

ELLERMAN BOMBS—EVIDENCE FOR MAGNETIC RECONNECTION IN THE LOWER SOLAR ATMOSPHERE

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ABSTRACT

The presence of photospheric magnetic reconnection has long been thought to give rise to short and impulsive events, such as Ellerman bombs (EBs) and Type II spicules. In this article, we combine high-resolution, high-cadence observations from the Interferometric Bidimensional Spectrometer and Rapid Oscillations in the Solar Atmosphere instruments at the Dunn Solar Telescope, National Solar Observatory, New Mexico, with co-aligned *Solar Dynamics Observatory* Atmospheric Imaging Assembly and *Hinode* Solar Optical Telescope (SOT) data to observe small-scale events situated within an active region. These data are then compared with state-of-the-art numerical simulations of the lower atmosphere made using the MURaM code. It is found that brightenings, in both the observations and the simulations, of the wings of the $H\alpha$ line profile, interpreted as EBs, are often spatially correlated with increases in the intensity of the $\text{Fe I } \lambda 6302.5$ line core. Bipolar regions inferred from *Hinode*/SOT magnetic field data show evidence of flux cancellation associated, co-spatially, with these EBs, suggesting that magnetic reconnection could be a driver of these high-energy events. Through the analysis of similar events in the simulated lower atmosphere, we are able to infer that line profiles analogous to the observations occur co-spatially with regions of strong opposite-polarity magnetic flux. These observed events and their simulated counterparts are interpreted as evidence of photospheric magnetic reconnection at scales observable using current observational instrumentation.

Key words: Sun: atmosphere – Sun: chromosphere – Sun: magnetic fields – Sun: photosphere

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1. INTRODUCTION

Ellerman bombs (EBs) were first observed by Ellerman (1917) and are small-scale ($1''$ or less), short-lived (2–15 minutes), impulsive events detected in the lower solar atmosphere (see, e.g., Zachariadis et al. 1987; Georgoulis et al. 2002; Watanabe et al. 2011; Nelson et al. 2013). It has been widely suggested that a link exists between EBs and photospheric vertical magnetic fields; Pariat et al. (2004, 2007) presented co-aligned magnetograms and observations showing the formation of EBs in the plage region trailing an emerging active region (AR). More recently, Nelson et al. (2013) found that strong $H\alpha$ line wing enhancements, identified as small-scale EBs, almost ubiquitously surrounded a complex penumbral structure in an emerging AR and are linked to strong magnetic fields and *G*-band magnetic bright points (MBPs). Due to the link between EBs and strong photospheric magnetic fields, it has often been asserted that EBs arise as a result of magnetic reconnection in the photosphere. Georgoulis et al. (2002) suggested three cartoon topologies that could excite magnetic reconnection in the photosphere, including the partial sinking of a flux tube due to intergranular downflows, flux loops emerging in a *serpentine* manner, and complex unipolar magnetic fields. Within any AR, many examples of cancellation can be observed by magnetogram data, implying the rapid change of magnetic field configuration within the photosphere, potentially consistent with the topologies suggested by Georgoulis et al. (2002). In this article, we address any potential links between EBs and cancellation events around the lead sunspot of a stable AR.

In a recent review, Rutten et al. (2013) suggested that certain brightenings in the wings of the $H\alpha$ line often regarded as EBs could, in fact, be formed due to the influence of high magnetic field concentrations on the $H\alpha$ profile. It was suggested that classical EBs, where energy release leads to increased intensity in the $H\alpha$ line wings, may often be confused with pseudo-EBs, where the line wings of the $H\alpha$ profile outline strong magnetic fields in the lower photosphere. The definitions of these two forms of brightenings advance the work of Watanabe et al. (2011), who suggested that an event must show “flaring,” rapid and small-scale topological variations associated with high energy, to be classified as an EB. Due to the sparseness of high-resolution magnetogram data, however, it has so far proved difficult to identify whether certain topologies lead to the classical EB form, i.e., whether only bipolar regions lead to “flaring” events.

Strong observational evidence supporting a photospheric magnetic reconnection model has, so far, proved difficult to establish. The spatial and temporal resolutions of modern magnetogram data are low compared to hypothesized reconnection events (such as EBs), as well as events within the simulated photosphere (see, e.g., Shelyag et al. 2007), meaning that unequivocal inferences about observed magnetic topologies and evolution are rare. In recent years, however, the possible importance of reconnection in the photosphere has been highlighted by the suggestion that such high-energy events could be driving mass into the chromosphere through the excitation of, for example, Type II spicules (De Pontieu et al. 2007), as well as, potentially, providing energy for heating. This, in turn, has highlighted the need for an analysis of the lower solar atmosphere

in order to infer if any observational signatures of reconnection exist.

Numerical simulations of the lower solar atmosphere have been used to investigate the link between brightening events in the photosphere and bipolar magnetic fields. Isobe et al. (2007), using the Coordinate Astronomical Numerical Softwares (CANS) code, suggested that an emerging flux loop could form a wide range of reconnection events by propagating upward from the photosphere, where EBs can be excited, into the corona, where X-ray jets are formed. A three-dimensional version of the CANS code was used by Archontis & Hood (2009), who found that U- and V-shaped magnetic topologies (similar to those suggested by Georgoulis et al. 2002) were co-spatial to an increase in temperature in the lower solar atmosphere. These temperature enhancements were hypothesized to be comparable to EB events in the $H\alpha$ line wings; however, line profiles were not simulated. Using a semi-empirical method, Fang et al. (2006) found that increases in the VALC temperature in the upper photosphere and lower chromosphere lead to the observed intensity enhancements in the wings of the $H\alpha$ and $\text{Ca II } \lambda 8542$ lines, supporting the conclusions of Archontis & Hood (2009).

Recently, Danilovič (2009) presented an extensive study of simulated magnetic flux cancellation in the solar photosphere using the low-photospheric absorption lines: $\text{Fe I } \lambda 6302.5$, $\text{Fe II } \lambda 5197.58$, and $\text{Fe II } \lambda 4923.92$. One magnetic reconnection event within MPS/University of Chicago Radiative MHD (MURaM) simulations was studied, showing increases in temperature at the inversion line between two opposite-polarity regions. Abnormal Stokes V profiles (where, for the $\text{Fe I } \lambda 6302.5$ line, non-rotationally symmetric lobes were returned), as well as signatures of emission in the Fe II line cores, were found to be co-spatial with such reconnection events. These variations were linked to the process of magnetic field cancellation, formation of current sheets, and local Joule heating. Interestingly, Shelyag et al. (2007) found intensity enhancements and splitting of the Fe I line core co-spatial with magnetic field concentrations in the photosphere, suggesting that signatures of reconnection could be observed in the $6302.5 \text{ \AA}$ line profile.

In this article, we combine observations and simulations to investigate the link between EBs and the underlying magnetic field. High-resolution, high-cadence multi-wavelength observations of a sunspot and the surrounding plasma show the link between the observational signatures of EBs in the wings of the $H\alpha$ line and enhancements in the core of the $\text{Fe I } \lambda 6302.5$ profile. The output of the simulations, made using the MURaM code, is compared with the observational line profiles. The simulated background magnetic fields co-spatial to these line profiles are then studied at high resolution, allowing for inferences about magnetic topologies of the field around EB events, which cannot be detected using current observational techniques. It is found that co-spatial increases in the $\text{Fe I } \lambda 6302.5$ line core and $H\alpha$ line wings occur over small, bipolar regions in both the observations and simulations, presenting the strongest indications to date that EBs are formed by photospheric magnetic reconnection. We present this work as follows: in Section 2, we discuss the observations used in this article; Section 3 outlines the code; Section 4 presents our findings; and conclusions are stated in Section 5.

2. OBSERVATIONS

The data used in this study were obtained with the Interferometric Bidimensional Spectrometer (IBIS) and the Rapid

Oscillations in the Solar Atmosphere (ROSA) instruments at the Dunn Solar Telescope in New Mexico, USA, during a period of good seeing between 14:50:59 and 15:05:34 on 2012 September 30. IBIS ran a 26-image sequence, repeated at a cadence of 5.5 s during the time series, sampling the atmosphere around AR 11579. This sequence included 17 $H\alpha$ wavelength points, taken with unequal step sizes ranging between $\pm 1.0 \text{ \AA}$ from the line center, and nine Fe I points, again taken in unequal steps, between $6302.4 \text{ \AA}$ and $6302.75 \text{ \AA}$. The pixel size of the IBIS instrument is $0''.098$, giving a spatial resolution of approximately $0''.2$.

Two ROSA CCDs were employed during this time series, sampling the granulated photosphere, using the G -band filter, and the chromosphere, as observed by Ca II K (with a band width of $1.2 \text{ \AA}$ centered at $3933.7 \text{ \AA}$). Interpolation time for these two wavelengths was 17 ms for G band and 120 ms for Ca II K , leading to a total, reconstructed cadence of 2.112 s and 8.448 s, respectively. Both wavelengths were observed with diffraction-limited pixel sizes of $0''.069$, giving a spatial resolution of approximately $0''.14$. Within the G -band data, numerous groups of MBPs can be observed, situated firmly within the intergranular lanes. Using these MBPs as a proxy for the vertical magnetic field, accurate co-alignment to the magnetogram data is possible.

We also made use of observations from two satellites, namely, the *Solar Dynamics Observatory* (SDO) and *Hinode*. Multi-wavelength analysis is undertaken using the Atmospheric Imaging Assembly (SDO/AIA). In this study, both the $1600 \text{ \AA}$ and $1700 \text{ \AA}$ wavelengths are exploited as these have previously been shown to display signatures of EBs (see, e.g., Qiu et al. 2000); however, higher in the atmosphere (e.g., $304 \text{ \AA}$, and hotter), no signal has currently been observed (as is shown in Figure 1). We, therefore, limit our analysis to the cooler SDO/AIA lines. Data from the SDO/AIA instrument have physical parameters as follows: pixel sizes of $0''.6$, meaning a spatial resolution of $1''.2$, and a cadence of 24 s (12 s for the hotter lines). Magnetic fields are inferred using Stokes V/I data taken by the Solar Optical Telescope (*Hinode*/SOT) with a pixel size of $0''.155$, a spatial resolution of approximately $0''.31$, and a cadence of around 1 minute. These data are reduced using the *fg_prep.pro* SSWIDL routine, and then compared to a co-aligned slit of spectropolarimetric data, reduced using the standard *sp_prep.pro* routine. For the quiet Sun, a linear conversion coefficient (as demonstrated by Lites & Ichimoto 2013) between Stokes V/I and Gauss was found. In this article, we plot *Hinode*/SOT images with respect to the Stokes V/I parameter. This is because the full field of view (FOV) is not converted to magnetic flux (both the sunspot and penumbra are left in to Stokes V/I form as these create complexities outlined by Lites & Ichimoto (2013) and are surplus to the requirements of this research). A zoomed FOV of the region under investigation is shown in Figure 1. A box highlighting a typical event (in its infancy) studied in this work is overlotted on the images.

In Figure 1, the solar atmosphere is sampled at a number of layers, from the granulated photosphere, observed with the G -band image, to the upper chromosphere, shown with the SDO/AIA $304 \text{ \AA}$ image. A box is overlaid on each wavelength within Figure 1 to highlight a typical event, as studied in this article. The EB event, identified by an increase in brightness, is seen in the center of the box in both the $H\alpha$ line wings (approximately $\pm 0.75 \text{ \AA}$ from the line core); however, it is not observed within the $H\alpha$ line core. The lower photosphere also

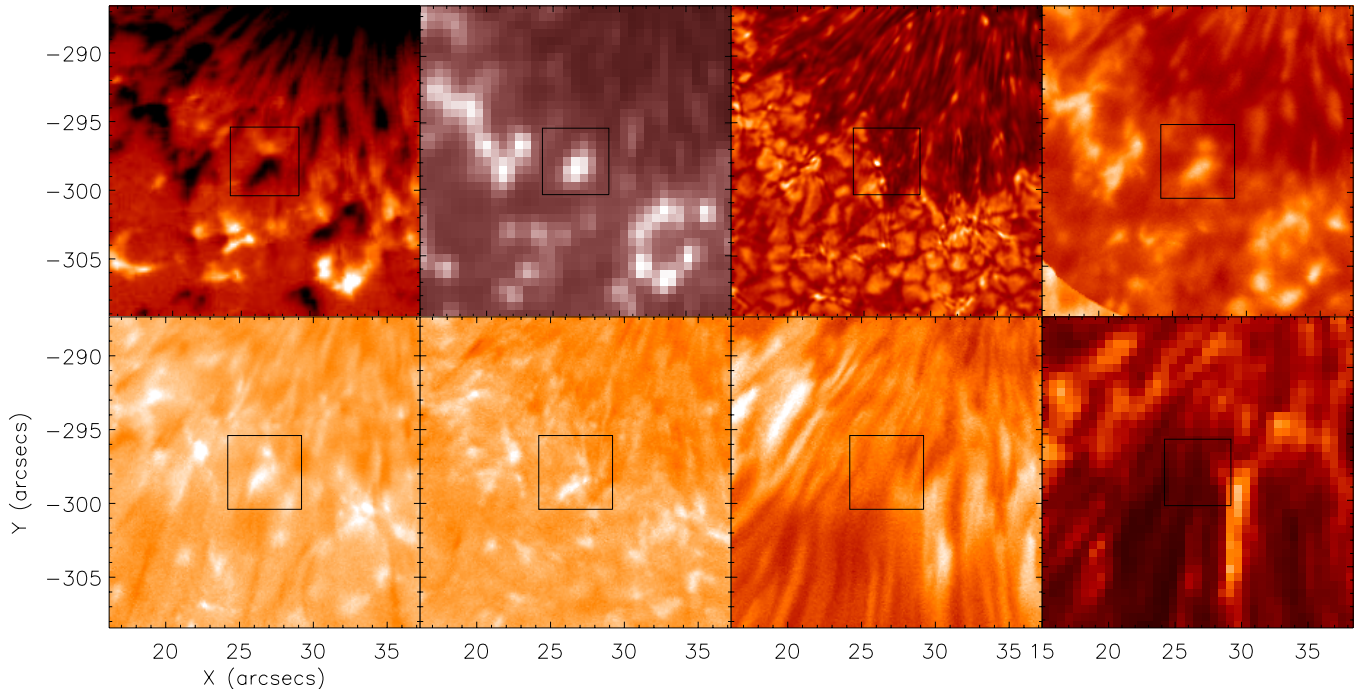


Figure 1. Top row (left to right): co-aligned *Hinode*/SOT Stokes *V/I*, *SDO*/AIA 1700 Å, ROSA *G* band, and IBIS Fe I line core. Bottom row (left to right): IBIS H α blue and red wing images (approximately ± 0.75 Å), IBIS H α line core, and *SDO*/AIA 304 Å. A box including the event analyzed in Figures 3, 6, and 7 is overplotted. This FOV is taken at the temporally closest image for each wavelength to 14:56:35 UT.

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shows evidence of a co-spatial event, with intensity increases being observed in the *SDO*/AIA 1700 Å and ROSA *G*-band images. Within the *G*-band data, a small group of MBPs are seen to approach one another in the time preceding this frame before forming this larger, bright event. Links between EB events and these continua have been widely studied in recent years due to the strong dependence that exists between, particularly, *G* band and the photospheric vertical magnetic field (see, e.g., Qiu et al. 2000; Jess et al. 2010; Nelson et al. 2013 for studies of EBs; Berger & Title 2001 for *G* band as a magnetic field proxy). Many of these researches have been used as evidence of a firm link between EBs and small-scale fields in support of the results obtained for larger, stronger fields, e.g., within a plage region (see, for example, Pariat et al. 2004) or around a complex penumbra (Zachariadis et al. 1987).

In the upper photosphere, the Fe I line core, and lower chromosphere, Ca II K images (not shown here), a co-spatial brightening event is observable; however, no signal of this event is observed in the *SDO*/AIA 304 Å images. These observations are typical of EBs and support the hypotheses of Matsumoto et al. (2008), who found evidence that these events were excited in the upper photosphere. These authors reported flows of opposite direction in the lower photosphere and lower chromosphere, indicating a high-energy driver situated between the two. This has been supported by simulations (for example, Isobe et al. 2007; Archontis & Hood 2009) and shall be discussed further with respect to these MURaM simulations.

3. NUMERICAL SIMULATIONS

The MURaM radiative MHD code (Vögler et al. 2005) has been used to carry out simulations of the photosphere for a bipolar active plage solar region. The code has been thoroughly tested and used for a wide variety of solar and stellar applications

(see, e.g., Shelyag et al. 2007, 2011; Rempel et al. 2009; Cegla et al. 2013). It solves the equations of radiative MHD for the large-eddy approximation on a Cartesian grid and employs a fourth-order central difference scheme for calculating spatial derivatives and a fourth-order Runge–Kutta scheme to advance the solution in time. The numerical scheme is stabilized against numerical instabilities using hyper-diffusion terms included in each of the MHD equations, with one exception: a small constant diffusion is added for the magnetic field, as in Vögler et al. (2005). The code also uses a short-characteristic, non-gray opacity binning radiation treatment scheme to account for radiative energy transport in the convection zone and the photosphere. The solar chemical composition is introduced through the equation of state which takes into account partial ionization for 11 of the most abundant elements in the solar photosphere.

The numerical box is set to have a horizontal extent of 12×12 Mm², while the vertical size of the box is 1.4 Mm. The corresponding numerical grid is 480×480 grid cells in the horizontal direction and 100 grid cells in the vertical direction; thus, the spatial resolution of the model is 25 km and 14 km, respectively. The side boundaries of the numerical domain are periodic. The top boundary is closed for in- and outflows, while it allows horizontal motions of plasma and magnetic field. The dispersion created by such horizontal flows, as well as the exponential decrease in density in the upper layers of the atmosphere, means that the effects of this closed boundary on these simulations are negligible (as was discussed by Beeck et al. 2012). The bottom boundary is transparent for both in- and outflows. The 500 nm continuum formation level is located at a height of approximately 800 km above the bottom boundary.

We start the simulation from a statistically stabilized and well-developed snapshot of non-magnetic photospheric convection.

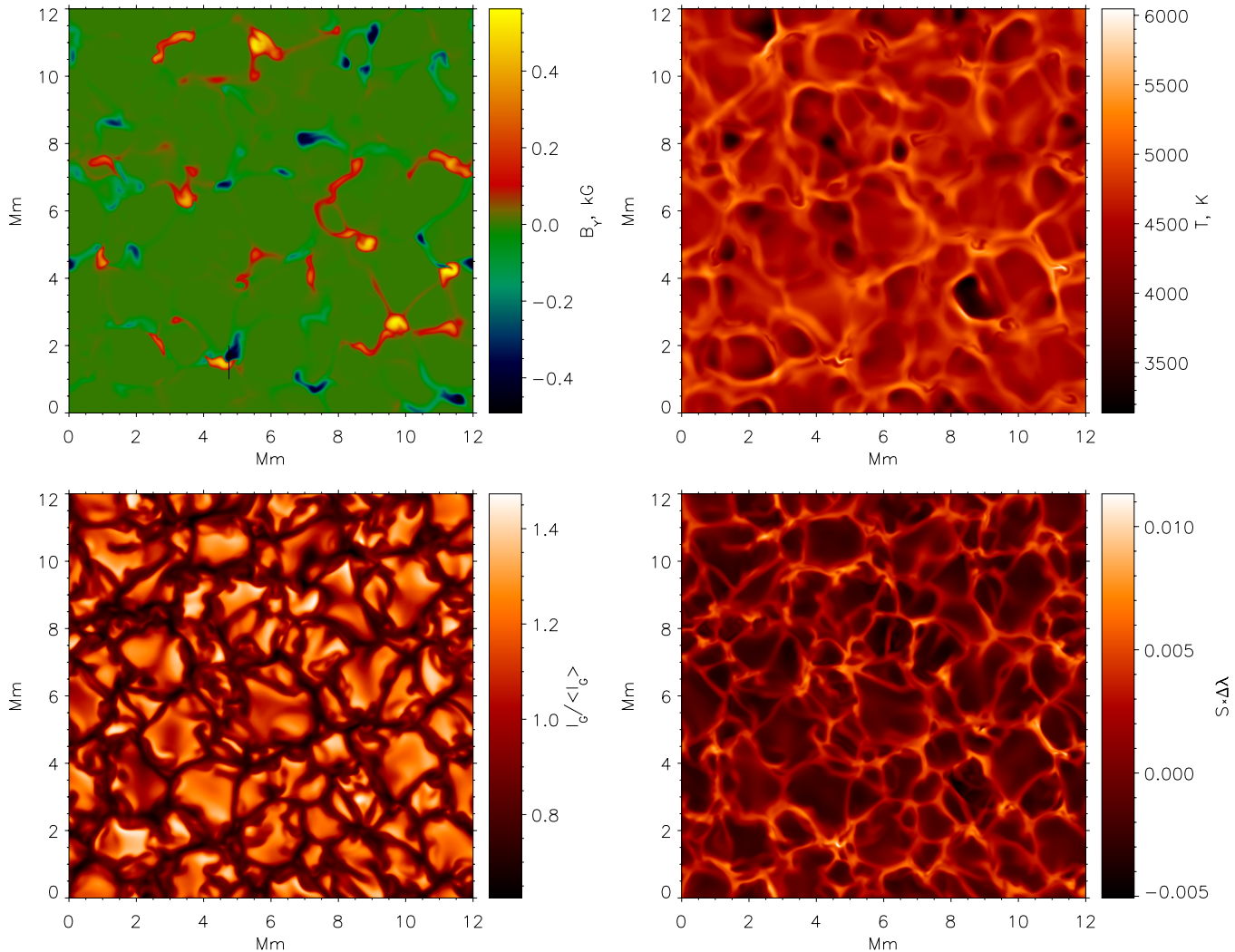


Figure 2. Top row: vertical component of magnetic field (left) and temperature (right) measured at 500 km height above the continuum formation level. Bottom row: normalized G -band intensity (left) and $\text{Fe I } \lambda 6302.5$ line strength, S , multiplied by the line Doppler shift, $\Delta\lambda$ (right). The reconnection event studied in this paper is marked by a line at (4.75, 1–2) Mm in the top left panel.

(A color version of this figure is available in the online journal.)

In a manner similar to that described by both Khomenko et al. (2005) and Danilović et al. (2010), a bipolar, checkerboard magnetic field structure with unsigned magnetic field strength of 200 G is introduced into the domain. The length of the individual squares of constant magnetic field is chosen to be 2 Mm, making a 6×6 box of opposite magnetic field regions. This structure is in agreement with the periodic side boundary conditions.

After the magnetic field was introduced into the domain, we let the simulated photosphere evolve for approximately 1.5 hr. During the first few minutes of the simulation, large parts of the magnetic field were advected into the intergranular lanes, while some magnetic flux was canceled due to the opposite-polarity initial configuration (see, e.g., Cameron et al. 2007). Magnetic field concentrations of opposite polarity with unsigned strength of up to 1.6 kG at the continuum formation level were subsequently formed. Since the turbulent magneto-convection process is effectively random, these intergranular opposite-polarity magnetic field concentrations move and sometimes come closer together, reconnect, and cancel out. One typical example of such an event was selected for further investigation.

Snapshots from our simulations are shown in Figure 2. The reconnection region studied in this paper is marked by a line in the top-left panel. The vertical components of magnetic field (top-left panel) and temperature (top-right panel) measured in the photosphere 500 km above the continuum formation level are plotted and show sharp and localized temperature enhancements in the regions where opposite-polarity magnetic field concentrations come close to each other and reconnect, leading to localized flows (as observed by Matsumoto et al. 2008). Reconnection events within these simulations appear to be of the Sweet–Parker type, whereby a current sheet is formed between two regions of opposite-polarity field. Although basic implementation of this model provides slow, non-physical reconnection rates, the inclusion of flux pileups and strong granular flows around the reconnection point can speed up the process (as has recently been shown by Litvinenko et al. 2007) to be comparable to physical phenomena, as observed here. Notably, the G -band intensity (bottom-left panel) does not show any peculiarities in the regions of magnetic reconnection, clearly demonstrating that the significant (compared to the background plasma) energy release from magnetic field reconnection occurs

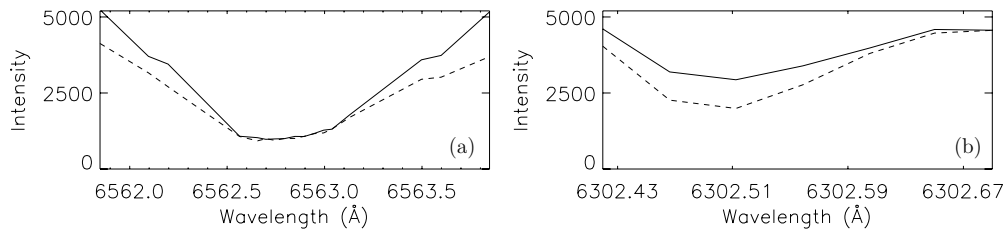


Figure 3. (a) $H\alpha$ line profile for two spatial positions within the box in Figure 1. The solid and dashed lines correspond to the EB event and quiet Sun, respectively. (b) Co-spatial Fe I line profiles.

higher than the location of the continuum radiation layer. The Fe I intensity (bottom-right panel) does show an increase in intensity co-spatial to the magnetic reconnection event and, therefore, with the temperature increase.

We performed spectropolarimetric diagnostics of the simulated event with the Fe I $\lambda 6302.5$ magnetically sensitive absorption line and the $H\alpha$ line. The Fe I line is known to form in the low photosphere, where, in general, the local thermodynamic equilibrium approximation holds; however, as reported by Shchukina & Trujillo Bueno (2001), application of an LTE approximation to neutral iron lines leads to a different temperature sensitivity when compared to a precise, but computationally difficult, non-LTE approach. Due to expected temperature effects during the reconnection process in the solar photosphere, we employed a full non-LTE line profile synthesis code RH (see Uitenbroek 2001 for information) to perform radiative diagnostics with the Fe I line.

The $H\alpha$ line is formed over the range of heights starting from the deep photosphere up to the chromosphere or transition region (see, e.g., Rutten 2008; Leenaarts et al. 2012). The wings of the $H\alpha$ line profile are formed in the photosphere and clearly show the presence of various structures, from large sunspot umbrae and penumbrae to smaller rapid blue excursions (see, e.g., Langangen et al. 2008; Rouppe van der Voort et al. 2009). What is not clear, however, is whether EBs, observed in the $H\alpha$ wings, are formed in the photosphere as has been widely hypothesized (by, e.g., Georgoulis et al. 2002; Watanabe et al. 2011; Nelson et al. 2013). In this work, using numerical modeling of a solar atmosphere that includes only the photosphere, we aim to model and analyze the behavior of the $H\alpha$ line wings, thereby proving the formation height of EBs. The same radiative diagnostics code RH was used to calculate the $H\alpha$ line profiles.

Since full three-dimensional calculations of the radiative field are still a computationally challenging task, we used a simplified approach of computing the emergent, wavelength-dependent intensities along the vertical rays in the selected regions of the simulation domain. Then, we degraded the spectral resolution of the obtained Stokes I profiles in order to directly compare the synthetic images with the observations.

4. RESULTS AND DISCUSSION

We identify a number of large EB events within these data, described in Section 2, by analyzing $H\alpha$ line profiles for intensity enhancements and assessing for any evidence of “flaring,” or an explosive nature, during bright events. Events that show significant increases in line wing intensity and no signature in the line core are studied (one representative example of such an event is shown in imaging form within the box in Figure 1 and line profile form by Figure 3). Small, short-lived intensity enhancement events (exhibiting line profiles analogous to EBs), with circular diameters less than $0''.5$ and lifetimes often below

2 minutes, are common in this data set (as predicted by Nelson et al. 2013); however, due to the tendency of line scan data to change quality rapidly, a large-scale statistical study of these short-lived events could provide spurious results. We, therefore, limit our analysis to larger EB events, which show sustained brightening and “flaring” as discussed by Watanabe et al. (2011). It should be noted that no events appear to be brighter than 150% of the background intensity of these data (possibly due to effects of seeing) and, hence, 140% of the background emission is used as a guide for identification of EBs.

In Figure 3, we plot the $H\alpha$ and Fe I line profiles for both the EB and a quiet-Sun region within the box studied in Figure 1. The $H\alpha$ line wings show significant intensity enhancements compared to the quiet Sun, whereas no signature within the line core is observed. This agrees with the definition of EBs put forward by Ellerman (1917), who suggested that these events were “a very brilliant and very narrow band extending four or five Å on either side of the line, but not crossing it.” Co-spatially, the Fe I line core shows increases in intensity. This line profile pairing (increases in the $H\alpha$ wings and Fe I core) is typical of the response of the Fe I $\lambda 6302.5$ line to the studied EB events in these observations, suggesting a possible link between the formation of the $H\alpha$ and Fe I lines. We present and analyze one event with these characteristics in more detail in Figures 6 and 7.

Similar, co-spatial behavior in the $H\alpha$ and Fe I $\lambda 6302.5$ line profiles is observed in numerical simulations at regions of magnetic field reconnection in the solar photosphere. In Figure 4, we plot the magnetic field geometry (panel 1), the temperature structure (panel 2), and the corresponding Fe I $\lambda 6302.5$ (panel 3) and $H\alpha$ (panel 4) line profiles for each vertical cut across the magnetic reconnection region identified in Figure 2. The magnetic field lines clearly outline the inversion line between the regions of the opposite magnetic field polarities (the color of the field lines represents their direction, which is white and black for upward- and downward-directed magnetic field, respectively). The temperature increase, co-spatial with the inversion line, corresponds to the Joule heating and energy release during reconnection (Danilović 2009). The Fe I line profile shows a prominent “gap” with strongly reduced absorption in the increased temperature region. Co-spatially, the $H\alpha$ line profile indicates a strong increase in intensity in both its wings, similar to an observed EB. These results support the semi-empirical assertions of Fang et al. (2006), who found that heightened emissions in the $H\alpha$ line wings were most likely a result of increased temperature within a magnetic inversion line. The above follows on from earlier radiative transfer modeling of hydrogen spectrum in the dMe stars by Houdebine & Doyle (1994). Also, a (stronger) narrowing of the profile, as in the observations, is simulated. This narrowing coincides with the uppermost part of the reconnection layer. The narrowing is formed due to the heightened electron density at the flare

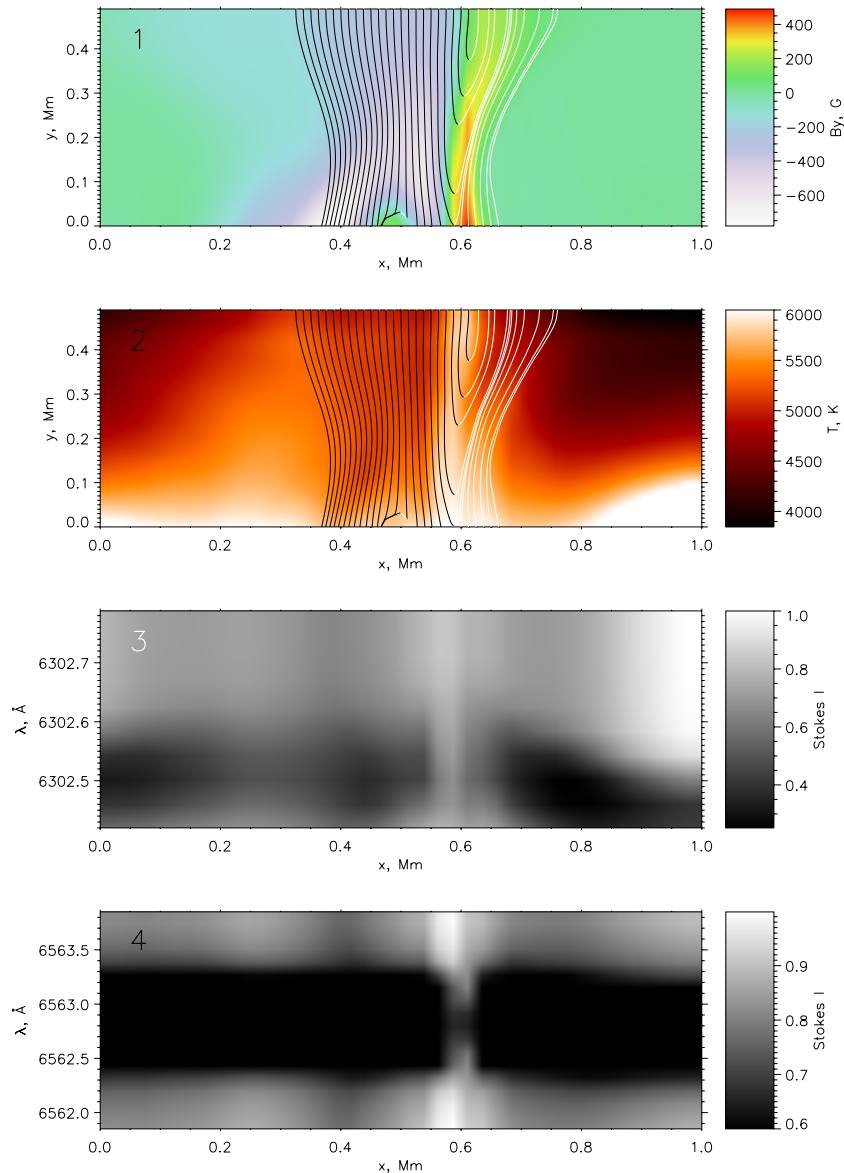


Figure 4. Reconnection event in the simulated photosphere. (1) The slice of the simulations box indicated by the slit in Figure 2 showing the vertical component of magnetic field in the reconnection region with height (y -axis where 0 Mm corresponds to the continuum layer), with the magnetic field lines overplotted. Black lines correspond to the negative (downward) direction of magnetic field, and white lines represent positive (upward) direction. The field lines originate in the points where the total magnetic field strength is greater than 150 G. (2) Temperature image with the field lines overplotted. The reconnection process is clearly seen as a strong temperature enhancement in the region of U-shaped magnetic field lines. (3) Fe I $\lambda 6302.5$ line profile with a “gap” in the reconnection region ($x \sim 0.6$ Mm), indicating an increase in intensity of the line core. (4) H α line profile shows brightenings in both wings in the reconnection region, which is identified as an EB.

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reconnection site increasing the atmospheric mass load, thus creating a deeper absorption core (Cram & Mullan 1979).

The time-dependent nature of reconnection, the evolution of the magnetic field geometry in the simulations, and the short-lived modification of the H α line profile in observed EBs suggest that further investigation into the temporal properties of these events is necessary. Namely, the expansion of the magnetic field concentrations toward the upper layers of the solar atmosphere due to decreasing gas pressure will cause the reconnection to happen earlier in the upper-simulated photosphere than in the lower regions. Thus, the thermal effect of reconnection will first be seen in the parts of the H α line profile sensitive to the higher solar atmosphere, followed by intensity increases in the H α wings and reduced opacity in the Fe I $\lambda 6302.5$ line

core. This process is demonstrated in Figure 5, where the temporal evolution of the Fe I (top panel) and H α (bottom panel) line profiles during the reconnection event is plotted. It is easy to see that the H α line core narrowing (due to increased electron density) is followed by the Fe I “gap” and the rise in the H α wings, which occur co-temporally.

It should be noted that the rise in the Fe I wing intensity following the EB event in Figure 5 does not necessarily correspond to the same driver. At around 5 minutes, there is a significant increase in the Fe I red wing intensity that is simultaneous with a very mild increase in the H α wings, while no strong rise in the Fe I core intensity is observed. This behavior can be linked to an increase in the background photospheric continuum intensity and reduced opacity in the

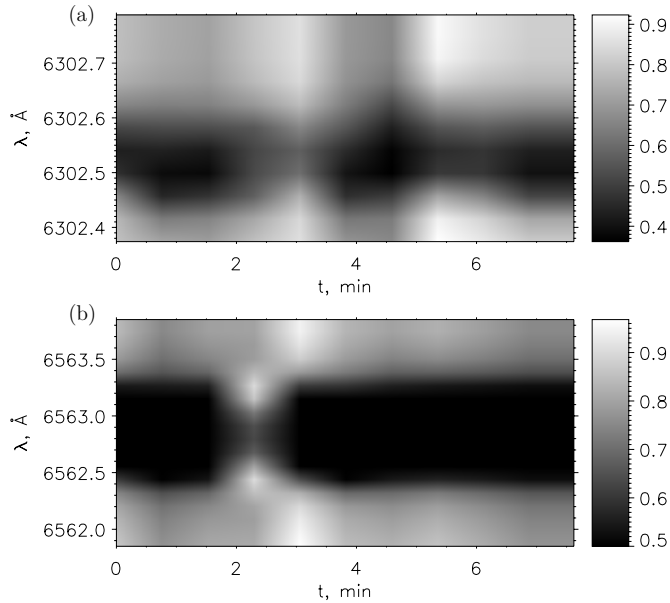


Figure 5. Temporal evolution of the simulated EB. (a) Fe I $\lambda 6302.5$ line profile. (b) H α line. Between 2.5 and 4.5 minutes, the Fe I line shows a “gap” corresponding to the reconnection event and H α shows line wing increases. Note, there is also a slight brightening in H α core about 1 minute before the EB event.

photospheric magnetic field concentration, which does not experience reconnection (see, e.g., Shelyag et al. 2007).

The temporal evolution of bright events in these observations was analyzed by identifying possible EB events within the H α line wings and then producing a sequence of images. These sequences included the H α wings, the Fe I line core, and *Hinode*/SOT Stokes V/I images. One excellent example is plotted in Figure 6. The H α line wings (blue and red for the top and second rows, respectively) and the Fe I line core (third row) show strong brightenings over the bipolar region (in the *Hinode*/SOT Stokes V/I data plotted in the bottom row). The EB shows increased brightness and area over the lifetime of the event, which is identified as “flaring” (as discussed by Watanabe et al. 2011). Rapid topology changes are easily observed through the sequence, implying a dynamic nature for this event. We find that this behavior is common in EBs found over bipolar regions and could be used as an indicator that these events are classical EBs rather than the possibly more common pseudo-EBs (which purely mimic the magnetic field in the photosphere; see, e.g., Rutten et al. 2013). Conducting an analysis of Dopplergrams (made using the [blue–red]/[blue+red+2] formula; see Madjarska et al. 2009), we find a common association between EBs and upward shifts in the H α line, analogous to those observed by Matsumoto et al. (2008); however, due to the small number of data points in the wings, we are unable to estimate accurately the velocities of these events.

Light curves for this event are plotted in Figure 7. These light curves were made by focusing the image used in Figure 6

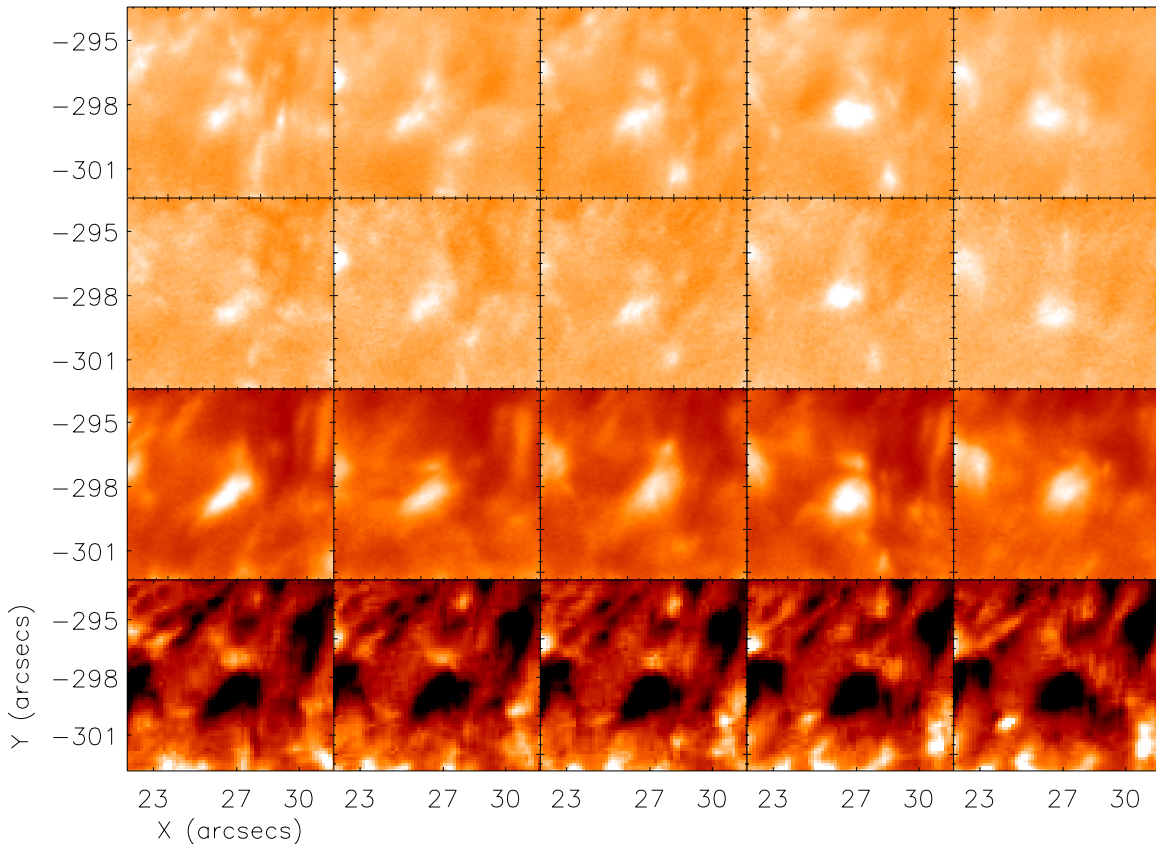


Figure 6. Zoomed FOV of a bipolar region and a corresponding EB event in the H α blue wing position (top row), H α red wing position (second row), Fe I line core (third row), and SOT Stokes V/I (bottom row) over time. Each column is the temporally closest frame for each wavelength to 14:52:36, 14:55:38, 14:58:39, 15:01:39, and 15:04:45, respectively. Light curves for the full observational period are plotted in Figure 7.

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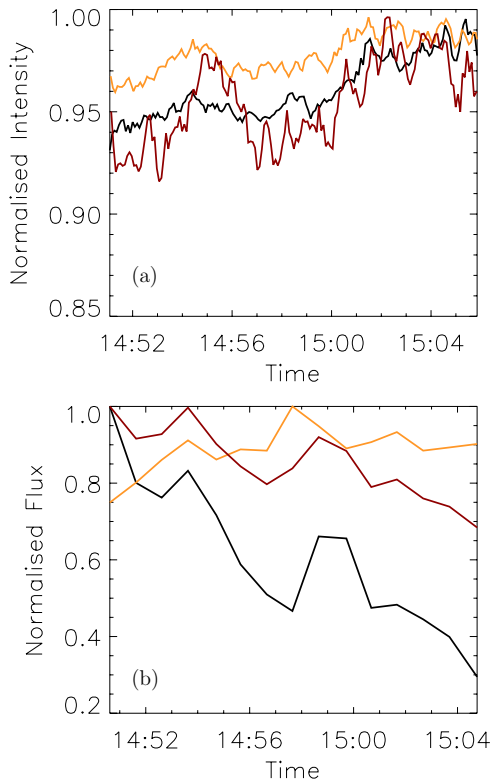


Figure 7. (a) Evolution of IBIS data over time. Black is for the H α blue wing, orange is for the H α red wing, and maroon is for the Fe I line core. (b) The evolution of the magnetic field over time. Black is for negative flux, orange for positive flux, and maroon for total flux. A decrease of flux to around 68% of the original strength is observed.

(A color version of this figure is available in the online journal.)

onto the opposite-polarity region such that no new, strong flux appears within the FOV in the lifetime of these observations. The total positive flux was then taken by summing all positive values within the box (vice versa for the negative polarity) for each frame before plotting the normalized total flux over time. In Figure 7(a) we plot the smoothed H α line wing (black and maroon for ± 0.75 Å, respectively) and the Fe I line core (orange) intensities. In Figure 7(b) the negative polarity (black), positive polarity (orange), and total flux (maroon) are shown. Over the period of the observations, there is a significant decrease in the flux of this small bipolar region, dropping to around 68% of the original flux, and only 30% of the negative-polarity flux. This decrease is easily observed through the sequence of images presented in Figure 6. Within a 15 minute period, a decrease of this order due to magnetic reconnection would suggest that any corresponding energy release must be strong and could produce large temperature increases in the surrounding atmosphere.

The IBIS data (in Figure 7(a)) reach a maximum intensity for each of these lines at approximately 15:02, maintaining a peak for around 2–4 minutes. This peak is identified as “flaring” within this EB, associated with both increased intensity and area. The short-lived peak is more evident within the Fe I line core, where an 8% increase in intensity is observed within around 1 minute. Co-temporally, a rapid decrease in negative (and total) polarity flux in Figure 7(b), which continues until the end of these data, is observed. A second, smaller peak is observable at approximately 14:54 and corresponds with a large drop in both negative polarity and total flux within this FOV. Although this co-temporal brightening and flux cancellation occur for

several examples within this data set, the cadence of the *Hinode*/SOT magnetograms is too low to make further inferences about whether this is a correlation or a coincidence; therefore, we shall leave this aside and note that it would be worthy of future study when sufficiently high cadence magnetograms become available.

H α wing brightening events co-spatial to unipolar regions, such as the strong field at (33, −302) in Figure 1, were also studied within these data. The average area of brightenings over unipolar fields appeared to be larger than bipolar EBs. Smaller, shorter-lived intensity variations were also common, implying that rapid, high-energy releases may not be leading to these brightenings and, hence, that the increase in intensity in the H α line wings and Fe I line core may be due to the strength of the magnetic field (analogous to the pseudo-EBs discussed by Rutten et al. 2013 and the reduced opacity simulated by Shelyag et al. 2007, respectively). Within such events, the shape of the intensity enhancement is often different between the imaged H α wings (which is not true of the bipolar EBs, such as the event presented in Figure 6). We find that brightenings occurring over unipolar regions still show significant intensity increases (often above thresholds currently applied within automated EB tracking, such as 130% or 150% of the background intensity) within the H α line wings; however, the line profiles frequently show excess intensity in one wing over the other.

The event at (33, −302) in Figure 1, for example, displays co-spatial intensity increases for short periods during these observations and would, using thresholds alone, be classified as an EB for these times; however, for the majority of these observations, the red wing exhibits traits associated with the network, whereas the blue wing is extremely bright (as can be seen in Figure 1). We suggest that comparable temporal and morphological evolutions between the H α line wings are an important factor in EB identification (similar to that shown in Figure 6) and that any events that show non-comparable evolutions between the wings are, in fact, pseudo-EBs.

5. CONCLUDING REMARKS

In this article, we have used multi-instrument, multi-wavelength observations and numerical simulations to investigate the formation mechanism of EBs, small-scale brightening events in the H α line wings, commonly found to occur in emerging ARs. Numerous possible EB events were found within these data by identifying instances of H α line wing intensity increases. The evolution of these events was then studied in order to identify whether they showed instances of “flaring.”

In order to investigate the link between EBs and magnetic fields, small opposite-polarity flux regions, observed to be spatially isolated from other strong fields, were analyzed. It was found that rapid changes in both magnetic flux polarity (apparent conversion from positive to negative, or vice versa) and magnetic flux strength (a total decrease in flux) were present during each of the studied brightening events. The evolution of a large, “flaring” EB event and the corresponding flux region is shown over the 15 minute observational period studied in this article, finding a quite remarkable decrease in total flux to approximately 68% of the initial strength. This rapid cancellation of flux, as well as the associated brightenings in this multi-wavelength analysis, implies that a release of magnetic energy into the surrounding plasma by a small-scale magnetic reconnection event is a more likely reason for the reduction in flux than a sinking of tubes back into the convective layer. Unfortunately, even relatively

high-cadence, high-resolution magnetograms, such as those studied here, are unable to detect the difference between these two scenarios.

A model box was constructed to simulate the solar photosphere using the MURaM code of radiative magneto-convection. Within the model box, the line profiles from small bipolar regions, similar to the representative example presented within this article, exhibited similar traits to EBs forming over bipolar regions. Further study of the evolution of these small-scale events over time in the simulations was undertaken, showing temperature increases co-spatially with magnetic inversion lines at bipolar regions. Interestingly, the relationship noted in the observations between the $H\alpha$ and Fe I line profiles also appeared in the simulations, which showed increased intensity in the line wings and line core, respectively. The reconnection events analyzed in this article appear to be of Sweet–Parker type; however, a more specialized study is required to fully understand the mechanism that is occurring.

The influence of the magnetic field on the Fe I lines has been discussed in previous researches (see, e.g., Shelyag et al. 2007; Danilović 2009). Shelyag et al. (2007) suggested that regions of strong flux, especially close to magnetic inversion lines, where opposite polarities converge, could lead to a heightened core, as well as splitting (where two intensity minima are observed within the line profile), whereas Danilović (2009) found that abnormal (non-rotationally symmetric) Fe I $\lambda 6302.5$ Stokes V profiles were co-spatial with both strong magnetic fields and magnetic reconnection in the photosphere. It has, therefore, been suggested that the formation of the Fe I $\lambda 6302.5$ line profile is intrinsically influenced by the magnetic field. A combination of these results could be used to intuitively expect co-aligned EB and Fe I brightenings as were found here; however, as the observations used in this article could not be used as a diagnostic tool for Stokes V and, due to both the spatial and spectral resolutions, did not find evidence of splitting within the Fe I line core, we suggest that further research be carried out when sufficiently high resolution data are available.

Overall, the observed and simulated photospheres analyzed in this article exhibit analogous line profiles co-spatial to bipolar regions, namely, brightenings within the $H\alpha$ line wings and an increase in intensity in the Fe I $\lambda 6302.5$ line core. We suggest that this analysis has presented the clearest evidence to date that the sub-class of brightening events known as EBs in the $H\alpha$ line wings are formed by magnetic reconnection in the solar photosphere. However, it should be noted that improved temporal, spatial, and spectral resolution when inferring both photospheric imaging lines and magnetic field is required before observations are capable of standing alone in identifying magnetic reconnection.

Finally, it should be noted that this article does not suggest that all identified brightening events are formed through magnetic reconnection in the photosphere. Several events situated over unipolar regions within the observations that form the required $H\alpha$ profiles to be classified as EBs (including significant brightness enhancements in the line wings) do not possess any evidence of “flaring,” which appears to be essential in identifying the more dynamic, higher-energy EBs. We suggest that to make any inferences on future data sets, the authors must have high-resolution, high-cadence data in a number of wavelengths, as presented here. Further research must be undertaken before conclusions about the wider range of brightening events in the $H\alpha$ line wings can be drawn.

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