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Topological susceptibility from $N_f = 2 + 1 + 1$ lattice QCD at nonzero temperature

Anton Trunin$^{1,2,*}$, Florian Burger$^3$, Ernst–Michael Ilgenfritz$^1$, Maria Paola Lombardo$^4$ and Michael Müller–Preussker$^3$

1 Bogoliubov Laboratory of Theoretical Physics, Joint Institute for Nuclear Research, Joliot–Curie Str. 6, 141980 Dubna, Russia
2 Samara State Aerospace University, Moskovskoye Shosse 34, 443086 Samara, Russia
3 Physics Department, Humboldt University Berlin, Newtonstr. 15, 12489 Berlin, Germany
4 Frascati National Laboratory, National Institute for Nuclear Physics, Via Enrico Fermi 40, 00044 Frascati (Rome), Italy
E-mail: amtrnn@gmail.com

Abstract. We present results for the topological susceptibility at nonzero temperature obtained from lattice QCD with four dynamical quark flavours. We apply different smoothing methods, including gradient Wilson flow and over–improved cooling, before calculating the susceptibility. It is shown that the considered smoothing techniques basically agree among each other, and that there are simple scaling relations between flow time and the number of cooling/smearing steps. The topological susceptibility exhibits a surprisingly slow decrease at high temperature.

The non–trivial topological structure of gauge fields and the computation of the topological susceptibility $\chi_{\text{top}}$ is discussed in lattice QCD since long time. Recent considerations (see [1] and references therein) ranging from the restoration of the $U_A(1)$ symmetry at high temperature (or density) to the abundance of cosmic axions [2] are calling for a better knowledge of $\chi_{\text{top}}$ as a function of the temperature and quark masses.

In this study we calculate $\chi_{\text{top}}$ in the temperature range $150 < T < 500$ MeV using lattice configurations generated with $N_f = 2 + 1 + 1$ dynamical Wilson twisted mass fermions at finite temperature [3]. The heavy doublet of $s$ and $c$ quarks has its mass parameters matching the physical $K$ and $D$ meson masses, while the HMC simulations still require light quarks being unnaturally heavy. The configurations are generated at three coupling values $\beta = 1.90, 1.95, \text{and } 2.10$ (for later continuum extrapolations) with the Wilson twisted mass fermionic action tuned to maximal twist, taking benefit from automatic $O(a)$ improvement. Several charged pion masses are considered, but here we restrict ourselves to one value $m_{\pi^\pm} \approx 370$ MeV.

We focus on the comparison between different smoothing methods for lattice gauge fields necessary to get $\chi_{\text{top}}$ by the gluonic method. The smoothing techniques under consideration are the Wilson flow [4], Wilson and over–improved cooling [5], and stout–link smearing [6]. We use the Wilson flow to set two stopping scales determined by

$$t^2\langle E \rangle |_{t=t_0} = 0.3 \quad \text{or} \quad t^2\langle E \rangle |_{t=t_1} = 0.66, \quad E = \frac{1}{2N_\tau N_\delta} \sum_x \text{Tr}[F_{\mu\nu} F^{\mu\nu}(x)], \quad (1)$$

* Presenter
where $F_{\mu\nu}(x)$ is the field strength tensor on the lattice. We match ensemble averages $\langle E \rangle$ obtained in other methods to the values measured with the Wilson flow at $t_0$ and $t_1$ in order to relate these to corresponding numbers of cooling steps. Empirically we find in a good agreement for both the stopping criteria $t_0$ and $t_1$ that $N^\text{Wilson}_{\text{cool}} \simeq 3\tau$ [7], $N^\text{ov.\text{-}imp.\text{-}cool} \simeq 5\tau$ and $N^\text{stout \text{-}smear} \simeq 12\tau$, the latter for our choice $\rho_{\text{smear}} = 0.06$ ($\tau$ defined by $t = a^2\tau$, $a$ is the lattice spacing).

If the diffusion radius $R = \sqrt{\kappa t}$ at $t_0$ or $t_1$ satisfies the condition $R \ll 1/T$ the topological susceptibility $\chi_{\text{top}}(T)$ can be safely determined as

$$\chi_{\text{top}} = \frac{\langle Q_{\text{top}}^2 \rangle}{V}, \quad Q_{\text{top}} = \frac{a^4}{32\pi^2} \sum_x \text{Tr}[F_{\mu\nu}F_{\rho\sigma}(x)], \quad V = a^4N_xN_y^3. \quad (2)$$

The spatial and temporal lattice sizes are $N_x = 24, 32$ and $N_y = 4 \ldots 16$, respectively, and the statistics is varying in the range 200 \ldots 1000 configurations, depending on the respective temperature. One can see from the left panel of Figure 1, where the methods are compared for the finest lattice spacing (at $\beta = 2.10$), that the pairs of Wilson flow and Wilson cooling, as well as over–improved cooling and stout–link smearing, give almost indistinguishable $\chi_{\text{top}}(T)$ values throughout the considered temperature range. In the same manner, the agreement between the two stopping criteria $t_0$ and $t_1$ can be checked independently for each algorithm. The results for $\chi_{\text{top}}$ from Wilson and over–improved cooling agree (within errors) for low temperatures. Starting from approximately $T \gtrsim 300 \text{ MeV}$, the over–improved result for $\chi_{\text{top}}$ turns out somewhat larger than the results for Wilson flow and Wilson cooling as one would expect.

In the right panel of Figure 1, $\chi_{\text{top}}(T)$ is shown as calculated at three different $\beta$ values (different lattice spacings $a$). The curves beyond $\sim 200 \text{ MeV}$ can be reasonably approximated with linear fits, with slopes visibly flattening with the growth of $\beta$ (towards $a \to 0$). Note that the value of crossover temperature obtained from the chiral susceptibility is $T_c = 184(4) \text{ MeV}$. Thus, we come to the preliminary conclusion that for $N_f = 2 + 1 + 1$ and $m_{\pi} \simeq \pm 370 \text{ MeV}$ the topological susceptibility decreases very slowly beyond $T_c$ in the continuum limit, in contrast to the rapid fall–off observed in the quenched approximation ($N_f = 0$), and to the gradual descent taking place in the $N_f = 2$ case [8].

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