler than those inferred from previous spectra-based approaches: the median temperature difference is  $80 \pm 90$  K considering the K-band index,  $10 \pm 60$  K considering Mann et al. (2013b), and  $30 \pm 70$  K considering Muirhead et al. (2014). Compared to Dressing & Charbonneau

(2013), the temperatures we infer are hotter by a median of  $40 \pm 90$  K. The errors we quote are the median absolute deviations in the temperature differences.

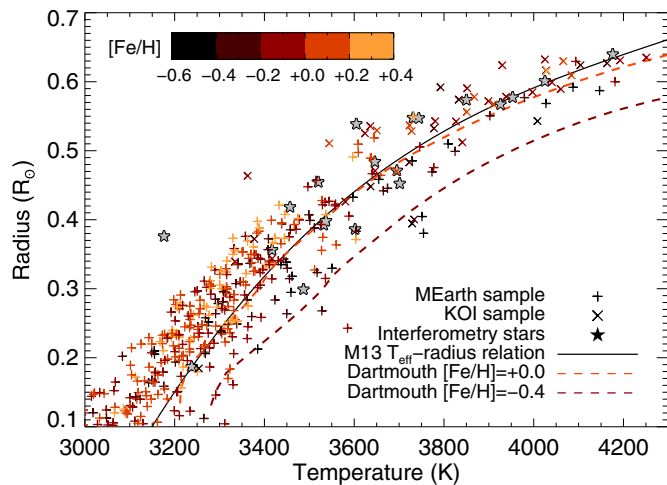
The radii we infer are larger than those found by Mann et al. (2013b), with a median difference of  $0.02 R_{\odot}$  and a median absolute deviation of  $0.04 R_{\odot}$ . This difference may be due to their use of a temperature–radius relation, which assumes that all stars at a given temperature have the same radius, but is within reasonable systematic uncertainties for the calibrations.

Our radii are also larger than the radii estimated using methods based on model isochrones by  $0.05 \pm 0.03 R_{\odot}$  relative to Muirhead et al. (2014) and  $0.09 \pm 0.04 R_{\odot}$  relative to Dressing & Charbonneau (2013). Because our temperatures are in agreement, the difference in the radii must result from their use of models to estimate radius from temperature. Indeed, Boyajian et al. (2012, Figure 14) found that for stars with interferometrically measured radii between  $0.4 < R/R_{\odot} < 0.6$ , the model radii—predicted by interpolating each stars’  $T_{\text{eff}}$  onto Dartmouth models—are smaller than the measured radii by about 10%, or  $0.05 R_{\odot}$ . We illustrate the disagreement between observations and models by comparing the temperatures and radii of the interferometric sample to Dartmouth isochrones in Figure 12. This effect is sufficient to explain the differences between our radii and those from Muirhead et al. (2014).

The discrepancy between models and observations also affects the radii from Dressing & Charbonneau (2013). However, the difference between our radii and those from Dressing & Charbonneau (2013) is larger because of the metallicities they estimate. In the Dartmouth models, an M dwarf with  $[\text{Fe}/\text{H}] = -0.2$  is about  $0.04 R_{\odot}$  smaller than a solar-metallicity dwarf of the same temperature; such a metallicity dependence is not seen in the temperatures and radii of the interferometry stars. Dressing & Charbonneau (2013) estimated sub-solar metallicities for most of the KOIs, whereas the metallicities from Muirhead et al. (2014) are closer to solar (see Figure 7 in



**Figure 11.** Comparison between the radii we infer for the KOIs and those inferred from previous works. Our estimates are on the horizontal axis, and those from the literature are on the vertical axis. The literature works are indicated on each panel. We show the one-to-one line and indicate  $\pm 0.05 R_{\odot}$  using the shaded region. Stars with  $T_{\text{eff}} > 3900$  K ( $> 3800$  K for Muirhead et al. 2014) are shown as open squares; due to selection effects, we limit our discussion to cooler stars.



**Figure 12.** Temperature–radius diagram for the stars in this work, showing the measured values for the interferometric sample (gray stars) and the values we infer using EWs for the MEarth sample (plus symbols) and the KOIs (crosses). We color the MEarth and KOI samples by their estimated metallicities. For MEarth, we estimate metallicities using relations from Mann et al. (2013a, 2014), and for the KOIs we use the metallicities from Muirhead et al. (2014). We also include the Mann et al. (2013b) temperature–radius relation (solid black line) and 5 Gyr, solar alpha-enhancement Dartmouth isochrones for  $[\text{Fe}/\text{H}] = 0.0$  (dashed orange line) and  $[\text{Fe}/\text{H}] = -0.4$  (dashed dark red line). Gl 725B (measured  $T_{\text{eff}} = 3142$  K,  $R = 0.32 R_{\odot}$ ) was used to calibrate our radius relation, but not our temperature relation, and is not shown. The largest outlier from the interferometric sample is Gl 876 (measured  $T_{\text{eff}} = 3176$  K,  $R = 0.38 R_{\odot}$ ).

Dressing & Charbonneau 2013), and Dressing & Charbonneau (2013) therefore find smaller radii. This can also be seen in Figure 12. The overall offset between the observed and theoretical temperatures and radii, and the sub-solar metallicities they

estimated are what cause our inferred radii to be substantially larger than those reported in Dressing & Charbonneau (2013).

## 5.2. Updated Stellar and Planetary Parameters

We use the data available on the NASA Exoplanet Archive<sup>6</sup> to update the properties of the planet candidates orbiting the cool KOIs. We compare the planet properties that we get using our stellar parameters to those that one would infer using the stellar parameters in the catalog from Huber et al. (2014). Huber et al. (2014) synthesized stellar parameters available in the literature for objects in the *Kepler* Input Catalog. For the M dwarfs, measurements primarily come from Muirhead et al. (2012b) and Dressing & Charbonneau (2013).

To update the planetary radii we use the planet-to-star radius ratio ( $r/R_*$ ) and calculate  $r$  using either our new stellar radius or that from Huber et al. (2014). To update planet equilibrium temperatures ( $T_{\text{eq}}$ ) we assume that the planet radiates the same amount of heat as it receives, that heat is distributed isotropically, that the planet has an albedo of  $a = 0.3$ , and that the star and planet radiate as blackbodies. This gives the familiar equation  $T_{\text{eq}} = T_{\odot} \times (1 - a)^{1/4} \sqrt{R_*/2d}$ . The ratio between the planet–star distance and the stellar radius ( $d/R_*$ ) is another directly measured transit parameter. We calculate  $T_{\text{eq}}$  using either our new stellar effective temperature or that from Huber et al. (2014).

We present our updated stellar and planetary parameters in Table 7. In Figure 13 we show how the planet radii and equilibrium temperatures change when using our updated parameters. The difference in equilibrium temperature is largely negligible, but our new stellar radii have a significant effect on the radii

<sup>6</sup> <http://exoplanetarchive.ipac.caltech.edu/>, accessed 2014 June 10.



**Table 7**  
Stellar Parameters of Cool Kepler Objects of Interest with High-fidelity Spectra

KOI Num.	Kepler ID	Huber et al. (2014)			This Work			Literature <sup>a</sup>		This Work	
		$T_{\text{eff}}^b$ (K)	Radius <sup>b</sup> ( $R_{\odot}$ )	Ref. <sup>c</sup>	$T_{\text{eff}}$ (K)	Radius ( $R_{\odot}$ )	Luminosity ( $\log L/L_{\odot}$ )	$R_p$ ( $R_{\oplus}$ )	$T_{\text{eq}}$ (K)	$R_p$ ( $R_{\oplus}$ )	$T_{\text{eq}}$ (K)
247.01	11852982	3741	0.49	SPE5	3850 $\pm$ 75	0.556 $\pm$ 0.027	−1.18 $\pm$ 0.06	1.85	507	2.10	522
250.01	9757613	3887	0.53	SPE5	3899 $\pm$ 100	0.572 $\pm$ 0.028	−1.10 $\pm$ 0.14	2.83	453	3.05	455
250.02	9757613	3887	0.53	SPE5	3899 $\pm$ 100	0.572 $\pm$ 0.028	−1.10 $\pm$ 0.14	2.70	405	2.91	406
250.03	9757613	3887	0.53	SPE5	3899 $\pm$ 100	0.572 $\pm$ 0.028	−1.10 $\pm$ 0.14	1.04	686	1.12	688
250.04	9757613	3887	0.53	SPE5	3899 $\pm$ 100	0.572 $\pm$ 0.028	−1.10 $\pm$ 0.14	2.73	290	2.94	291
251.01	10489206	3810	0.52	SPE5	3827 $\pm$ 84	0.542 $\pm$ 0.029	−1.20 $\pm$ 0.07	2.61	665	2.72	668
251.02	10489206	3810	0.52	SPE5	3827 $\pm$ 84	0.542 $\pm$ 0.029	−1.20 $\pm$ 0.07	0.83	595	0.86	598
254.01	5794240	3820	0.55	SPE43	3867 $\pm$ 82	0.578 $\pm$ 0.028	−0.97 $\pm$ 0.10	10.23	691	10.71	700
255.01	7021681	3770	0.51	SPE5	4027 $\pm$ 81	0.616 $\pm$ 0.027	−1.01 $\pm$ 0.12	2.46	334	2.98	357
255.02	7021681	3770	0.51	SPE5	4027 $\pm$ 81	0.616 $\pm$ 0.027	−1.01 $\pm$ 0.12	0.74	497	0.90	531
314.01	7603200	3841	0.50	SPE5	3841 $\pm$ 73	0.512 $\pm$ 0.027	−1.33 $\pm$ 0.05	1.58	467	1.62	467
314.02	7603200	3841	0.50	SPE5	3841 $\pm$ 73	0.512 $\pm$ 0.027	−1.33 $\pm$ 0.05	1.68	392	1.72	392
314.03	7603200	3841	0.50	SPE5	3841 $\pm$ 73	0.512 $\pm$ 0.027	−1.33 $\pm$ 0.05	0.64	515	0.66	515
463.01	8845205	3387	0.30	SPE5	3377 $\pm$ 79	0.372 $\pm$ 0.028	−1.71 $\pm$ 0.06	1.50	244	1.87	244
478.01	10990886	3744	0.50	SPE5	3727 $\pm$ 74	0.529 $\pm$ 0.027	−1.27 $\pm$ 0.06	2.63	468	2.78	466
531.01	10395543	4030	0.60	SPE5	4065 $\pm$ 76	0.630 $\pm$ 0.028	−0.92 $\pm$ 0.07	3.40	490	3.57	495
571.01	8120608	3761	0.46	SPE5	3624 $\pm$ 79	0.525 $\pm$ 0.029	−1.32 $\pm$ 0.07	1.44	705	1.64	679
571.02	8120608	3761	0.46	SPE5	3624 $\pm$ 79	0.525 $\pm$ 0.029	−1.32 $\pm$ 0.07	1.59	576	1.81	555
571.03	8120608	3761	0.46	SPE5	3624 $\pm$ 79	0.525 $\pm$ 0.029	−1.32 $\pm$ 0.07	1.21	870	1.39	839
571.04	8120608	3761	0.46	SPE5	3624 $\pm$ 79	0.525 $\pm$ 0.029	−1.32 $\pm$ 0.07	1.45	484	1.66	466
571.05	8120608	3761	0.46	SPE5	3624 $\pm$ 79	0.525 $\pm$ 0.029	−1.32 $\pm$ 0.07	1.02	182	1.17	175
596.01	10388286	3678	0.47	SPE5	3635 $\pm$ 78	0.536 $\pm$ 0.028	−1.35 $\pm$ 0.08	1.29	874	1.47	864
818.01	4913852	3721	0.52	SPE5	3723 $\pm$ 86	0.537 $\pm$ 0.028	−1.36 $\pm$ 0.10	2.34	581	2.41	582
854.01	6435936	3593	0.47	SPE5	3694 $\pm$ 85	0.470 $\pm$ 0.032	−1.12 $\pm$ 0.06	2.05	271	2.05	279
898.01	7870390	3893	0.52	SPE5	4025 $\pm$ 99	0.632 $\pm$ 0.038	−0.71 $\pm$ 0.09	2.42	530	2.94	548
898.02	7870390	3893	0.52	SPE5	4025 $\pm$ 99	0.632 $\pm$ 0.038	−0.71 $\pm$ 0.09	1.86	656	2.27	679
898.03	7870390	3893	0.52	SPE5	4025 $\pm$ 99	0.632 $\pm$ 0.038	−0.71 $\pm$ 0.09	2.02	419	2.45	433
899.01	7907423	3568	0.42	SPE5	3636 $\pm$ 77	0.448 $\pm$ 0.030	−1.28 $\pm$ 0.06	1.34	554	1.43	565
899.02	7907423	3568	0.42	SPE5	3636 $\pm$ 77	0.448 $\pm$ 0.030	−1.28 $\pm$ 0.06	1.02	715	1.09	729
899.03	7907423	3568	0.42	SPE5	3636 $\pm$ 77	0.448 $\pm$ 0.030	−1.28 $\pm$ 0.06	1.31	428	1.40	436
936.01	9388479	3581	0.44	SPE5	3544 $\pm$ 78	0.511 $\pm$ 0.028	−1.40 $\pm$ 0.07	2.28	520	2.65	515
936.02	9388479	3581	0.44	SPE5	3544 $\pm$ 78	0.511 $\pm$ 0.028	−1.40 $\pm$ 0.07	1.30	1143	1.51	1132
947.01	9710326	3750	0.46	SPE5	3780 $\pm$ 92	0.543 $\pm$ 0.034	−1.24 $\pm$ 0.07	2.23	450	2.63	454
961.01	8561063	3068	0.17	SPE41	3254 $\pm$ 106	0.185 $\pm$ 0.043	−2.73 $\pm$ 0.07	0.86	570	0.94	605
961.02	8561063	3068	0.17	SPE41	3254 $\pm$ 106	0.185 $\pm$ 0.043	−2.73 $\pm$ 0.07	0.79	790	0.86	838
961.03	8561063	3068	0.17	SPE41	3254 $\pm$ 106	0.185 $\pm$ 0.043	−2.73 $\pm$ 0.07	0.77	494	0.83	524
1085.01	10118816	3939	0.52	SPE5	3777 $\pm$ 97	0.532 $\pm$ 0.030	−1.07 $\pm$ 0.09	1.05	634	1.07	608
1397.01	9427402	3957	0.54	PHO2	4104 $\pm$ 107	0.624 $\pm$ 0.032	−1.21 $\pm$ 0.16	2.01	501	2.31	520
1408.01	9150827	4023	0.57	SPE5	4192 $\pm$ 84	0.631 $\pm$ 0.029	−0.78 $\pm$ 0.07	1.31	462	1.45	482
1408.02	9150827	4023	0.57	SPE5	4192 $\pm$ 84	0.631 $\pm$ 0.029	−0.78 $\pm$ 0.07	0.80	260	0.89	271
1422.01	11497958	3517	0.37	SPE5	3580 $\pm$ 98	0.426 $\pm$ 0.031	−1.50 $\pm$ 0.08	1.41	457	1.63	465
1422.02	11497958	3517	0.37	SPE5	3580 $\pm$ 98	0.426 $\pm$ 0.031	−1.50 $\pm$ 0.08	1.45	303	1.67	309
1422.03	11497958	3517	0.37	SPE5	3580 $\pm$ 98	0.426 $\pm$ 0.031	−1.50 $\pm$ 0.08	1.13	373	1.30	380
1422.04	11497958	3517	0.37	SPE5	3580 $\pm$ 98	0.426 $\pm$ 0.031	−1.50 $\pm$ 0.08	1.19	206	1.37	210
1422.05	11497958	3517	0.37	SPE5	3580 $\pm$ 98	0.426 $\pm$ 0.031	−1.50 $\pm$ 0.08	1.03	254	1.19	258
1649.01	11337141	3767	0.48	PHO2	3833 $\pm$ 102	0.574 $\pm$ 0.030	−1.29 $\pm$ 0.13	0.98	580	1.18	590
1681.01	5531953	3608	0.40	PHO2	3722 $\pm$ 132	0.483 $\pm$ 0.033	−1.46 $\pm$ 0.09	1.00	470	1.21	485
1681.02	5531953	3608	0.40	PHO2	3722 $\pm$ 132	0.483 $\pm$ 0.033	−1.46 $\pm$ 0.09	0.88	1193	1.07	1231
1681.03	5531953	3608	0.40	PHO2	3722 $\pm$ 132	0.483 $\pm$ 0.033	−1.46 $\pm$ 0.09	0.78	647	0.95	667
1702.01	7304449	3304	0.26	PHO2	3334 $\pm$ 99	0.339 $\pm$ 0.031	−1.91 $\pm$ 0.08	0.82	796	1.07	803
1843.01	5080636	3584	0.45	PHO2	3650 $\pm$ 92	0.529 $\pm$ 0.031	−1.40 $\pm$ 0.09	1.16	536	1.37	546
1843.02	5080636	3584	0.45	PHO2	3650 $\pm$ 92	0.529 $\pm$ 0.031	−1.40 $\pm$ 0.09	0.73	463	0.85	472
1867.01	8167996	3799	0.49	PHO2	3938 $\pm$ 112	0.578 $\pm$ 0.029	−1.20 $\pm$ 0.12	1.13	737	1.33	764
1867.02	8167996	3799	0.49	PHO2	3938 $\pm$ 112	0.578 $\pm$ 0.029	−1.20 $\pm$ 0.12	2.01	415	2.36	430
1867.03	8167996	3799	0.49	PHO2	3938 $\pm$ 112	0.578 $\pm$ 0.029	−1.20 $\pm$ 0.12	1.01	579	1.18	600
1868.01	6773862	3950	0.56	PHO2	4163 $\pm$ 133	0.627 $\pm$ 0.035	−0.94 $\pm$ 0.12	2.14	316	2.39	334
1902.01	5809954	3763	0.46	PHO16	3737 $\pm$ 114	0.490 $\pm$ 0.032	−1.36 $\pm$ 0.12	18.43	172	19.77	171
1907.01	7094486	3901	0.54	PHO2	3851 $\pm$ 109	0.591 $\pm$ 0.033	−1.24 $\pm$ 0.17	2.01	476	2.19	470
2006.01	10525027	3809	0.46	PHO2	3792 $\pm$ 93	0.592 $\pm$ 0.030	−1.24 $\pm$ 0.10	0.77	732	1.00	729
2036.01	6382217	3903	0.52	PHO2	4060 $\pm$ 112	0.589 $\pm$ 0.031	−1.03 $\pm$ 0.15	1.49	504	1.68	525
2036.02	6382217	3903	0.52	PHO2	4060 $\pm$ 112	0.589 $\pm$ 0.031	−1.03 $\pm$ 0.15	0.96	568	1.09	591
2057.01	9573685	3900	0.54	PHO2	3997 $\pm$ 118	0.585 $\pm$ 0.029	−1.02 $\pm$ 0.11	1.14	636	1.24	652
2130.01	2161536	3972	0.56	PHO2	4251 $\pm$ 130	0.635 $\pm$ 0.037	−0.75 $\pm$ 0.09	1.72	367	1.94	392

**Table 7**  
(Continued)

KOI Num.	Kepler ID	Huber et al. (2014)			This Work			Literature <sup>a</sup>		This Work	
		$T_{\text{eff}}^b$ (K)	Radius <sup>b</sup> ( $R_{\odot}$ )	Ref. <sup>c</sup>	$T_{\text{eff}}$ (K)	Radius ( $R_{\odot}$ )	Luminosity ( $\log L/L_{\odot}$ )	$R_p$ ( $R_{\oplus}$ )	$T_{\text{eq}}$ (K)	$R_p$ ( $R_{\oplus}$ )	$T_{\text{eq}}$ (K)
2191.01	5601258	3724	0.46	PHO2	3910 $\pm$ 107	0.567 $\pm$ 0.030	−1.24 $\pm$ 0.11	1.15	584	1.42	613
2306.01	6666233	3878	0.52	PHO2	4029 $\pm$ 89	0.616 $\pm$ 0.029	−0.92 $\pm$ 0.10	0.94	1546	1.11	1607
2329.01	11192235	3815	0.50	PHO2	3929 $\pm$ 135	0.624 $\pm$ 0.036	−1.32 $\pm$ 0.13	1.26	983	1.58	1012
2347.01	8235924	3972	0.56	PHO2	4084 $\pm$ 108	0.609 $\pm$ 0.029	−1.25 $\pm$ 0.14	1.07	1637	1.16	1684
2542.01	6183511	3339	0.29	PHO2	3417 $\pm$ 113	0.344 $\pm$ 0.041	−1.65 $\pm$ 0.11	0.60	1365	0.72	1397
2650.01	8890150	3735	0.40	PHO2	4040 $\pm$ 102	0.599 $\pm$ 0.028	−1.12 $\pm$ 0.10	0.98	319	1.47	346
2650.02	8890150	3735	0.40	PHO2	4040 $\pm$ 102	0.599 $\pm$ 0.028	−1.12 $\pm$ 0.10	0.86	545	1.29	590
2662.01	3426367	3410	0.34	PHO2	3646 $\pm$ 128	0.471 $\pm$ 0.031	−1.30 $\pm$ 0.11	0.69	816	0.94	872
2704.01	9730163	3327	0.19	PHO54	3134 $\pm$ 102	0.274 $\pm$ 0.034	−2.17 $\pm$ 0.10	2.02	414	2.92	390
2704.02	9730163	3327	0.19	PHO54	3134 $\pm$ 102	0.274 $\pm$ 0.034	−2.17 $\pm$ 0.10	1.36	529	1.97	499
2705.01	11453592	3400	0.27	PHO54	3592 $\pm$ 134	0.534 $\pm$ 0.044	−1.52 $\pm$ 0.18	1.39	876	2.79	925
2715.01	9837661	4385	0.71	KIC0	4150 $\pm$ 129	0.660 $\pm$ 0.029	−0.90 $\pm$ 0.22	6.83	588	6.38	557
2715.02	9837661	4385	0.71	KIC0	4150 $\pm$ 129	0.660 $\pm$ 0.029	−0.90 $\pm$ 0.22	3.69	1035	3.45	980
2715.03	9837661	4385	0.71	KIC0	4150 $\pm$ 129	0.660 $\pm$ 0.029	−0.90 $\pm$ 0.22	3.16	879	2.95	832
2764.01	10073672	3952	0.55	PHO2	4124 $\pm$ 117	0.608 $\pm$ 0.029	−1.02 $\pm$ 0.15	1.59	1204	1.76	1257
2839.01	6186964	3900	0.54	PHO2	3935 $\pm$ 136	0.555 $\pm$ 0.031	−1.13 $\pm$ 0.14	1.29	740	1.33	747
2845.01	10591855	3954	0.55	PHO2	4066 $\pm$ 146	0.618 $\pm$ 0.031	−0.78 $\pm$ 0.22	0.84	940	0.95	966
2926.01	10122538	3903	0.52	PHO2	4208 $\pm$ 157	0.603 $\pm$ 0.034	−0.95 $\pm$ 0.16	2.28	515	2.63	555
2926.02	10122538	3903	0.52	PHO2	4208 $\pm$ 157	0.603 $\pm$ 0.034	−0.95 $\pm$ 0.16	2.00	610	2.30	658
2926.03	10122538	3903	0.52	PHO2	4208 $\pm$ 157	0.603 $\pm$ 0.034	−0.95 $\pm$ 0.16	2.45	364	2.82	392
2926.04	10122538	3903	0.52	PHO2	4208 $\pm$ 157	0.603 $\pm$ 0.034	−0.95 $\pm$ 0.16	2.36	310	2.72	335
2992.01	8509442	3952	0.55	PHO2	4088 $\pm$ 141	0.578 $\pm$ 0.034	−0.90 $\pm$ 0.19	2.04	184	2.14	190
3090.01	6609270	3854	0.53	PHO2	3850 $\pm$ 128	0.560 $\pm$ 0.030	−1.29 $\pm$ 0.16	1.17	748	1.24	747
3090.02	6609270	3854	0.53	PHO2	3850 $\pm$ 128	0.560 $\pm$ 0.030	−1.29 $\pm$ 0.16	2.14	694	2.26	693
3282.01	12066569	3894	0.54	PHO2	3944 $\pm$ 127	0.571 $\pm$ 0.031	−1.23 $\pm$ 0.13	2.24	305	2.36	309
3414.01	6023859	3900	0.54	PHO2	3834 $\pm$ 159	0.563 $\pm$ 0.039	−1.32 $\pm$ 0.16	18.48	239	19.37	235
3749.01	11547869	3311	0.22	PHO2	3362 $\pm$ 112	0.348 $\pm$ 0.039	−1.81 $\pm$ 0.10	8.17	314	12.62	319
4252.01	10525049	3842	0.53	PHO2	3809 $\pm$ 117	0.586 $\pm$ 0.033	−1.36 $\pm$ 0.22	0.67	344	0.74	341
4427.01	4172805	3668	0.43	PHO2	4037 $\pm$ 155	0.573 $\pm$ 0.050	−1.20 $\pm$ 0.14	1.46	176	1.94	194

**Notes.**<sup>a</sup> Planet properties one infers using the stellar parameters from Huber et al. (2014) and the planet parameters from the NASA Exoplanet Archive.<sup>b</sup> Stellar parameter in catalog from Huber et al. (2014).<sup>c</sup> Reference for Huber et al. (2014) data. SPE5: Muirhead et al. (2014); PHO2: Dressing & Charbonneau (2013); SPE41: Muirhead et al. (2012b); SPE43: Johnson et al. (2012); PHO16: Gaidos (2013); PHO54: Huber et al. (2014); and KIC0: Brown et al. (2011).

(This table is available in its entirety in machine-readable form.)

of orbiting planets—the typical planet is 15% larger with our stellar parameters than with those in Huber et al. (2014).

### 5.3. Comments on Individual Systems

Two KOIs stand out because our new radii are smaller than those in Huber et al. (2014). These are the candidate planets orbiting KOI 2715, for which the previous best stellar parameters come from the *Kepler* Input Catalog (Brown et al. 2011).

Three *Kepler* targets in our sample have previously received significant attention. KOI 961 (*Kepler*-42) hosts a suite of sub-Earth-sized planets and was analyzed by Muirhead et al. (2012a), who inferred this star’s properties by tying models to Barnard’s star, which has a directly measured radius. While the temperature we estimate is nearly 200 K hotter than the temperature from their analysis, the two estimates are consistent (3068  $\pm$  174 K versus 3254  $\pm$  110 K). The radii are also in very good agreement: 0.17  $\pm$  0.04  $R_{\odot}$  in their analysis versus 0.19  $\pm$  0.04 in this work.

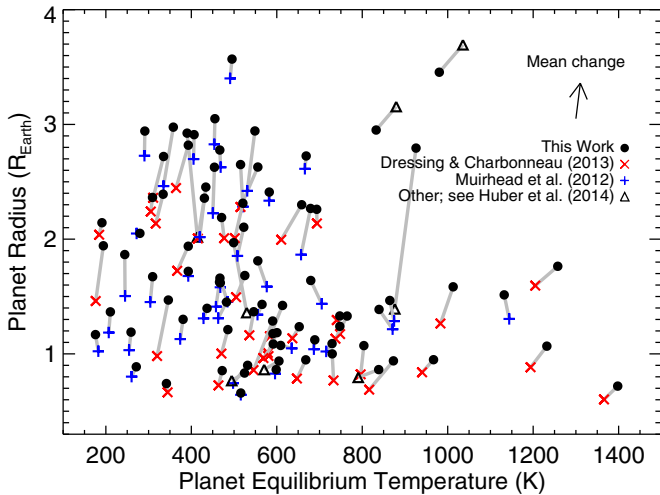
Johnson et al. (2012) presented KOI 254 (*Kepler*-45), an early M dwarf hosting a hot Jupiter, and found  $R = 0.55 \pm 0.11 R_{\odot}$

and  $T_{\text{eff}} = 3820 \pm 90$ , which is again consistent with our results of  $R = 0.58 \pm 0.03 R_{\odot}$  and  $T_{\text{eff}} = 3870 \pm 80$ .

KOI 571 (*Kepler*-186) was recently announced as hosting an Earth-sized planet in its habitable zone (*Kepler*-186f; Quintana et al. 2014). The stellar parameters listed for this star in Huber et al. (2014) and in our Table 7 are from Muirhead et al. (2012b). Quintana et al. (2014) separately determined the radius for this star by finding the Dartmouth model best matching the mean stellar density, determined from transit photometry, and the metallicity and  $T_{\text{eff}}$  from Muirhead et al. (2012b). Their radius of  $0.47 \pm 0.05 R_{\odot}$  is smaller than our estimate of  $0.53 \pm 0.03 R_{\odot}$ , but is still consistent. We also revise the radius of *Kepler*-186f upward, from 1.02  $R_{\oplus}$  to 1.17  $R_{\oplus}$ ; a planet of this size is still likely to be rocky (e.g., Rogers 2014). The  $T_{\text{eff}}$  (3624  $\pm$  80) and luminosity ( $0.048 \pm 0.008 L_{\odot}$ ) we infer are also consistent with the properties from Quintana et al. (2014), who estimated  $T_{\text{eff}} = 3788 \pm 54$  K and  $L = 0.041 \pm 0.009 L_{\odot}$ . Therefore, our results support the conclusion that *Kepler*-186f is a rocky, habitable-zone planet.

## 6. SUMMARY

We presented an empirical calibrations for the effective temperatures, radii, and luminosities of cool dwarfs. We used



**Figure 13.** Planet radius vs. planet equilibrium temperature when using stellar properties from Huber et al. (2014) or updated parameters from this work (filled circles). The Huber et al. (2014) catalog primarily draws M dwarf stellar parameters from Dressing & Charbonneau (2013, red crosses) and from Muirhead et al. (2012b, blue pluses) with small numbers of objects from other works (black open triangles). When stellar properties are from Dressing & Charbonneau (2013), red open squares indicate planet properties from that work, while red crosses indicate planet properties from NExSci. Gray lines connect the previous and updated values for each planet.

24 M dwarfs with interferometrically measured parameters (25 for our radius calibration) to calibrate our relationships, which are based on EWs of  $H$ -band spectral features. Our relationships are applicable to dwarfs with  $3100 < T_{\text{eff}}(\text{K}) < 4800$ ,  $0.2 < R/R_{\odot} < 0.8$ , and  $-2.5 < \log L/L_{\odot} < -0.5$ . The standard deviations in the residuals of the best fits are 73 K,  $0.027 R_{\odot}$ , and 0.049 dex (11%). From our bootstrap analysis, our luminosity calibration is the only one for which systematic error is important, but comparing temperature, radius, and luminosities indicates that there may be additional sources of systematic uncertainty. Our calibrations can be applied to stars without parallaxes and to non-flux calibrated spectra, and can be used very effectively for early M dwarfs with  $3700 < T_{\text{eff}}(\text{K}) < 4000$ , where the  $H_2O$ -K2 index used by Muirhead et al. (2012a, 2014) to estimate temperature saturates. This is an important regime for understanding planets orbiting cool stars because these are the late-type dwarfs with the greatest representation in *Kepler*. Figure 12 summarizes our results—we show the measured parameters for the interferometric stars on which the calibration is based, and the properties we infer for the MEarth and *Kepler* samples.

Our investigation of  $H$ -band spectral features also revealed that the EWs we measure of our features in our observed spectra that are not strongly dependent on metallicity—in particular, Mg I features—show the best agreement with the EWs we measure from synthetic spectra. Conversely, the EWs we measure for features with an apparent metallicity dependence are the most discrepant.

We applied our calibrations to the MEarth sample of M dwarfs and validated the stellar parameters we infer by demonstrating that they display a clear relationship with  $M_K$ , which is an independent tracer of the stars’ underlying physical properties. By comparing  $M_K$  to inferred stellar luminosity, we identified 31 candidate multiples. Eighteen of the objects we identified in this manner are known binaries. We also used the luminosities we estimated to demonstrate that the Newton et al. (2014) metallicity calibration overestimates the metallicities of late

M dwarfs, and updated the metallicities of the sample using the calibrations from Mann et al. (2013a, 2014).

Using spectra from Muirhead et al. (2014), we applied our calibration to the cool stars from *Kepler* that host candidate planets. The temperatures that we find agree remarkably well with the temperatures reported in previous works, particularly given the different methods used, although our new stellar radii are larger. The largest discrepancy (median difference  $0.09 R_{\odot}$ ) arises when we compare our radii to Dressing & Charbonneau (2013), who fit photometry to Dartmouth models to estimate stellar parameters. The primary cause of the discrepancy is that at a given  $T_{\text{eff}}$  the interferometrically measured radii are larger than those predicted by models by about 10% (Boyajian et al. 2012), so the model-based radii from Dressing & Charbonneau (2013) are too small. The sub-solar metallicities they infer for the KOIs also contribute. Using our new stellar parameters, we updated the properties of the candidate planets, finding that the typical planet is larger than what one would calculate using the recent catalog from Huber et al. (2014) by 15%. The properties we infer for KOIs 961 (*Kepler*-42), KOI 254 (*Kepler*-45), and KOI 571 (*Kepler*-186) are consistent with the results from previously published in-depth studies of those objects.

Our new calibrations have the benefit of being independent of stellar models. Because of the discrepancies between theoretical and observed stellar parameters, methods that rely on fixing stellar parameters to models will be subject to systematic errors. We note two important considerations that were discussed by Boyajian et al. (2012). First, stellar radii for these low-mass stars generally are measured to be larger than predicted. Second, the effect of metallicity on radius does not appear to be as strong in the interferometric sample as is predicted by models, so metal-poor M dwarfs may be particularly misrepresented by such methods. We note, however, that we see some evidence of a metallicity dependence in the temperature–radius plane when considering the larger sample of stars to which we have applied our calibrations (see Section 4.2 and Figure 12). Additional interferometric measurements—particularly of mid to late M dwarfs and stars of extreme metallicity—would improve both our empirical calibrations and our general understanding of the physical properties of low-mass stars.

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