# REALLY COOL STARS AND THE STAR FORMATION HISTORY AT THE GALACTIC CENTER

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### ABSTRACT

We present  $\lambda/\Delta\lambda = 550-1200$  near-infrared H and K spectra for a magnitude-limited sample of 79 asymptotic giant branch and cool supergiant stars in the central  $\approx 5$  pc (diameter) of the Galaxy. We use a set of similar spectra obtained for solar neighborhood stars with known  $T_{eff}$  and  $M_{bol}$  that is in the same range as the Galactic center (GC) sample to derive  $T_{eff}$  and  $M_{bol}$  for the GC sample. We then construct the H-R diagram for the GC sample. Using an automated maximum likelihood routine, we derive a coarse star formation history of the GC. We find that (1) roughly 75% of the stars formed in the central few parsecs are older than 5 Gyr; (2) the star formation rate (SFR) is variable over time, with a roughly 4 times higher SFR in the last 100 Myr compared to the average SFR; (3) our model can match dynamical limits on the total mass of stars formed only by limiting the initial mass function to masses above 0.7  $M_{\odot}$  (this could be a signature of mass segregation or of the bias toward massive star formation from the unique star formation conditions in the GC); (4) blue supergiants account for 12% of the total sample observed, and the ratio of red to blue supergiants is roughly 1.5; and (5) models with isochrones with [Fe/H] = 0.0 over all ages fit the stars in our H-R diagram better than models with lower [Fe/H] in the oldest age bins, consistent with the finding of Ramírez et al. that stars with ages between 10 Myr and 1 Gyr have solar [Fe/H].

*Subject headings:* Galaxy: center — stars: AGB and post-AGB — stars: late-type — supergiants *On-line material:* color figures, machine-readable tables

# 1. INTRODUCTION

The properties of the stellar population at the Galactic center (GC) suggest that the nucleus is distinct from the other main structural components of the Galaxy (the Galactic disk, bulge, and halo), although each of these components may contribute to the integrated population. We would like to distinguish between extensions of these populations and a unique GC population that has formed and evolved there. OH/IR stars distributed in an inner disk between  $\sim 1$  and 100 pc (Lindqvist, Habing, & Winnberg 1992) show higher rotational velocities than expected for a "hot" bulge component, suggesting a disklike population. Near-infrared surface brightness measurements indicate that the bright nucleus joins the bulge discontinuously at a radius of about 150 pc (see the discussion by Kent 1992), which relies in part on the 4  $\mu$ m minor-axis surface brightness profile presented by Little & Price (1985). Recent work on detailed abundance determinations in the central 60 pc

(Carr, Sellgren, & Balachandran 2000; Ramírez et al. 2000) also reveal differences between the Galaxy's nucleus and bulge components. Ramírez et al. (2000) find a narrow distribution in [Fe/H] in the GC peaked around the solar value, while the bulge has a very broad distribution in [Fe/H] with a mean less than the solar value (McWilliam & Rich 1994; Sadler, Rich, & Terndrup 1996).

Stars as young as  $\lesssim 5$  Myr are now known to exist in the central parsec. Very recent star formation was clearly established by Forrest et al. (1987) and Allen, Hyland, & Hillier (1990) with the discovery of a bright, evolved, and massive emission-line star (the "AF" star). Krabbe et al. (1991) further showed that a significant component of the strong recombination lines of H and He seen toward the GC arises in spatially compact sources, particularly the "IRS 16" cluster of massive stars (Najarro et al. 1994, 1997; Libonate et al. 1995; Blum et al. 1995a, 1995b; Krabbe et al. 1995; Tamblyn et al. 1996). Krabbe et al. (1995) have modeled the IRS 16 cluster as the evolved descendants of the most massive stars ( $\sim 100 M_{\odot}$ ) belonging to a  $\lesssim 7$  Myr old burst.

In their review of the global phenomena ongoing in the GC region, Morris & Serabyn (1996) described the properties of the "central molecular zone" (CMZ). The CMZ is a "disk" of enhanced molecular density about 200 pc in

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radius centered on the GC. The gas is confined to a region near the plane of the Galaxy, but with significant noncircular motions. The distribution and presence of molecular gas in the CMZ may in large part be due to the effects of the inner Galactic stellar bar (Liszt & Burton 1980; Mulder & Liem 1986; Binney et al. 1991; Blitz & Spergel 1991; Weiland et al. 1994; Dwek et al. 1995; Stanek et al. 1994). The material in the CMZ is fueling current star formation on this large scale at a rate of about 0.5  $M_{\odot}$  yr<sup>-1</sup> (Güsten 1989), but it may also be the ultimate source of material that is processed into stars within a few parsecs of the GC (Morris & Serabyn 1996). If so, angular momentum losses must funnel the gas down to the circumnuclear disk (CND) at radii between  $\sim 2$  and 8 pc (see the extensive reviews by Genzel, Hollenbach, & Townes 1994 and Morris & Serabyn 1996). This molecular structure is probably not a long-lived one, but rather periodically forms and supplies the GC with star-forming material through instabilities that cause material to fall from its inner radius into the central parsec (Sanders 1999); at present the CND may be accreting about  $0.5 \times 10^{-2} M_{\odot} \text{ yr}^{-1}$  (Güsten et al. 1987; Jackson et al. 1993).

In this paper we continue the exploration of the stellar content of the central few parsecs of the Milky Way begun by Blum, DePoy, & Sellgren (1996a, hereafter BDS96) and Blum, Sellgren, & DePoy (1996b, hereafter BSD96). Using J, H, and K photometry, BSD96 identified a bright component to the dereddened K-band luminosity function relative to the Galactic bulge population seen toward Baade's window (BW), which is predominantly old ( $\geq 10$  Gyr; Terndrup 1988; Lee 1992; Holtzman et al. 1993). Specifically, BSD96 compared K-band counts in the central few parsecs with those in the BW presented by Tiede, Frogel, & Terndrup (1995). BSD96 presented a small sample of nearinfrared spectra, which they used to begin a detailed investigation of the properties of the cool stellar population in the GC, including the ages of individual stars that trace multiple epochs of star formation there. This work, in turn, is built on earlier work, most notably that of Lebofsky, Rieke, & Tokunaga (1982b), who investigated recent star formation in the GC using the luminous and young M-type supergiants that they identified there.

Our goal is to determine  $T_{\rm eff}$  and  $M_{\rm bol}$  for a magnitudelimited sample of GC stars, using the two-dimensional classification provided by the measured strengths of the CO and H<sub>2</sub>O absorption features present in modest resolution *K*and *H*-band spectra. The technique, described in § 3, is calibrated using a sample of comparison stars with known  $T_{\rm eff}$  and  $M_{\rm bol}$ , selected from the literature to match the  $T_{\rm eff}$ and  $M_{\rm bol}$  of the GC stars. After  $T_{\rm eff}$  and  $M_{\rm bol}$  are determined for the GC stars, we place the GC stars in the H-R diagram (§ 4.1) and use this to constrain the star formation history (SFH) within the central few parsecs of our Galaxy. The SFH calculation is described in § 4.2 and discussed in § 5. A brief summary is given in § 6.

#### 2. OBSERVATIONS AND DATA REDUCTION

Spectroscopic observations of the comparison and GC stars were made using the facility Infrared Spectrometer (IRS) and OSIRIS<sup>2</sup> spectrometer mounted on the Cerro

Tololo Interamerican Observatory (CTIO) 4 m Blanco Telescope over several runs beginning in 1997 and ending in 2000 (see Tables 1 and 2). In addition, comparison stars were observed at the Michigan-Dartmouth-MIT (MDM) 2.4 m telescope on Kitt Peak using the Ohio State University MOSAIC infrared camera/spectrometer (Table 1). The IRS, OSIRIS, and MOSAIC are described by DePoy et al. (1990, 1993) and Pogge et al. (1998), respectively. The IRS employed a  $0.7 \times 12.5$  slit, OSIRIS employed a  $1.2 \times 30^{\prime\prime}$ slit, and MOSAIC used a  $0.6 \times 150^{\circ}$  slit. The detector pixel scales are 0.32, 0.40, and 0.30 pixel<sup>-1</sup> for the IRS, OSIRIS, and MOSAIC, respectively. The IRS and OSIRIS were used in cross-dispersed mode, giving essentially full coverage of the J, H, and K bands. MOSAIC was used in JHK grism mode (1.22–2.29  $\mu$ m, where the 2.29  $\mu$ m cutoff is due to the JHK blocking/order sorting filter). For the MOSAIC spectra, an extra, independent K-band segment was obtained for each star covering the red portion of the K band. This was accomplished using the same setup, but with a K filter instead of the JHK blocker. The spectral resolutions are approximately  $\lambda/\Delta\lambda = 1200$ , 560, and 750 for OSIRIS, the IRS, and MOSAIC, respectively.

Observing conditions varied over the course of different observing runs. Data were obtained in photometric and nonphotometric conditions. The *K*-band seeing at the CTIO 4 m telescope was typically between 0".5 and 1". At the MDM 2.4 m telescope no effort was made to keep the bright comparison stars in good focus, and in fact, sometimes the telescope was intentionally defocused (see below).

All basic data reduction was accomplished using IRAF.<sup>3</sup> Each spectrum was flat-fielded using dome flat-field images and then sky-subtracted using a median combined image formed from the data themselves or from a set of independent sky frames obtained off the source (for the GC stars, typically on dark clouds  $\sim 30''-90''$  away). Nearby sky apertures ( $\sim 1''-2''$  on either side of the object) were defined on the long-slit images and used to correct for over- and undersubtraction of the night-sky OH lines and the unresolved background light in the case of the GC stars. For the comparison stars, the situation varied depending on the brightness of the star. Some of the stars had to be defocused and/ or placed on the edge of the slit in order not to saturate the detector. In a number of cases at the 4 m Blanco Telescope, the mirror covers were partially closed. These procedures typically produced considerable wings to the point-spread function. For the IRS, the slit was not long enough in such cases to provide blank sky, so sky frames were obtained off the source, some 5'' to 10'' away. These special procedures do not affect the spectral resolution as confirmed by comparing the night-sky line widths with similar spectra taken under normal conditions.

Following sky subtraction, the object spectra were extracted from the long-slit images by summing the dispersed light over  $\pm 3$ -5 spatial pixels (depending on seeing and source crowding in the GC) and then divided by the spectrum of an O, B, or A star to correct for the telluric absorption. For the case of the bright comparison stars that were defocused, the extraction apertures were up to

<sup>&</sup>lt;sup>2</sup> OSIRIS is a collaborative project between the Ohio State University and CTIO. OSIRIS was developed through NSF grants AST 90-16112 and AST 92-18449.

<sup>&</sup>lt;sup>3</sup> IRAF is distributed by the National Optical Astronomy Observatory, which is operated by the Association of Universities for Research in Astronomy, Inc., under cooperative agreement with the National Science Foundation.

TABLE 1
COMPARISON STAR TABLE OF OBSERVATIONS

HR	HD	Other	T <sub>eff</sub> (K)	$M_{ m bol}$	Spectral Type	$T_{ m eff}$ Reference <sup>a</sup>	$M_{ m bol}$ Reference <sup>a</sup>	Observation Date	Instrument <sup>1</sup>
					Giants				
HR 6418	HD 156283	$\pi$ Her	4100	-3.34	K3 II	Dy98	Hip+La91	1999 May 5	MOSAIC
HR 4299	HD 095578	61 Leo	3700	-3.04	M0 III	Fe90	Hip+La91	1999 Mar 3	MOSAIC
	HD 119667	$BD-02^{\circ}3730$	3700	-5.10	M1 Swk	SL90	SL90	1999 Mar 1	MOSAIC
HR 5154	HD 119228	IQ UMa	3600	-3.60	M2 III	Fe90	Hip+La91	1999 Mar 2	MOSAIC
	HD 147923	$BD + 57^{\circ}1671$	3600	-5.20	M2 S	SL90	SL90	1999 May 4	MOSAIC
		$BD + 06^{\circ}2063$	3550	-5.70	M3 S	SL90	SL90	1999 Mar 4	MOSAIC
	HD 189581		3500	-4.90	M3S4*2	SL90	SL90	1997 Jul 20	IRS
HR 8714	HD 216672	HR Peg	3500	-3.30	M3 S5,3	SL86	SL86	1997 Jul 19	IRS
	HD 096360	HL UMa	3550	-5.70	M3 Swk	SL90	SL90	1999 Mar 1	MOSAIC
HR 6039	HD 145713	LQ Her	3460	-3.08	M4 III	Dy98	Hip+La91	1999 May 5	MOSAIC
HR 7139	HD 175588	δLyr	3650	-6.15	M4 II	Dy96, Dy98	Hip+La91	1999 May 6	MOSAIC
HR 7009	HD 172380	XY Lyr	3400	-4.93	M4 II:	Dy96, Dy98	Hip+La91	1999 May 4	MOSAIC
	HD 167539		3450	-5.60	M4 Swk	SL90	SL90	1999 May 5	MOSAIC
HR 8818	HD 218655	DL Gru	3520	-3.70	M4 III	Dy98 T <sub>eff</sub> v Sp	Hip+La91	1997 Jul 20	IRS
HR 5299	HD 123657	BY Boo	3500	-3.10	M4.5 III	Dy96, Dy98, SL85	SL85	1999 Mar 1	MOSAIC
HR 5512	HD 130144	EK Boo	3610	-4.32	M5 III	Dy96, Dy98	Hip+La91	1999 May 6	MOSAIC
HR 4909	HD 112264	TU CVn	3320	-3.79	M5 III	Dy96, Dy98	Hip+La91	1999 May 5	MOSAIC
HR 4949	HD 113866	FS Com	3420	-3.54	M5 III	Dy98	Hip+La91	1999 May 4	MOSAIC
HR 6702	HD 163990	OP Her	3450	-3.40	M5 II-III	Dy96, Dy98, SL85	SL85	1999 May 4	MOSAIC
HR 8637	HD 214966	19 PsA	3410	-3.50	M5 III M5 III	Dy98 $T_{\rm eff}$ v Sp	Hip+La91	1997 Jul 20	IRS
HR 0085	HD 001760	T Cet	3360	-4.50	M5-6 Ib-II	Dy98 $T_{\rm eff}$ v Sp	Be98	2000 Jul 22	OSIRIS
HR 6146	HD 148783	g Her	3380	-5.50	M6 III	Dy96, Dy98, SL85	SL85	1999 Mar 3	MOSAIC
HR 3639	HD 078712	RS Cnc	3190	-5.50	M6 IIIas	Dy96, Dy98, SL86	SL86	1999 Mar 3	MOSAIC
HR 1492	HD 29712	R Dor	3060	-4.00	M8 IIIe	Dy98 $T_{\rm eff}$ v Sp	Dumm98	2000 Jul 22	OSIRIS
111(14)2	HD 207076	EP Aqr	3240	-3.80	M8 IIIvar	Dy98 $T_{\rm eff}$ v Sp	Dumm98	2000 Jul 22 2000 Jul 22	OSIRIS
					Supergiants				
HR 9103	HD 225212	3 Cet	3860	-5.25	K3 Iab	Dy98 T <sub>eff</sub> v Sp	Lu82b, La91	1997 Jul 19	IRS
HR 8726	HD 216946		3650	-4.90	K5 Iab	Dy98 $T_{\rm eff}$ v Sp	EFH85+Hip	BSD96 <sup>c</sup>	
	HD 163428		3800	-6.81	K5II	Lu82a	Hu78+La91	1998 May 18	IRS
		$CD - 60^{\circ} 3621$	3720	-6.29	M0 Ib	Dy98 $T_{\rm eff}$ v Sp	Hu78+La91	1998 May 18	IRS
	HD 316496	KW Sgr	3620	-8.90	MI	Dy98 $T_{\rm eff}$ v Sp	Hu78	1999 Jul 21	OSIRIS
HR 2197	HD 042543	BU Gem	3800	-7.88	M1 Ia-ab	LB80	Hu78+La91	1999 Mar 3	MOSAIC
HR 2061	HD 039801	$\alpha$ Ori	3540	-6.90	M1 I	CSB00	Hip+La91	BSD96 <sup>c</sup>	
HR 2190	HD 042475	TV Gem	3520	-7.88	M1 Iab	Ri98	Hu78+La91	1999 Mar 3	MOSAIC
	HD 143183		3560	-9.00	M1-2 I	Dy98 $T_{\rm eff}$ v Sp	Hu78, La91	1999 Jul 21	OSIRIS
HR 8316	HD 206936	μCep	3510	-10.30	M2 Ia	Dy98 $T_{\rm eff}$ v Sp	Hip+EFH85	BSD96 <sup>c</sup>	
	HD 14469	SU Per	3350	-7.90	M3.5 Ia	Dy98 $T_{\rm eff}$ v Sp	Hip+EFH85	BSD96 <sup>c</sup>	
		KY Cyg	3310	-9.30	M3.9 Iab	Dy98 $T_{\rm eff}$ v Sp	EFH85+MJE95	LVR92 <sup>c</sup>	
	HD 172804	кт су <u></u> д	3400	-6.00	M4S5/6-	SL90	SL90	1999 May 06	MOSAIC
		BC Cyg	3300	-9.30	M4 Ia	Dy98 $T_{\rm eff}$ v Sp	EFH85+MJE95	LVR92 <sup>c</sup>	
					Miras				
HR 5894	HD 141850	R Ser	2800	-3.63	M7 IIIe	Ba98	Hip+BL96	1999 May 06	MOSAIC
HR 5080	HD 117287	R Hya	2660	$-8.30^{d}$	M7 IIIe	Ba98	vL97	1999 Jul 21	OSIRIS
HR 0681	HD 014386	o Cet	2400	-4.80	M7 IIIe	Ba98	vL97	2000 Jul 23	OSIRIS
HR 8992	HD 222800	R Aqr	2570	-4.10	M7 IIIpevar	Ba98	vL97	1999 Jul 21	OSIRIS

Note.—The choice for inclusion with giants or supergiants for stars with luminosity class II or with no luminosity class (the S stars) was made on the basis of their measured CO strength (see  $\S$  3.2.1).

<sup>a</sup> Dy98 = Dyck et al. 1998; Fe90 = Fernandez-Villacanas, Rego, & Cornide 1990; SL90 = Smith & Lambert 1990; SL86 = Smith & Lambert 1986; Dy96 = Dyck et al. 1996; Ram97 = Ramírez et al. 1997; SL85 = Smith & Lambert 1985; Lu82a = Luck 1982a; LB80 = Luck & Bond 1980; CSB00 = Carr et al. 2000; Ri98 = Richichi et al. 1998; Ba98 = Barthes 1998; La86 = Lambert et al. 1986; vB97 = van Belle et al. 1997; Hip = Perryman et al. 1997; La91 = Lafon & Berruyer 1991; Be98 = Bedding & Zijlstra 1998; Dumm98 = Dumm & Schild 1998; Lu82b = Luck 1982b; Hu78 = Humphreys 1978; EFH85 = Elias, Frogel, & Humphreys 1985; MJE95 = Massey, Johnson, & Degioia-Eastwood 1995; BL96 = Benson & Little-Marenin 1996; vL97 = van Leeuwen et al. 1997.

<sup>b</sup> IRS (CTIO 4 m):  $\lambda/\Delta\lambda = 560$ , OSIRIS (CTIO 4 m):  $\lambda/\Delta\lambda = 1200$ , MOSAIC (MDM 1.3 m):  $\lambda/\Delta\lambda = 750$ .

<sup>c</sup> Analysis uses spectrum for this star is presented by Blum et al. 1996b or Lançon & Rocca-Volmerange 1992.

<sup>d</sup> The  $M_{bol}$  for R Hya is more luminous than the predicted maximum for AGB stars given by Paczyński 1970; see also the discussion by Blum et al. 1996b. However, the 1  $\sigma$  uncertainty on its distance encompasses the luminosity limit.

TABLE 2Galactic Center Stars

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	Offseta								
		R.A.	Decl.				OBSERVATION		
Number <sup>b</sup>	NAME <sup>c</sup>	(arcsec)	(arcsec)	$K^{\mathrm{d}}$	$J-K^{\rm d}$	$H-K^{d}$	DATE/REFERENCE	INSTRUMENT <sup>e</sup>	Notes <sup>f</sup>
1		-40.1	-8.0	$10.40\pm0.08$		$2.54\pm0.11$	2000 Jul 22	OSIRIS	
2		-38.3	-22.6	$10.27\pm0.05$		$2.51\pm0.07$	2000 Jul 22	OSIRIS	
4		-34.4	-18.7	$10.36\pm0.04$	$6.85\pm0.07$	$2.74\pm0.06$	2000 Jul 22	OSIRIS	
5		-29.4	-23.2	$10.26\pm0.03$	$6.61\pm0.06$	$2.62\pm0.05$	2000 Jul 22	OSIRIS	
6		-28.5	-40.4	$9.67 \pm 0.03$		$3.02 \pm 0.06$	1998 May 16 IRS		
7		-26.9	18.9	$9.80 \pm 0.03$	$7.15 \pm 0.06$	$2.76 \pm 0.05$	1998 May 15	IRS	
8 9		-24.0 -23.0	18.4 16.4	$10.24 \pm 0.06$	$6.65 \pm 0.07$	$2.51 \pm 0.07$ $2.42 \pm 0.05$	1998 May 18	IRS IRS	
9 10		-23.0 -23.0	10.4	$9.92 \pm 0.04$ $10.44 \pm 0.04$	$6.29 \pm 0.05$ $7.11 \pm 0.09$	$2.42 \pm 0.03$ $2.99 \pm 0.05$	1998 May 17 2000 Jul 22	OSIRIS	
11		-23.0 -22.4	41.9	$9.22 \pm 0.14$	7.11 ± 0.09 	$2.99 \pm 0.03$ $2.28 \pm 0.14$	1998 May 16	IRS	
13		-20.1	-32.2	$10.06 \pm 0.03$	$7.42\pm0.07$	$2.20 \pm 0.14$ $3.33 \pm 0.05$	2000 Jul 22	OSIRIS	
14		-19.3	25.4	$10.48 \pm 0.03$	$6.75 \pm 0.05$	$2.68 \pm 0.05$	2000 Jul 22	OSIRIS	
17		-16.1	18.5	$9.74 \pm 0.03$	$6.07 \pm 0.04$	$2.48 \pm 0.05$	1998 May 16	IRS	
18		-13.7	15.3	$10.18\pm0.04$		$3.16\pm0.05$	1998 May 17	IRS	
19		-13.3	-16.9	$10.14\pm0.03$	$7.57\pm0.15$	$3.01\pm0.05$	1998 May 17	IRS	
23		-9.6	6.5	$9.75\pm0.03$	$6.14\pm0.04$	$2.31\pm0.05$	1998 May 17	IRS	
27		-9.1	-34.8	$9.02\pm0.03$		$2.66\pm0.05$	1998 May 15	IRS	
28	IRS11	-8.4	8.0	$9.17\pm0.07$	$5.95\pm0.07$	$1.99\pm0.08$	1998 May 17	IRS	RAM00
35	IRS 30	-6.6	0.2	$10.49\pm0.05$	$7.11\pm0.16$	$2.49\pm0.11$			
38	IRS 6E	-5.53	-5.08	$10.06\pm0.06$		$4.26 \pm 0.17$	Krabbe et al. 1995		Hot/young star
40		-4.9	-33.8	$10.09 \pm 0.05$		$2.69 \pm 0.06$	1998 May 18	IRS	
41		-4.5	-30.3	$10.14 \pm 0.09$		$2.39 \pm 0.10$	1998 May 17	IRS	
43		-4.3	-21.8	$10.07 \pm 0.03$	$6.13\pm0.05$	$2.29 \pm 0.05$	1998 May 17	IRS	
45 46	IRS 2 IRS 34	-4.2 -4.13	$-10.2 \\ -4.09$	$\begin{array}{c} 10.57 \pm 0.06 \\ 10.48 \pm 0.08 \end{array}$		$3.65 \pm 0.13$ $2.83 \pm 0.16$	Krabbe et al. 1995		Hot/young star
47	IRS 12S	-4.13 -4.1	-14.6	$9.95 \pm 0.05$	$5.85\pm0.06$	$2.03 \pm 0.10$ $2.18 \pm 0.06$	1998 May 17	IRS	110t/ young star
48	IRS 22	-3.9	-32.0	$8.03 \pm 0.03$	$5.13 \pm 0.04$	$1.77 \pm 0.05$	1997 Jul 20	IRS	RAM00
49	IRS 2L	-3.9	-9.6	$11.68 \pm 0.20$			1997 0 41 20	1115	11111100
50	IRS 12N	-3.9	-12.9	$8.58\pm0.04$	$6.95\pm0.05$	$2.83\pm0.06$	1997 Jul 20	IRS	BSD96 LPV
51		-3.5	-7.2	$10.14\pm0.20$		$3.49\pm0.20$			
52	IRS 13E	-3.37	-7.51	$9.82\pm0.13$	$5.71\pm0.14$	$2.34\pm0.14$	Blum et al. 1995a		Hot/young star
53	IRS 3	-2.45	-2.01	$11.16\pm0.11$		$3.06\pm0.07$	Krabbe et al. 1995		Hot/young star
56	IRS 29N	-1.79	-4.41	$9.96\pm0.11$		$4.36\pm0.13$	Krabbe et al. 1995		Hot/young star
60	IRS 14SW	-0.7	-15.1	$10.15 \pm 0.04$	$6.72 \pm 0.07$	$2.59 \pm 0.06$			
64	ID C Z	-0.2	31.5	$10.13 \pm 0.03$	$6.53 \pm 0.05$	$2.51 \pm 0.05$	1998 May 18	IRS	CODOO DAMOO
66	IRS 7	0.0	0.0	$6.40 \pm 0.03$	$6.64 \pm 0.04$	$2.42 \pm 0.10$	1997 Jul 20	IRS	CSB00, RAM00
68 69	IRS 14NE	0.4	$-14.2 \\ -34.0$	$9.75 \pm 0.04$ $9.79 \pm 0.03$	$6.80 \pm 0.06$	$2.64 \pm 0.06$	1998 May 17 1998 May 16	IRS IRS	
70		0.4 0.4	-34.0 -28.2	$9.79 \pm 0.03$ $9.88 \pm 0.06$	$7.02 \pm 0.08$ $6.30 \pm 0.07$	$2.75 \pm 0.05$ $2.49 \pm 0.08$	1998 May 16 1998 May 16	IRS	
70	IRS 16SW	0.4	-28.2	$9.88 \pm 0.00$ $9.60 \pm 0.05$	$5.15 \pm 0.07$	$2.49 \pm 0.03$ $2.00 \pm 0.07$	Krabbe et al. 1995	IKS	Hot/young star
72 <sup>g</sup>	F95B	0.8	-36.3	$9.00 \pm 0.03$ $9.05 \pm 0.04$	$3.74 \pm 0.05$	$1.27 \pm 0.05$	1998 May 15	IRS	RAM00
75	IRS 15NE	1.4	5.6	$8.96 \pm 0.04$	$6.03 \pm 0.05$	$2.41 \pm 0.05$	1997 Jul 20	IRS	He I emission <sup>h</sup>
77	MPE 1.6-6.8	1.58	-7.21	$9.98 \pm 0.06$	$6.00 \pm 0.08$	$2.42 \pm 0.08$	Krabbe et al. 1995		Hot/young star
78	IRS 8	1.88	23.9	$10.49 \pm 0.06$		$4.02 \pm 0.09$	18 May 98	IRS	Hot/young star
79		2.0	42.7	$9.73\pm0.13$	$6.60\pm0.13$	$2.28\pm0.13$	1998 May 16	IRS	
81	IRS 21	2.22	-8.83	$10.40\pm0.05$		$3.95\pm0.04$	Krabbe et al. 1995		Hot/young star
83	IRS 16NE	2.89	-4.90	$9.01\pm0.05$	$5.00\pm0.06$	$1.93\pm0.06$	Tamblyn et al. 1996		Hot/young star
84		3.1	-17.5	$9.88\pm0.04$	$7.38\pm0.08$	$3.06\pm0.07$	1998 May 15	IRS	
91	IRS 9	5.4	-12.6	$8.61 \pm 0.03$	$7.33 \pm 0.07$	$2.24\pm0.04$	1997 Jul 20	IRS	BSD96 LPV
92	IRS 1W	5.42	-5.61	$8.81 \pm 0.04$	$6.21 \pm 0.06$	$3.13 \pm 0.07$	Blum et al. 1995a		Hot/young star
96	IRS INE	7.3	-4.3	$10.00 \pm 0.07$		$2.50 \pm 0.09$	BSD96 <sup>i</sup>		K band only
97	IRS 1SE	7.5	-6.6	$10.23 \pm 0.04$	$6.66\pm0.11$	$2.46\pm0.06$	BSD96 <sup>i</sup>		K band only <sup>j</sup>
98	IRS 10EL	8.1	-1.8	$10.75 \pm 0.09$ 0.76 ± 0.03	$\dots$ 6 50 $\pm$ 0 05		1008 May 14	IDS	
99 102	IRS 28	8.6 10.6	-24.8 -12.1	$\begin{array}{c} 9.76 \pm 0.03 \\ 9.36 \pm 0.03 \end{array}$	$\begin{array}{c} 6.50 \pm 0.05 \\ 6.94 \pm 0.05 \end{array}$	$\begin{array}{c} 2.45 \pm 0.05 \\ 2.81 \pm 0.05 \end{array}$	1998 May 16 1998 May 15	IRS IRS	BSD96 LPV
102	OSU C2	10.8	-12.1 -5.0	$9.36 \pm 0.03$ $10.10 \pm 0.04$	$6.94 \pm 0.05$ $6.20 \pm 0.06$	$2.81 \pm 0.03$ $2.35 \pm 0.06$	1998 May 15 1998 May 17	IRS	D3D70 LF V
105	050 02	13.3	-0.7	$10.10 \pm 0.04$ $8.91 \pm 0.03$	$6.20 \pm 0.00$ $6.28 \pm 0.04$	$2.33 \pm 0.00$ $2.32 \pm 0.05$	1997 Jul 20	IRS	
107		13.7	17.5	$10.10 \pm 0.03$	$6.27 \pm 0.04$	$2.32 \pm 0.03$ $2.35 \pm 0.04$	1997 Jul 20 1998 May 17	IRS	
108	IRS 19	14.4	-25.7	$8.22 \pm 0.04$	$6.59 \pm 0.04$	$2.61 \pm 0.04$	1997 Jul 20	IRS	RAM00
109	IRS 18	14.9	-17.4	$9.50 \pm 0.03$	$6.36 \pm 0.04$	$2.40 \pm 0.04$	1998 May 16	IRS	
110		15.6	-28.6	$10.07\pm0.04$	$7.09\pm0.07$	$2.95\pm0.05$	1998 May 17	IRS	
112	OSUC 3	18.3	44.6	$10.73\pm0.04$		$3.06\pm0.05$	BSD96 <sup>i</sup>		K band only <sup>j</sup>

TABLE 2—Continued

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Notaria         R.A         Deci (arresc)         K <sup>1</sup> $J - K^4$ $H - K^4$ $H - K^4$ Derry Retrieves         Normal           113         11.3         4.6.4         6.80 ± 0.03         5.74 ± 0.02         179 ± 0.01         1997 May 18         IRS         RAM00           114         1.8.3         4.6.4         6.80 ± 0.03         5.74 ± 0.02         177 ± 0.05         1997 May 18         IRS         RAM00           123         -2.3.2         10.40 ± 0.03         5.74 ± 0.04         2.33 ± 0.04         1989 May 15         IRS         RAM00           124         -2.5.5         3.09         9.10 ± 0.03         5.34 ± 0.04         1989 May 15         IRS         RAM00           126         3.53         2.88 ± 0.03         5.47 ± 0.04         1.198 ± 0.04         1989 May 15         IRS         RAM00           128         .040         9.4 ± 0.03         5.34 ± 0.04         2.14 ± 0.05         1997 May 18         IRS         Notaria         IRS         RAM00           131         .040         9.4 ± 0.03         5.34 ± 0.04         2.14 ± 0.04         1998 May 15         IRS         RAM00           135         .041         .057 ± 0.04         2.14 ± 0.04         1998 May 15		Offseta								
			R.A.	Decl.				OBSERVATION		
	Number <sup>b</sup>	NAME <sup>c</sup>	(arcsec)	(arcsec)	K <sup>d</sup>	$J-K^{\rm d}$	$H-K^{d}$	DATE/REFERENCE	Instrument <sup>e</sup>	Notesf
	113		17.7	-0.9	$10.01\pm0.04$		$2.59\pm0.07$	1998 May 18	IRS	
			18.3	44.6	$8.60\pm0.03$	$5.74\pm0.05$	$2.17\pm0.05$	1997 Jul 20	IRS	RAM00
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$					$10.38\pm0.03$					
										<b>D A M</b> (00)
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		IPS 24						•		BSD06 I PV
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$        \begin{array}{ccccccccccccccccccccccccccccc$	135	OSUC 4						•		K band only <sup>j</sup>
	136	IRS 23	42.5	8.2	$8.62\pm0.03$	$6.51\pm0.04$	$2.58\pm0.04$	1997 Jul 20	IRS	
	140		46.8	15.8	$9.44\pm0.03$	$5.73\pm0.04$		1998 May 15	IRS	RAM00
					$10.45\pm0.03$			2		
					$10.45\pm0.03$	$7.24\pm0.09$	$2.89\pm0.05$			
								2000 Jul 23	OSIRIS	
								2000 X 122	OCIDIC	
								2000 Jul 23	OSIRIS	
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	159		-29.0	-33.2	$11.43\pm0.04$		$3.20\pm0.06$			
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	160		-28.3	-47.7	$11.16\pm0.04$		$3.11\pm0.06$			
					$10.84\pm0.17$			2000 Jul 24	OSIRIS	
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$						$6.90\pm0.08$				
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$								2000 Jul 22	OSIDIS	
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$ \begin{array}{cccccccccccccccccccccccccccccccccccc$								2000 Jul 23	OSIRIS	K band only <sup>j</sup>
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$								2000 0 01 20	obiitib	11 ound only
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$					$11.09\pm0.04$					
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	173			-37.1				2000 Jul 23	OSIRIS	K band only <sup>j</sup>
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	174		-7.9	8.0						-
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$										
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$							$2.82\pm0.05$	2000 Jul 24	OSIRIS	
180 $-2.8$ $-49.6$ $10.76 \pm 0.04$ $\dots$ $4.10 \pm 0.06$ $2000 \text{ Jul } 23$ OSIRIS181 $1.6$ $-21.2$ $11.30 \pm 0.04$ $\dots$ $3.39 \pm 0.11$ $11.1022$ OSIRIS182 $2.6$ $35.3$ $10.52 \pm 0.04$ $7.19 \pm 0.08$ $2.84 \pm 0.06$ $2000 \text{ Jul } 22$ OSIRIS185 $5.0$ $-16.2$ $11.30 \pm 0.05$ $\dots$ $3.31 \pm 0.16$ $000 \text{ Jul } 24$ OSIRIS186 $6.0$ $8.8$ $10.94 \pm 0.04$ $\dots$ $3.25 \pm 0.10$ $2000 \text{ Jul } 24$ OSIRIS188 $7.5$ $-38.1$ $12.40 \pm 0.04$ $\dots$ $3.76 \pm 0.10$ $000 \text{ Jul } 24$ OSIRIS199 $11.1$ $-15.3$ $11.03 \pm 0.03$ $7.38 \pm 0.15$ $3.09 \pm 0.05$ $000 \text{ Jul } 24$ OSIRIS190 $13.5$ $-28.4$ $11.20 \pm 0.04$ $\dots$ $3.25 \pm 0.07$ $000 \text{ Jul } 24$ OSIRIS191 $14.6$ $-27.5$ $11.66 \pm 0.06$ $\dots$ $4.51 \pm 0.30$ $023$ $022.6$ $-44.8$ $12.75 \pm 0.05$ $\dots$ $4.03 \pm 0.23$ 194 $25.4$ $-49.7$ $12.75 \pm 0.04$ $\dots$ $3.01 \pm 0.06$ $000 \text{ Jul } 24$ OSIRIS195 $26.3$ $-2.9$ $10.75 \pm 0.04$ $\dots$ $3.01 \pm 0.06$ $000 \text{ Jul } 24$ OSIRIS198 $27.3$ $-24.8$ $10.76 \pm 0.03$ $7.07 \pm 0.08$ $2.89 \pm 0.04$ $2000 \text{ Jul } 24$ OSIRIS199 $30.3$ $-38.8$ $12.21 \pm 0.03$ $\dots$ $3.95 \pm 0.08$ $000 \text{ Jul } 24$ OSIRIS										
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201 $37.8$ $3.7$ $12.01 \pm 0.04$ $3.52 \pm 0.09$	201		51.0	5.7	$12.01 \pm 0.04$		$5.52 \pm 0.09$			

TABLE 2-Continued

		Off	SET <sup>a</sup>						
Number <sup>b</sup>	Name <sup>c</sup>	R.A. (arcsec)	Decl. (arcsec)	K <sup>d</sup>	J-K <sup>d</sup>	$H-K^{d}$	Observation Date/Reference	Instrument <sup>e</sup>	Notes
202		37.9	-42.9	$12.28\pm0.04$		$3.74\pm0.07$			
203		38.8	46.6	$10.73\pm0.08$					
204		40.4	-47.8	$11.51\pm0.04$		$3.80\pm0.06$			
205		40.7	2.7	$10.82\pm0.03$		$2.90\pm0.05$			
206		41.3	-6.2	$11.92\pm0.04$		$4.30\pm0.15$			
208		47.5	-37.7	$11.91\pm0.03$		$4.12\pm0.07$			
209		54.2	-24.1	$12.98\pm0.04$		$4.15\pm0.24$			
210		54.7	7.3	$11.43\pm0.03$		$3.71\pm0.06$			
211		57.8	-36.8	$10.54\pm0.03$	$7.09\pm0.07$	$2.95\pm0.04$	2000 Jul 22	OSIRIS	
212		59.7	15.0	$11.51\pm0.04$		$3.48\pm0.09$			
213		62.4	-34.6	$11.51\pm0.04$		$3.63\pm0.06$			
214		63.3	12.4	$10.61\pm0.04$		$4.06\pm0.08$	2000 Jul 23	OSIRIS	<i>K</i> band only <sup>j</sup>
	VR 5–7 <sup>k</sup>			$7.30\pm0.15$		$2.30\pm0.18$	1998 May 18	IRS	RAM00

Note.—Table 2 is also available in machine-readable form in the electronic edition of the Astrophysical Journal.

<sup>a</sup> Offset is from IRS 7:  $\alpha$  (2000.0) = 17<sup>h</sup>45<sup>m</sup>40<sup>s</sup>,  $\delta$  (2000.0) = -29°00'22".7; see Blum et al. 1996a.

<sup>b</sup> Number corresponds to sequence in Blum et al. 1996a. Data tables in Blum et al. 1996a were truncated at star 147.

<sup>c</sup> Object names are as given by Becklin & Neugebauer 1975, Lebofsky et al. 1982a, Tollestrup, Becklin, & Capps 1989, Krabbe et al. 1995, and Blum et al. 1996a.

 $^{d}$  K magnitude, J-K, and H-K taken from Blum et al. 1996a. For stars 1, 2, 6, 11, 40, 51, 53, 78, 81, 148, 162, 170, 174, 175, and 183 photometry was derived using new unpublished J and/or H images that were not available to Blum et al. 1996a. These new data were obtained with OSIRIS on 1999 March 3, but the field of view did not completely cover the same field as for BSD96, and thus some stars without H magnitudes given by BSD96 were not observed. Improved H magnitudes could not be extracted for several stars that were in the BSD96 field of view, primarily because of crowding.

<sup>e</sup> IRS (CTIO 4 m):  $\lambda/\Delta\lambda = 560$ , OSIRIS (CTIO 4 m):  $\lambda/\Delta\lambda = 1200$ , MOSAIC (MDM 1.3 m):  $\lambda/\Delta\lambda = 750$ .

<sup>f</sup> Objects indicated with CSB00 or RAM00 have been observed at high spectral resolution by Carr et al. 2000 and Ramírez et al. 2000, respectively. LPV candidates are given by Blum et al. 1996b and BSD96. Stars labeled as "hot/young" are previously identified massive young stars that lie in the  $K_0$  selected sample of this paper (see text).

<sup>g</sup> Star 72 is included because it was observed and analyzed at high spectral resolution by Ramírez et al. 2000. This star was not used in the star formation history calculation because its  $K_0$  is below the cutoff adopted (see § 3.3).

<sup>h</sup> The spectrum of IRS 15NE shows strong characteristic cool star features (see Table 6), but also broad He I and Br $\gamma$  emission, suggesting that it is a close blend. Because the cool star features are quite strong and the emission-line stars are generally fainter at K than the brighter M stars in the GC, we assume that the hot star contribution to this object is not sufficient to remove it from our brightness selected sample.

<sup>i</sup> Spectrum taken from Blum et al. 1996b and BSD96.

<sup>j</sup> Analysis based on *K*-band spectrum only; see text (§ 2.3).

<sup>k</sup> This star is located in the Quintuplet cluster (Moneti, Glass, & Moorwood 1994) and is included for reference. Photometry is taken from Moneti et al. 1994.

approximately  $\pm 10$  pixels. Brackett absorption features in the telluric standards were "fixed" by drawing a line across the feature from continuum points on either side.

The wavelength dispersion solutions were determined from OH lines observed in the *H* and *K* bands and the line positions given by Oliva & Origlia (1992). The wavelength zero point was set by moving the observed position of the CO 2.3  $\mu$ m band head to 2.2935  $\mu$ m (Kleinmann & Hall 1986).

These spectra were then multiplied by a  $\lambda^{-4}$  spectrum approximating the blackbody curve for the telluric standards. For the GC stars, a correction was made for the interstellar reddening assuming the extinction law given by Mathis (1990) and the derived  $A_K$  from BSD96.

#### 3. SPECTRAL CLASSIFICATION

### 3.1. Description of the Technique

BSD96 used K-band spectroscopic indexes for CO and  $H_2O$  to provide a two-dimensional classification yielding  $T_{eff}$  and  $M_{bol}$  for cool, luminous GC stars (following Kleinmann & Hall 1986; see also Ramírez et al. 1997). The CO index is a measure of the strength of the 2.2935  $\mu$ m 2–0 <sup>12</sup>CO rovibrational band head (Kleinmann & Hall 1986). The H<sub>2</sub>O feature is a broad depression of the continuum between the H and K

bands due to myriad blended steam absorption lines. There are similar steam absorption bands between the J and H bands and between the K and L bands (Strecker, Erickson, & Witteborn 1978). The latter band can affect the CO band head region for stars with extreme H<sub>2</sub>O absorption strength (see § 3.2.1). The combination of H<sub>2</sub>O and CO features has been used in the past to break the degeneracy in luminosity class and  $T_{\rm eff}$  versus absorption strength that exists for each feature alone (see the extensive discussion in Kleinmann & Hall 1986 and BSD96). The correlation of band strength with luminosity is positive for CO (CO increases in stars of higher luminosity) and negative for H<sub>2</sub>O (H<sub>2</sub>O decreases for stars of higher luminosity).

Several improvements have been made in the present work, relative to the analysis in BSD96. Both *H*- and *K*-band coverage are used to define the H<sub>2</sub>O absorption, whereas BSD96 had only *K*-band spectra. In that case, the derived H<sub>2</sub>O strength was sensitive to the interstellar extinction and reddening for any given star since the relative depression of the blue end of the *K* band depends sensitively on  $A_K$  and the assumed wavelength ( $\lambda^{-1.7}$ ) dependence of the reddening law. Using both *H*- and *K*-band spectra, the H<sub>2</sub>O absorption is seen as a broad "feature" spanning the blue end of the *K* band and the red end of the *H* band (as well as a downturn in stellar flux at the blue end of the

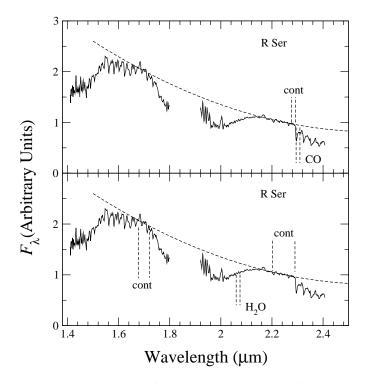


FIG. 1.—*H* and *K* spectra of the AGB star R Ser (M7 III, Mira) used to demonstrate the CO and H<sub>2</sub>O measurements for all spectra. The CO strength is determined by the ratio of flux in the band centered at 2.302  $\mu$ m compared with the flux in the continuum band (in the star, not the ft) at 2.2875  $\mu$ m (the bands are indicated with vertical dashed lines in the upper panel). The dashed curves are quadratic fits to the continuum in bands at 1.68–1.72  $\mu$ m and 2.20–2.29  $\mu$ m (*lower panel*). The H<sub>2</sub>O strength is measured using the flux in a band at 2.060–2.075  $\mu$ m relative to the flux in the fit at the same position (*lower panel*).

*H* band and the red end of the *K* band). By combining spectra of both the *H* and *K* bands, in continuum regions that are not affected by stellar steam absorption, the intrinsic stellar steam absorption at 2.07  $\mu$ m can be distinguished from the interstellar reddening, which produces a monotonic decrease in flux toward bluer wavelengths. In this paper the CO versus  $T_{\text{eff}}$  relation has also been improved, placing it on a more quantitative basis that relies less on spectral types and more on  $T_{\text{eff}}$  measurements. In terms of sample size, we are presenting a GC sample more than 3 times larger than the one presented by BSD96. A larger sample is crucial for constraining the SFH through theoretical models.

The CO index is defined (BSD96) as the percentage of flux in the CO 2.3  $\mu$ m feature relative to a continuum band centered at 2.284  $\mu$ m [ $(1 - F_{band}/F_{cont}] \times 100$ ] The CO band and continuum band were 0.015  $\mu$ m wide, and the CO band was centered at 2.302  $\mu$ m. The CO index is only marginally affected by extinction since the CO and continuum bands are closely spaced. A typical CO strength of 20% changes by about 1% for a change in  $A_K$  of 1 mag. This is similar to the typical uncertainty in the derived CO strength, which is taken as the 1  $\sigma$  uncertainty in feature strength derived from the pixel-to-pixel variation in the nearby continuum. The CO and associated continuum band are graphically represented in the upper panel of Figure 1.<sup>4</sup>

The H<sub>2</sub>O strength is defined similarly to the CO index, but with a quadratic fit to the continuum using bands at 1.68–1.72 and 2.20–2.29 µm (Fig. 1, lower panel) and a band 0.015  $\mu$ m wide centered at 2.0675  $\mu$ m (Fig. 1, lower panel). The difference between this index and that used by BSD96 is that the latter index used the same continuum band as the CO index (hence the sensitivity in that work to the adopted extinction for any given star). The formal uncertainty in the H<sub>2</sub>O strength measured here is a fraction of a percent; the actual uncertainty, 3%, is derived from the scatter of H<sub>2</sub>O measurements for supergiants, which have no measurable H<sub>2</sub>O. Small changes in the choice of continuum bands used for the fits can lead to systematic changes in the derived H<sub>2</sub>O strength of  $\leq 5\%$ -10%, but these tend to affect all the spectra similarly. The systematic uncertainty should be much smaller than that given by BSD96 since the continuum fit spans the H and K bands and is thus insensitive to the details of the reddening.

#### 3.2. Comparison Stars

The comparison star list is given in Table 1. These stars were chosen to span the range of expected GC star  $M_{bol}$  and  $T_{eff}$  (BSD96; Carr et al. 2000; Ramírez et al. 2000). The literature was surveyed for cool giant, asymptotic giant branch (AGB), and cool supergiant stars with derived  $T_{eff}$  and  $M_{bol}$ . References for these two quantities are given in Table 1. In addition, we have used five digital spectra previously presented by BSD96 and Lançon & Rocca-Volmerange (1992) for several comparison stars, which we were not able to reobserve (see Table 1) but which met our criteria on  $M_{bol}$ and  $T_{eff}$ . We observed two stars (FL Ser and Z UMa) that we later removed from our sample because they had discrepant  $T_{eff}$  compared with stars of similar spectral type (Dyck, van Belle, & Thompson 1998).

#### 3.2.1. CO Index

The comparison stars were used to define a CO index versus  $T_{\rm eff}$  relation, which we then used to determine the  $T_{\rm eff}$  of the GC stars (see below). In order to produce a relationship between CO and  $T_{\text{eff}}$ , comparison stars were selected with fundamental (e.g., Dyck et al. 1996)  $T_{\text{eff}}$  determinations or  $T_{\text{eff}}$  determined from detailed spectroscopic analysis. Unfortunately, the list of such stars matching the GC  $T_{\rm eff}$ range and also with  $M_{\rm bol}$  determined was too small (only four supergiants). Thus, for a number of stars in Table 1 the  $T_{\rm eff}$  versus spectral type relation of Dyck et al. (1998) was used to extend our sample of comparison stars (mainly for the supergiants). A small offset in  $T_{\text{eff}}$  (-400 K for K-type stars and -220 K for M-type stars) was made for supergiants (Dyck et al. 1996).  $T_{\rm eff}$  as determined from the relationship in Dyck et al. (1998) was substituted for two stars whose independent  $T_{\rm eff}$  published in the literature gave substantially larger residuals relative to the derived CO versus  $T_{\rm eff}$  relationship (3 Cet [Luck & Bond 1980] and CD -60°3621 [ Luck & Bond 1989]).

Following Dyck et al. (1996), two stars of luminosity class II were treated as giants (XY Lyr and  $\pi$  Her), and this is consistent with their measured CO values compared with similar stars of the same spectral type. Conversely, HD 163428 (luminosity class II) was included in the supergiant category since its CO is consistent with that group. We made analogous assignments for the S stars shown in Table 1 that have no luminosity class given explicitly. In all cases, the

<sup>&</sup>lt;sup>4</sup> The data at wavelengths between 1.80 and 1.92  $\mu$ m have been omitted from the plots in Figs. 1, 3, 4, and 7–10 because of the low telluric transmission in these regions.

associated  $M_{bol}$  is consistent with the luminosity class chosen, although there is overlap between luminous giants (AGB stars) and less luminous supergiants (see the extensive discussion in BSD96).

The tabulated  $T_{\text{eff}}$  and computed CO indexes for the comparison stars (see Table 3) were used to derive a linear relationship for  $T_{\text{eff}}$  versus CO strength. The relationship (a least-squares fit to the data) is shown for both giants

TABLE 3
COMPARISON STAR INDEXES

					СО	H <sub>2</sub> O					
HR	HD	Other	Spectral Type	$T_{\rm eff}$	(%) <sup>a</sup>	(%) <sup>a</sup>					
Giants											
HR 6418	HD 156283	$\pi$ -Her	K3 II	4100	$13\pm0.7$	$1\pm3$					
HR 4299	HD 095578	61-Leo	M0 III	3700	$15\pm0.4$	$2\pm3$					
	HD 119667	$BD - 02^{\circ}3726$	M1 Swk	3700	$17\pm0.5$	$1\pm 3$					
HR 5154	HD 119228	IQ-UMa	M2 III	3600	$16\pm0.5$	$-1\pm 3$					
	HD 147923	$BD + 57^{\circ}1671$	M2 S	3600	$17 \pm 0.5$	$0\pm 3$					
		$BD + 06^{\circ}2063$	M3 S	3550	$15\pm0.5$	$0\pm 3$					
	HD 189581		M3 S4*2	3500	$16 \pm 0.5$	$3\pm3$					
HR 8714	HD 216672	HR Peg	M3 S5,3	3500	$20\pm0.5$	$-2 \pm 3$					
	HD 096360	HL UMa	M3 Swk	3550	$16 \pm 0.6$	$2\pm 3$					
HR 6039	HD 145713	LQ Her	M4 III	3460	$10 \pm 0.0$ $17 \pm 0.6$	$0\pm 3$					
HR 7139	HD 175588	δLyr	M4 II	3650	$17 \pm 0.0$ $15 \pm 0.8$	$-4 \pm 3$					
HR 7009	HD 172380	XY Lyr	M4 II:	3400	$10 \pm 0.0$ $19 \pm 0.5$	$-2 \pm 3$					
III (700)	HD 167539		M4 Swk	3450	$17 \pm 0.0$ $17 \pm 0.6$	$1\pm 3$					
HR 8818	HD 218655	DL Gru	M4 III	3520	$17 \pm 0.0$ $16 \pm 0.7$	$-1 \pm 3$					
HR 5299		BY Boo									
	HD 123657		M4.5 III	3500	$18 \pm 0.7$	$0 \pm 3$					
HR 5512	HD 130144	EK Boo	M5 III	3610	$17 \pm 0.6$	$-4 \pm 3$					
HR 4909	HD 112264	TU CVn	M5 III	3320	$18 \pm 0.6$	$4 \pm 3$					
HR 4949	HD 113866	FS Com	M5 III	3420	$17 \pm 0.4$	$-7 \pm 3$					
HR 6702	HD 163990	OP Her	M5 II-III	3450	$20 \pm 0.5$	$-1 \pm 3$					
HR 8637	HD 214966	19 PsA	M5 III	3410	$16 \pm 0.7$	$-3\pm3$					
HR 0085	HD 001760	T Cet	M5-6 Ib-II	3360	$20 \pm 1.1$	$-9\pm3$					
HR 6146	HD 148783	g Her	M6 III	3380	$17 \pm 0.7$	$1\pm3$					
HR 3639	HD 078712	RS Cnc	M6 IIIas	3190	$19 \pm 0.5$	$5\pm3$					
HR 1492	HD 29712	R Dor	M8 IIIe	3060	$22 \pm 1.4$	<sup>b</sup>					
	HD 207076	EP Aqr	M8 IIIvar	3240	$19\pm1.4$	<sup>b</sup>					
		Supergi	ants								
HR 9103	HD 225212	3 Cet	K3 Iab	3860	$19\pm0.6$	$1\pm3$					
HR 8726	HD 216946		K5 Iab	3650	$20 \pm 0.5$	$2\pm3$					
	HD 163428		K5 II	3800	$19\pm0.5$	$0\pm 3$					
		CD -60-3621	M0 Ib	3720	$22 \pm 0.6$	$-3\pm3$					
	HD 316496	KW Sgr	MI	3620	$23 \pm 1.4$	$-2\pm 3$					
HR 2197	HD 042543	BU Gem	M1 Ia-ab	3800	$22 \pm 0.6$	$1\pm 3$					
HR 2061	HD 039801	$\alpha$ Ori	M1 I	3540	$\frac{22}{23} \pm 0.5$	$2\pm 3$					
HR 2190	HD 042475	TV Gem	M1 Iab	3520	$23 \pm 0.3$ $23 \pm 0.7$	$1\pm 3$					
IIIX 2190	HD 143183		M1-2I	3570	$23 \pm 0.7$ $24 \pm 1.4$	$-3\pm3$					
HR 8316	HD 206936	 	M1-21 M2 Ia	3510	$24 \pm 1.4$ $25 \pm 0.5$	$-3 \pm 3$ $4 \pm 3$					
ПК 8510	HD 200930 HD 14469	$\mu$ Cep	M3.5 Ia								
	ПD 14409	SU Per		3350	$26 \pm 0.5$	$6 \pm 3$					
		KY Cyg	M3.9 Iab	3310	$25 \pm 1.7$	$8 \pm 3$					
	HD 172804		M4S5/6-	3400	$26 \pm 0.7$	$1 \pm 3$					
	•••	BC Cyg	M4 Ia	3300	$25 \pm 1.7$	$-2 \pm 3$					
		Mira Var	iables								
HR 5894	HD 141850	R Ser	M7 IIIe	2800	$21\pm0.6$	$15\pm3$					
HR 5080	HD 117287	R Hya	M7 IIIe	2660	$25\pm1.1$	$43\pm3$					
HR 0681	HD 014386	o Cet	M7 IIIe	2400	$16\pm1.7$	$37\pm3$					
HR 8992	HD 222800	R Aqr	M7 IIIpevar	2570	$23\pm1.2$	$34\pm3$					

<sup>a</sup> Indexes are percent absorption defined as  $100 \times (1 - \text{flux/cont})$  where "flux" and "cont" are the integrated fluxes in 0.015  $\mu$ m bands in the spectra. For CO, the "flux" band is centered at 2.302  $\mu$ m and the "cont" band is taken at 2.284  $\mu$ m. For H<sub>2</sub>O, "flux" is for a band centered at 2.0675  $\mu$ m and "cont" is the same but derived from a fit to the spectral continuum (see text). The uncertainty in the measured H<sub>2</sub>O strength is taken from the scatter in the measurements of IIIs and Is; see text.

<sup>b</sup> Continuum may have been affected by data taking procedure; see text. Strong expected H<sub>2</sub>O absorption not evident.

<sup>c</sup> *K*-band spectrum only; continuum fit made in the same way as for stars with *H* and *K*, but fitted only in the region between 2.20 and 2.29  $\mu$ m.

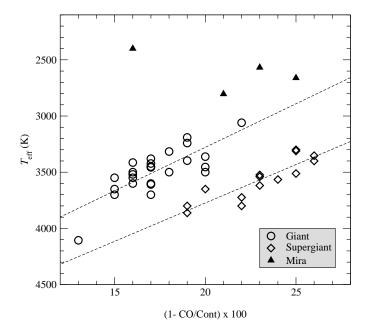


FIG. 2.—The 2.3  $\mu$ m CO absorption strength for the comparison stars. CO strength increases with decreasing  $T_{\text{eff}}$  but also with increasing luminosity (see text). The correlation appears to break down for some LPVs (Mira variables; see text).

 $(T_{\rm eff} = 4828.0-77.5 \times {\rm CO})$  and supergiants  $(T_{\rm eff} = 5138.7-68.3 \times {\rm CO})$  in Figure 2. The offset between the two luminosity classes is due to the effects described in Kleinmann & Hall (1986) and BSD96. Figures 3 and 4 show several supergiant and giant spectra from our sample.

The four points for long-period variables (LPVs) shown in Figure 2 effectively constitute a third relationship for  $T_{\text{eff}}$ 

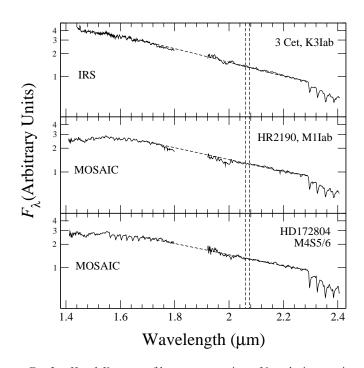
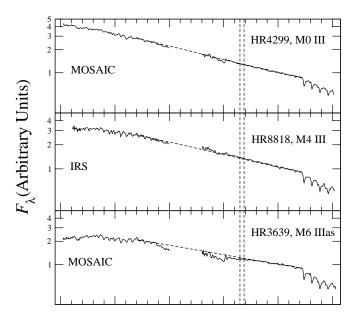


FIG. 3.—*H* and *K* spectra of late-type supergiants. Note the increase in CO (2.3  $\mu$ m) absorption strength for later types. The dashed curves are the fits to the continua used to measure H<sub>2</sub>O at the position of the vertical dashed lines (see text and Fig. 1). The *y*-axis is scaled logarithmically.



# Wavelength $(\mu m)$

FIG. 4.—Same as Fig. 3, but for late-type giants. HR 3639 has similar CO strength compared with 3 Cet, but visibly stronger  $H_2O$  (Fig. 3). The *y*-axis is scaled logarithmically.

versus CO strength; however, the CO strength versus  $T_{\rm eff}$  for LPVs does not correlate in the same way as for giants and supergiants (increase in CO strength for lower  $T_{\rm eff}$ ). This is most likely because there is such strong H<sub>2</sub>O absorption in these stars that the CO continuum band near 2.29  $\mu$ m is affected (depressed; see Figs. 1 and 2). The coolest LPVs have weaker CO. We discuss the determination of  $T_{\rm eff}$  for the LPVs in the next section.

## 3.2.2. H<sub>2</sub>O Index

Given a measured CO strength,  $T_{\rm eff}$  can be determined using the calibration from the comparison stars (Fig. 2), provided the luminosity class is known. For purposes of determining  $M_{\rm bol}$ , we take "LPV" as a distinct luminosity class because we assign different bolometric corrections to giants (AGB, III), supergiants (I), and LPVs (Mira variables, tip of the AGB).

The H<sub>2</sub>O strengths for our comparison stars are generally small for both the giants and supergiants because the giants are dominated by earlier M spectral types (a consequence of our  $M_{bol}$  and  $T_{eff}$  selection), and the supergiants do not exhibit strong H<sub>2</sub>O as described above. H<sub>2</sub>O strengths are given in Table 3 for the comparison giants and supergiants. The quoted uncertainty is derived from the scatter of all the giant and supergiant stars. The LPV comparison stars, on the other hand, exhibit strong H<sub>2</sub>O absorption, 32% on average for the four stars listed in Table 3.

For purposes of classification, we will assume that the  $H_2O$  strengths of late-M stars on the AGB lie between those for giants and LPVs. We take the  $H_2O$  strength of R Ser  $(H_2O = 15\%)$  to be the lower bound for stars that are likely to be LPVs (see Fig. 5). This star is classified as a Mira (i.e., an LPV) but has less  $H_2O$  absorption than the other luminous LPVs in our sample. Thus any star with  $H_2O > 15\%$  in the GC will be classified as an LPV candidate. Table 3 shows

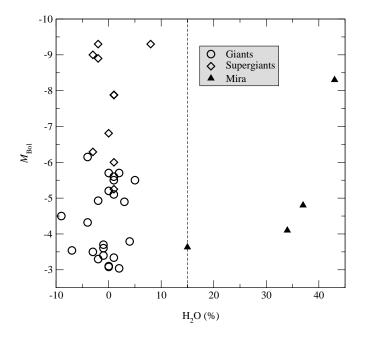


FIG. 5.—H<sub>2</sub>O absorption strength for comparison stars. Stars with H<sub>2</sub>O  $\geq$  15% are known Mira variables (LPVs). All GC stars with H<sub>2</sub>O greater than 15% were classified as LPV candidates with a correspondingly lower  $T_{\text{eff}}$  (see text).

that  $T_{\rm eff}$  depends on the H<sub>2</sub>O index for LPVs, in the sense that H<sub>2</sub>O is stronger for cooler stars. We use this relationship ( $T_{\rm eff} = 2893.0-8.8 \times H_2$ O) to derive  $T_{\rm eff}$  for GC stars that are likely to be LPVs based on their H<sub>2</sub>O strength. LPVs are assigned an uncertainty in  $T_{\rm eff}$  based on the full range of values for the comparison stars ( $\pm 200$  K).

For non-LPVs in the GC, the relations between CO and  $T_{\rm eff}$  for giants and supergiants, illustrated in Figure 2, are used to derive  $T_{\rm eff}$ . The distinction between giants and supergiants, however, is not entirely straightforward. The classifications were initially made by eye, paying attention to the appearance or absence of H<sub>2</sub>O between the *H* and *K* bands for a given CO strength. However, we then determined what quantitative values of CO, H<sub>2</sub>O, and  $M_{\rm bol}$  reproduced our classifications by eye. We used  $M_{\rm bol}$  to make an initial distinction (see § 4.1 for the derivation of  $M_{\rm bol}$ ): GC stars with  $M_{\rm bol} < -7.2$  must be supergiants (see discussion in BSD96), while GC stars with  $M_{\rm bol} > -4.9$  (the faintest  $M_{\rm bol}$  known for supergiants; see Table 1 and BSD96) are likely to be giants. The GC stars with  $-4.9 > M_{\rm bol} > -7.2$  could be either giants or supergiants, based solely on  $M_{\rm bol}$ .

could be either giants or supergiants, based solely on  $M_{bol}$ . For the GC stars with  $-4.9 > M_{bol} > -7.2$ , we use a combination of CO and H<sub>2</sub>O to distinguish giants from supergiants. This is shown in Figure 6. Supergiants have CO > 20% and a linear relationship between H<sub>2</sub>O and CO defining the upper envelope of supergiants, such that H<sub>2</sub>O < -5 + (0.5 × CO), between CO indexes of 20% and 26%, for supergiants. Giants are all the remaining GC stars that have not previously been classified as LPVs or supergiants by these techniques. The bolometric correction (§ 4.1) depends on assigned luminosity class, and most of the giants with  $M_{bol} > -4.9$  could be more luminous than  $M_{bol} = -4.9$  if the luminosity class was assumed originally to be supergiant instead. However, each case would then result in a giant classification based on the linear relation between CO and H<sub>2</sub>O. Thus, the original assignments based

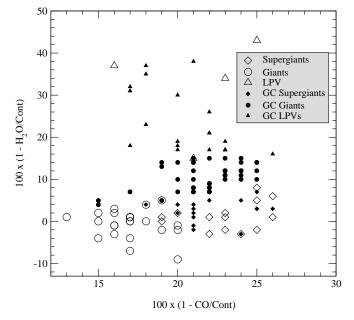


FIG. 6.—H<sub>2</sub>O strength vs. CO strength for the GC and comparison stars. The comparison stars are shown as open symbols, while GC stars are plotted with filled symbols. For a given CO strength, GC stars classified as III have larger H<sub>2</sub>O than those classified as I. The GC III stars lie between the warmer comparison stars and cooler comparison star LPVs. GC stars with H<sub>2</sub>O > 15% are classified as LPVs or LPV candidates (LPV? in Table 6).

on the appearance of  $H_2O$  are consistent with the assignments based on the measured CO,  $H_2O$ , and derived  $M_{bol}$ . For stars in common with BSD96, we have arrived at the same luminosity classes, and also for star VR 5-7 of the Quintuplet cluster. Moneti, Glass, & Moorwood (1994) classified VR 5-7 as a late-type M supergiant.

As an aside, we note that for two of the comparison stars (R Dor and EP Aqr) we could not compute reliable  $H_2O$  indexes owing to possible low-order variations and/or slope changes in the continuum. This may be due to the nonstandard data-taking procedures that were used to keep these bright stars from saturating (see above), although most of the stars appear to have normal spectra independent of how they were obtained. These variations should not affect the relatively narrow and closely spaced CO and associated continuum measurements. R Dor and EP Aqr are the two coolest giants not identified as Mira variables or LPVs in our sample, yet our spectra show no evidence of  $H_2O$  absorption when fitted for a continuum like the other stars.

# 3.3. Galactic Center Sample

The GC stars are listed in Table 2. The complete GC sample was chosen as all stars in the GC *K*-band luminosity function (KLF; taken here to be the *dereddened* luminosity function) as derived by BSD96, which have  $K_0 \leq 7.0$  (where  $K_0$  is the dereddened *K* magnitude). We have revised several stars to  $K_0$  magnitudes fainter than 7.0 on the basis of new, unpublished *H*-band images used to determine  $A_K$  (see the footnotes to Tables 2 and 4). With these adjustments, there are 136 stars in the complete sample. However, 22 stars in this original list were not observed because of severe crowding with neighboring stars, and we were unable to observe an additional 24 stars because of cloudy weather at the telescope. Eleven stars in the list are known emission-line stars

TABLE 4	
Galactic Center Stars Younger than $10\ \mathrm{Myr}$	

Name <sup>a</sup>	Number <sup>b</sup>	$T_{\rm eff}$ / Spectral Type <sup>c</sup>	$K_0{}^{\mathrm{d}}$	$A_K^{d}$	$M_K^{\rm e}$	$M_{\rm bol}^{\rm c}$
IRS 6E	38	WC9	$3.77\pm0.28$	$6.28\pm0.28$	-10.75	
IRS 34	46	WC9	$6.47\pm0.12$	$4.01\pm0.22$	-8.05	
IRS 13E	52	29000	$6.85\pm0.18$	$2.97\pm0.12$	-7.67	-11.1
IRS 3	53	Featureless	$6.78\pm0.12$	$4.38\pm0.11$	-7.74	
IRS 29N	56	WC9	$3.56\pm0.25$	$6.40\pm0.10$	-10.96	
IRS 7	66	$3430\pm240$	$2.92\pm0.16$	$3.48\pm0.09$	-11.60	-9.0
IRS 16SW	71	24000	$7.00\pm0.08$	$2.60\pm0.06$	-7.52	-11.3
MPE 1.6-6.8	77	WC9	$6.65\pm0.11$	$3.28\pm0.08$	-7.87	
IRS 8	78	Featureless	$4.59\pm0.15$	$5.90\pm0.14$	-9.93	
IRS 21	81	Featureless	$4.61\pm0.08$	$5.79\pm0.06$	-9.91	
IRS 16NE	83	24000	$6.55\pm0.06$	$2.45\pm0.05$	-7.97	-11.0
IRS1W	92	Featureless	$6.29\pm0.07$	$2.42\pm0.06$	-8.23	
VR 5-7		$3570\pm150$	$4.13\pm0.30$	$3.17\pm0.28$	-10.22	-7.6

NOTE.—Only stars with  $K_0 \leq 7.0$  from the list of Blum et al. 1996b are included in this table.

<sup>a</sup> Object names are as given by Becklin & Neugebauer 1975, Lebofsky et al. 1982a, Tollestrup, Becklin, & Capps 1989, Krabbe et al. 1995, and Blum et al. 1996a.

<sup>b</sup> Number corresponds to sequence in Blum et al. 1996a. Data tables in Blum et al. 1996a were truncated at star 147.

 $^{\circ}T_{\text{eff}}/M_{\text{bol}}$  for hot stars IRS 16NE, 16SW, and 13E from Najarro et al. 1997.  $T_{\text{eff}}/M_{\text{bol}}$  for IRS 7 and VR 5-7 from present work. Spectral types from reference listed in Table 2.

 ${}^{d}K_0 (=K-A_K)$  and  $A_K$  taken from Blum et al. 1996a. Stars 53, 78, and 81 use new unpublished *H*-band magnitudes that were unavailable to Blum et al. 1996a.  $K_0$  and  $A_K$  for VR 5-7 derived using photometry in Table 2.

<sup>e</sup>  $M_K$  is derived from  $K_0$  (see Blum et al. 1996a) and a distance modulus of 14.52 (8 kpc; Reid 1993).

or have featureless continua (BSD96; Tables 2 and 4). The featureless stars are apparently young massive stars that are embedded in ionized gas filaments in the central parsec (Tanner et al. 2003). We thus take them to be part of the youngest burst of star formation in the GC. All 136 stars are listed in Table 2. Observation dates and instruments are given for the stars for which we obtained new spectra.

Table 5 indicates the level of completeness in the spectroscopic sample as a function of luminosity. In what follows we assume that the stars in our complete list that we were not able to observe are distributed in the H-R diagram like the cool stars we have observed. This is a good assumption since previously known emission-line objects (apart from those in Table 5) are fainter than  $K_0 = 7.0$ , and all of the stars that we observed in this work from the complete sample were cool stars.

Example spectra of GC supergiant, giant, AGB, and LPV stars are shown in Figures 7, 8, 9, and 10, respectively. The complete sample of cool stars observed at high spectral resolution by Carr et al. (2000) and Ramírez et al. (2000) is shown in these figures, including all stars for which the values of  $T_{\rm eff}$  in this paper were used in the abundance analysis of Ramírez et al. (2000).

Table 6 lists the measured CO and  $H_2O$  strengths for each GC star. Four stars in our list are taken from BSD96 and have only *K*-band spectra (IRS 1NE, IRS 1SE, OSUC 3, and OSUC 4). Classifications for three other stars (stars 170, 173, and 214) were made based only on their *K*-band spectra because use of the *H*-band spectra (which was much fainter) in these cases resulted in poor fits for the continuum and hence  $H_2O$  indexes, which were clearly in error.

The GC star luminosity class is given in Table 6 for each GC star. As noted above, any GC star with  $H_2O > 15\%$  is classified as an LPV (or LPV candidate [LPV?] in Table 6). We retained the classification of LPV for stars IRS 9, IRS 12N, IRS 24, and IRS 28 of BSD96. These are known

photometric variables (Haller et al. 1992; Tamura et al. 1996; BSD96; Ott, Eckart, & Genzel 1999). For LPVs and LPV?s, the average  $H_2O$  was  $23\% \pm 7\%$ . Twenty GC stars of the spectroscopic sample of 79 were classified as LPVs.

GC stars with lower  $H_2O$  were binned into III and I classes on the basis of the appearance (by eye) of the overall

 TABLE 5

 Completeness and Characteristics of the GC Sample

$K_0{}^a$	Percent Complete <sup>b</sup> (%)	Number Observed	Number in K <sub>0</sub> LF
$K_0 < 4.0$	100	3	3
$4.0 < K_0 < 4.5$	100	0	0
$4.5 < K_0 < 5.0$	100	7	7
$5.0 < K_0 < 5.5$	50	5	10
$5.5 < K_0 < 6.0$	85	17	20
$6.0 < K_0 < 6.5$	78	18	23
$6.5 < K_0 < 7.0$	55	40	73
Total (cool+hot/young)		90	136
Stars by type:			
Number of GC Is			15 <sup>c</sup>
Number of GC IIIs			43
Number of GC LPVs			20
Stars younger than 10 Myr			12 <sup>c</sup>
Stars not observed			46
Total			136

<sup>a</sup> The original sample is selected from the complete list of stars with  $K_0 < 7.0$  mag in the central few parsecs of the Galaxy presented by Blum et al. 1996a.

<sup>b</sup> Percentage of stars observed spectroscopically from the list of Blum et al. 1996a, which includes known emission-line stars or featureless young stars; see text and Tables 2 and 4.

<sup>c</sup> IRS 7, M I, is counted among the young stars. VR 5-7 of the Quintuplet cluster is not counted in this table, nor is star 72, which is listed in Tables 2 and 6.

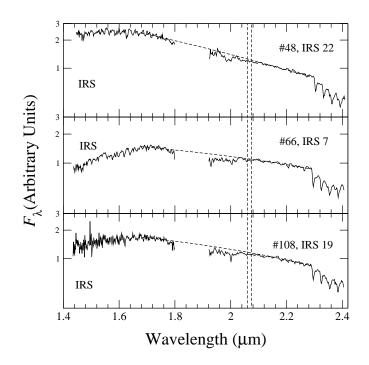


FIG. 7.—Example spectra of stars classified as supergiants in the GC. These stars were analyzed at high spectral resolution by Ramírez et al. (2000) and Carr et al. (2000). The dashed curves are the fits to the continua used to measure  $H_2O$  at the position of the vertical dashed lines (see text and Fig. 1). The *y*-axis is scaled logarithmically.

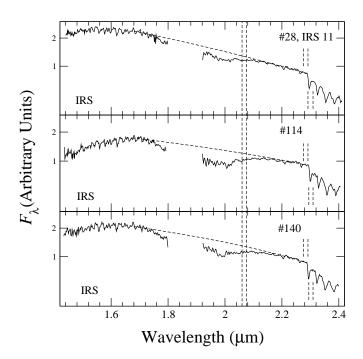
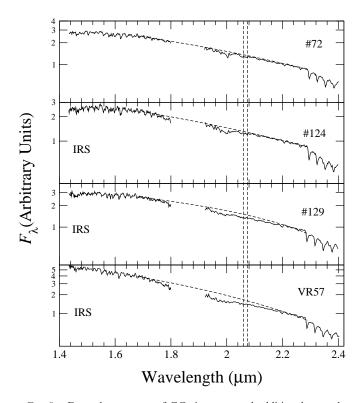


FIG. 9.—Example spectra of stars classified as asymptotic giant branch (referred to in the text as AGB, giant, or III stars) stars in the GC. Stars 28, 114, and 140 were analyzed at high spectral resolution by Ramírez et al. (2000). The dashed curves are the fits to the continua used to measure  $H_2O$  at the position of the vertical dashed lines (see text and Fig. 1). The *y*-axis is scaled logarithmically.



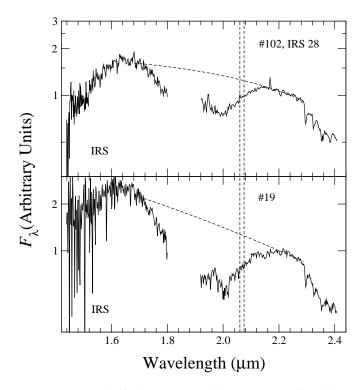


FIG. 8.—Example spectrum of GC giant star and additional examples of GC supergiants (see Fig. 7). Star 72 is classified as a giant; all others are supergiants. These stars were analyzed at high spectral resolution by Ramírez et al. (2000).

FIG. 10.—Stars classified as LPVs or candidate LPVs (LPV? in Table 6) in the GC. IRS 28 is a known photometric variable (see text). The emission line near 2.17  $\mu$ m evident in the spectrum of IRS 28 is likely due to incomplete subtraction of the local nebular background. The dashed curves are the fits to the continua used to measure H<sub>2</sub>O at the position of the vertical dashed lines (see text and Fig. 1). The *y*-axis is scaled logarithmically.

TABLE 6Galactic Center Star Properties

Number	Name	CO (%)	H <sub>2</sub> O (%)	Luminosity Class <sup>a</sup>	$T_{\rm eff}{}^{\rm b}$	$K_0^{c}$	$M_K^{\rm d}$	$A_{K}^{c}$	$BC_K^e$	$\Delta A_K^{\rm f}$	$M_{\rm bol}{}^{ m g}$
1		$21\pm0.6$	-2	Ι	$3710\pm130$	$6.85\pm0.12$	-7.67	$3.55\pm0.09$	2.6	0.0	$-5.07 \pm 0.42$
2		$20\pm0.8$	17	LPV?	$2740\pm200$	$6.76\pm0.12$	-7.76	$3.51\pm0.11$	3.2	0.5	$-4.06\pm0.42$
4		$26\pm0.9$	16	LPV?	$2750\pm200$	$6.70\pm0.10$	-7.82	$3.67\pm0.06$	3.2	0.5	$-4.12\pm0.41$
5		$24\pm0.9$	10	III	$2930\pm150$	$6.74\pm0.06$	-7.78	$3.51\pm0.05$	3.2	0.4	$-4.18\pm0.40$
6		$17 \pm 2.1$	32	LPV?	$2610 \pm 200$	$5.36 \pm 0.11$	-9.16	$4.31 \pm 0.11$	3.2	0.5	$-5.46 \pm 0.41$
7 8		$23 \pm 1.9$	15 21	III LPV?	$3040 \pm 250$	$5.97 \pm 0.06$	$-8.55 \\ -7.79$	$3.83 \pm 0.05$ $3.50 \pm 0.06$	3.1 3.2	0.3 0.5	$\begin{array}{c} -5.15 \pm 0.40 \\ -4.09 \pm 0.41 \end{array}$
8 9		$\begin{array}{c} 22\pm1.2\\ 20\pm0.9 \end{array}$	10	III	$2710 \pm 200$ $3280 \pm 150$	$6.73 \pm 0.08 \\ 6.63 \pm 0.06$	-7.79 -7.89	$3.30 \pm 0.00$ $3.29 \pm 0.05$	3.2 2.9	0.3	$-4.09 \pm 0.41$ $-4.79 \pm 0.40$
10		$\frac{20 \pm 0.9}{23 \pm 0.9}$	12	III	$3050 \pm 150$ $3050 \pm 150$	$6.97 \pm 0.10$	-7.55	$3.47 \pm 0.10$	3.1	0.2	$-4.15 \pm 0.41$
11		$21 \pm 1.2$	12	III	$3190 \pm 180$	$6.08 \pm 0.17$	-8.44	$3.14 \pm 0.09$	3.0	0.2	$-5.24 \pm 0.43$
13		$23\pm1.0$	17	LPV	$2740\pm200$	$6.62\pm0.08$	-7.90	$3.44\pm0.07$	3.2	0.5	$-4.20\pm0.41$
14		$24\pm0.9$	5	Ι	$3470\pm150$	$6.88\pm0.06$	-7.64	$3.60\pm0.05$	2.6	0.0	$-5.04\pm0.40$
17		$21 \pm 1.3$	7	III	$3200 \pm 190$	$6.55 \pm 0.06$	-7.97	$3.20 \pm 0.05$	3.0	0.2	$-4.77 \pm 0.40$
18		$23 \pm 0.7$	19	LPV? LPV?	$2740 \pm 200$	$5.83 \pm 0.09$	-8.69	$4.36 \pm 0.09$	3.2	0.5	$-4.99 \pm 0.41$
19 23		$21 \pm 1.1 \\ 21 \pm 0.7$	38 8	III	$2560 \pm 200$ $3240 \pm 140$	$\begin{array}{c} 6.03 \pm 0.09 \\ 6.56 \pm 0.05 \end{array}$	$-8.49 \\ -7.96$	$4.11 \pm 0.09$ $3.18 \pm 0.04$	3.2 3.0	0.5 0.2	$-4.79 \pm 0.41$ $-4.76 \pm 0.40$
27		$21 \pm 0.7$ $20 \pm 1.5$	18	LPV?	$3240 \pm 140$ $2740 \pm 200$	$5.49 \pm 0.09$	-9.03	$3.53 \pm 0.09$	3.2	0.2	$-5.33 \pm 0.41$
28	IRS 11	$22 \pm 1.6$	9	III	$3130 \pm 220$	$6.18 \pm 0.09$	-8.34	$3.00 \pm 0.06$	3.0	0.3	$-5.04 \pm 0.41$
40		$22\pm2.0$	21	LPV?	$2700\pm200$	$5.83\pm0.11$	-8.69	$4.20\pm0.10$	3.2	0.5	$-4.99\pm0.41$
41		$15\pm1.7$	5	III	$3680\pm230$	$6.92\pm0.19$	-7.60	$3.22\pm0.16$	2.7	0.0	$-4.90\pm0.44$
43	ID G 10G	$17 \pm 1.9$	7	III	$3540 \pm 260$	$6.89 \pm 0.06$	-7.63	$3.18 \pm 0.05$	2.8	0.0	$-4.83 \pm 0.40$
47	IRS 12S	$15 \pm 1.1$	4	III	$3630 \pm 170$	$6.95 \pm 0.07$	-7.57	$3.00 \pm 0.06$	2.7	0.0	$-4.87 \pm 0.41$
48 50	IRS 22 IRS 12N	$21 \pm 0.7$ $20 \pm 2.1$	4 30	I LPV	$3710 \pm 140 \\ 2630 \pm 200$	$5.43 \pm 0.14$ $4.74 \pm 0.14$	$-9.09 \\ -9.78$	$2.52 \pm 0.13$ $3.79 \pm 0.14$	2.6 3.2	0.0 0.5	$-6.49 \pm 0.42$ $-6.08 \pm 0.42$
64	113 121	$20 \pm 2.1$ $23 \pm 2.1$	17	LPV?	$2030 \pm 200$ $2740 \pm 200$	$4.74 \pm 0.14$ $6.69 \pm 0.05$	-7.83	$3.79 \pm 0.14$ $3.44 \pm 0.05$	3.2	0.5	$-0.03 \pm 0.42$ $-4.13 \pm 0.40$
66	IRS 7	$25 \pm 1.8$	3	I	$3430 \pm 240$	$2.92 \pm 0.16$	-11.60	$3.48 \pm 0.09$	2.6	0.0	$-9.00 \pm 0.43$
68	IRS 14NE	$23\pm1.3$	17	LPV?	$2740\pm200$	$6.13\pm0.07$	-8.39	$3.61\pm0.05$	3.2	0.5	$-4.69\pm0.41$
69		$20\pm1.7$	14	III	$3290\pm230$	$6.03\pm0.07$	-8.49	$3.76\pm0.06$	2.9	0.2	$-5.39\pm0.41$
70		$22\pm1.2$	8	III	$3160\pm180$	$6.57\pm0.09$	-7.95	$3.32\pm0.07$	3.0	0.3	$-4.65\pm0.41$
72 <sup>h</sup>	F95B	$17 \pm 0.6$	4	III L DV/9	$3960 \pm 130$	$7.38 \pm 0.06$	-7.14	$1.67 \pm 0.04$	2.6	0.0	$-4.54 \pm 0.40$
75 79	IRS 15NE	$17 \pm 1.1 \\ 21 \pm 0.9$	18 - 1	LPV? I	$2730 \pm 200$ $3740 \pm 150$	$5.80 \pm 0.06$ $6.32 \pm 0.17$	$-8.72 \\ -8.20$	$3.16 \pm 0.05$ $3.41 \pm 0.11$	3.2 2.6	0.5 0.0	$-5.02 \pm 0.40$ $-5.60 \pm 0.43$
84		$21 \pm 0.9$ $23 \pm 1.4$	$\frac{-1}{10}$	III	$3740 \pm 130$ $3050 \pm 200$	$0.32 \pm 0.17$ $5.85 \pm 0.08$	-8.20 -8.67	$3.41 \pm 0.11$ $4.03 \pm 0.07$	3.1	0.0	$-5.27 \pm 0.41$
91	IRS 9	$18 \pm 0.9$	37	LPV	$2570 \pm 200$	$5.21 \pm 0.06$	-9.31	$3.36 \pm 0.06$	3.2	0.5	$-5.61 \pm 0.40$
96	IRS 1NE	$21\pm2.1$	14	III	$3220\pm270$	$6.52\pm0.16$	-8.00	$3.48\pm 0.14$	3.0	0.2	$-4.80\pm0.43$
97	IRS 1SE	$21\pm1.6$	14	III	$3220\pm220$	$6.78\pm0.08$	-7.74	$3.47\pm0.07$	3.0	0.2	$-4.54\pm0.41$
99	ID C AG	$22 \pm 0.7$	7	III	$3160 \pm 170$	$6.35 \pm 0.05$	-8.17	$3.41 \pm 0.04$	3.0	0.3	$-4.87 \pm 0.40$
102	IRS 28	$20 \pm 1.6$	19	LPV	$2730 \pm 200$	$5.62 \pm 0.06$	-8.90	$3.74 \pm 0.05$	3.2	0.5	$-5.20 \pm 0.40$
103 105	OSU C2	$\begin{array}{c} 23\pm1.0\\ 22\pm0.6\end{array}$	11 17	III LPV?	$3040 \pm 160 \\ 2740 \pm 200$	$6.88 \pm 0.07$ $5.65 \pm 0.05$	-7.64 - 8.87	$3.22 \pm 0.06$ $3.26 \pm 0.04$	3.1 3.2	0.3 0.5	$-4.24 \pm 0.41$ $-5.17 \pm 0.40$
107		$19 \pm 1.2$	5	III	$3330 \pm 180$	$6.84 \pm 0.05$	-7.68	$3.26 \pm 0.04$ $3.26 \pm 0.04$	2.9	0.2	$-4.58 \pm 0.40$
108	IRS 19	$18 \pm 1.7$	4	Ι	$3910\pm230$	$4.64\pm0.11$	-9.88	$3.50\pm0.10$	2.6	0.0	$-7.28\pm0.41$
109	IRS 18	$20\pm1.6$	4	Ι	$3760\pm220$	$6.18\pm0.05$	-8.34	$3.32\pm0.04$	2.6	0.0	$-5.74\pm0.40$
110		$21\pm0.9$	15	III	$3200\pm150$	$6.22\pm0.07$	-8.30	$3.85\pm0.05$	3.2	0.5	$-4.60\pm0.41$
112	OSUC 3	$24 \pm 2.0$	12	III	$2980 \pm 270$	$6.48 \pm 0.06$	-8.04	$4.25 \pm 0.04$	3.1	0.4	$-4.54 \pm 0.40$
113 114		$20 \pm 1.0$ $24 \pm 1.1$	7 12	III III	$3310 \pm 160 \\ 2970 \pm 170$	$\begin{array}{c} 6.50 \pm 0.12 \\ 5.65 \pm 0.06 \end{array}$	$-8.02 \\ -8.87$	$\begin{array}{c} 3.51 \pm 0.11 \\ 2.95 \pm 0.04 \end{array}$	2.9 3.2	0.2	$\begin{array}{c} -4.92 \pm 0.42 \\ -5.27 \pm 0.40 \end{array}$
114		$24 \pm 1.1$ $22 \pm 1.1$	7	III III	$2970 \pm 170$ $3140 \pm 170$	$5.03 \pm 0.06$ $6.83 \pm 0.06$	-8.87 -7.69	$2.93 \pm 0.04$ $3.56 \pm 0.05$	3.0	0.4 0.3	$-3.27 \pm 0.40$ $-4.39 \pm 0.40$
120		$22 \pm 1.1$ $23 \pm 1.0$	12	III	$3080 \pm 160$	$6.73 \pm 0.06$	-7.79	$3.75 \pm 0.06$	3.1	0.3	$-4.39 \pm 0.40$
121		$24 \pm 1.1$	14	III	$3000 \pm 170$	$5.68 \pm 0.05$	-8.84	$3.83 \pm 0.04$	3.1	0.4	$-5.34 \pm 0.40$
123		$21\pm1.1$	8	III	$3170\pm170$	$6.37\pm0.05$	-8.15	$3.67\pm0.04$	3.0	0.3	$-4.85\pm0.40$
124		$21\pm1.2$	3	Ι	$3720\pm180$	$6.42\pm0.05$	-8.10	$2.68\pm0.04$	2.6	0.0	$-5.50\pm0.40$
126	ID C A I	$22 \pm 1.0$	9	III	$3110 \pm 160$	$5.99 \pm 0.06$	-8.53	$3.70 \pm 0.05$	3.1	0.3	$-5.13 \pm 0.40$
128	IRS 24	$22 \pm 2.6$	26	LPV	$2660 \pm 200$	$4.92 \pm 0.09$	-9.60	$3.34 \pm 0.07$	3.2	0.5	$-5.90 \pm 0.41$
129 131		$\begin{array}{c} 20\pm0.7\\ 22\pm0.9\end{array}$	2 15	I III	$3750 \pm 130 \\ 3090 \pm 150$	$\begin{array}{c} 6.59 \pm 0.06 \\ 6.26 \pm 0.05 \end{array}$	$-7.93 \\ -8.26$	$2.79 \pm 0.05$ $2.65 \pm 0.04$	2.6 3.1	0.0 0.3	$-5.33 \pm 0.40$ $-4.86 \pm 0.40$
132		$\begin{array}{c} 22 \pm 0.9 \\ 23 \pm 0.8 \end{array}$	11	III	$3050 \pm 130$ $3050 \pm 140$	$0.20 \pm 0.03$ $6.98 \pm 0.05$	-7.54	$2.03 \pm 0.04$ $3.27 \pm 0.04$	3.1	0.3	$-4.30 \pm 0.40$ $-4.14 \pm 0.40$
133		$19 \pm 1.1$	14	III	$3340 \pm 170$	$5.72 \pm 0.05$	-8.80	$3.48 \pm 0.04$	2.9	0.2	$-5.70 \pm 0.40$
135	OSUC 4	$24 \pm 1.9$	15	III	$2940\pm260$	$7.00 \pm 0.08$	-7.52	$3.67 \pm 0.07$	3.2	0.4	$-3.92 \pm 0.41$
136	IRS 23	$18\pm2.1$	35	LPV?	$2580\pm200$	$5.18\pm0.14$	-9.34	$3.44\pm0.13$	3.2	0.5	$-5.64\pm0.42$
140		$23 \pm 0.5$	12	III	$3050 \pm 120$	$6.50\pm0.05$	-8.02	$2.93\pm0.04$	3.1	0.3	$-4.62 \pm 0.40$
145		$18 \pm 1.3$	23	LPV?	$2690 \pm 200$	$5.91 \pm 0.09$	-8.61	$4.54 \pm 0.09$	3.2	0.5	$-4.91 \pm 0.41$
146 148		$19 \pm 2.2$ 21 ± 1.0	13	III	$3330 \pm 290$ $3170 \pm 160$	$6.54 \pm 0.07$ $6.97 \pm 0.05$	-7.98	$3.91 \pm 0.06$ $3.55 \pm 0.09$	2.9	0.2	$-4.88 \pm 0.41$ $4.25 \pm 0.40$
140		$21 \pm 1.0$	7	III	$3170\pm160$	$6.97\pm0.05$	-7.55	$3.55\pm0.09$	3.0	0.3	$-4.25\pm0.40$

TABLE 6—Continued

Number	Name	CO (%)	H <sub>2</sub> O (%)	Luminosity Class <sup>a</sup>	$T_{\rm eff}{}^{\rm b}$	$K_0^{c}$	$M_K^{\rm d}$	$A_K^{\rm c}$	$\mathrm{BC}_{K}^{e}$	$\Delta A_K^{\rm f}$	$M_{ m bol}{}^{ m g}$
150		$22\pm1.9$	8	III	$3130\pm250$	$6.97\pm0.12$	-7.55	$3.84\pm0.11$	3.0	0.3	$-4.25\pm0.42$
151		$25\pm0.9$	3	Ι	$3400\pm150$	$5.73\pm0.16$	-8.79	$3.30\pm0.15$	2.6	0.0	$-6.19\pm0.43$
155		$24\pm1.5$	11	III	$2930\pm210$	$6.79\pm0.09$	-7.73	$3.82\pm0.08$	3.2	0.4	$-4.13\pm0.41$
162		$26\pm0.8$	3	Ι	$3360\pm150$	$5.92\pm0.17$	-8.60	$4.93\pm0.00$	2.6	0.0	$-6.00\pm0.43$
167		$17\pm2.8$	31	LPV?	$2620\pm200$	$5.66\pm0.10$	-8.86	$5.19\pm0.09$	3.2	0.5	$-5.16\pm0.41$
168		$25\pm0.9$	7	Ι	$3430\pm150$	$6.92\pm0.07$	-7.60	$3.77\pm0.06$	2.6	0.0	$-5.00\pm0.41$
170		$22 \pm 1.0$	5	Ι	$3660\pm160$	$5.83\pm0.12$	-8.69	$4.68\pm0.11$	2.6	0.0	$-6.09\pm0.42$
173		$21\pm1.0$	2	Ι	$3740 \pm 160$	$5.62\pm0.14$	-8.90	$5.33\pm0.13$	2.6	0.0	$-6.30\pm0.42$
176		$25\pm0.9$	12	III	$2880 \pm 150$	$6.79\pm0.09$	-7.73	$3.88\pm0.08$	3.2	0.4	$-4.13\pm0.41$
180		$25\pm1.0$	14	III	$2910\pm160$	$4.99\pm0.10$	-9.53	$5.77\pm0.10$	3.2	0.4	$-5.93\pm0.41$
182		$21 \pm 1.0$	10	III	$3180\pm160$	$6.65\pm0.07$	-7.87	$3.87\pm0.06$	3.0	0.2	$-4.67\pm0.41$
186		$24\pm1.6$	-3	Ι	$3500\pm220$	$6.40\pm0.16$	-8.12	$4.54\pm0.15$	2.6	0.0	$-5.52\pm0.43$
197		$21 \pm 1.4$	1	Ι	$3680\pm200$	$6.82\pm0.06$	-7.70	$3.77\pm0.05$	2.6	0.0	$-5.10\pm0.40$
198		$25\pm1.3$	10	III	$2860\pm190$	$6.93\pm0.06$	-7.59	$3.82\pm0.06$	3.2	0.4	$-3.99\pm0.40$
211		$25\pm1.1$	14	III	$2930\pm170$	$6.70\pm0.06$	-7.82	$3.85\pm0.05$	3.2	0.4	$-4.22\pm0.40$
214		$23 \pm 1.0$	11	III	$3080 \pm 160$	$4.91\pm0.13$	-9.61	$5.69\pm0.12$	3.1	0.3	$-6.21\pm0.42$
	VR 5-7	$23\pm0.9$	6	Ι	$3570\pm150$		-10.22	$3.10\pm\ldots$	2.6	0.0	$-7.62\pm0.40$

NOTE.—Indexes are percent absorption defined as  $100 \times (1 - \text{flux/cont})$  where "flux" and "cont" are the integrated fluxes in 0.015  $\mu$ m bands in the spectra. For CO, the "flux" band is centered at 2.302  $\mu$ m and the "cont" band is taken at 2.284  $\mu$ m. For H<sub>2</sub>O, "flux" is for a band centered at 2.0675  $\mu$ m and "cont" is the same, but derived from a fit to the spectral continuum; see text. The uncertainty in the measured H<sub>2</sub>O strength (3%) is taken from the scatter in the measurements for comparison star IIIs and Is (see text). Table 6 is also available in machine-readable form in the electronic edition of the *Astrophysical Journal*.

<sup>a</sup> The luminosity class is set by the CO index and the H<sub>2</sub>O index (see text).

<sup>b</sup>  $T_{\rm eff}$  is derived from the CO index using the relationship of CO vs.  $T_{\rm eff}$  derived for the comparison stars. See the discussion in the text and Fig. 2.

 $^{c}K_{0}$  (=*K*-*A<sub>K</sub>*) and *A<sub>K</sub>* taken from Blum et al. 1996a). For stars 1, 2, 6, 11, 40, 148, 162, and 170, *A<sub>K</sub>* were derived in the same way, but using new unpublished *H* magnitudes as given in Table 2.

<sup>d</sup>  $M_K$  is derived from  $K_0$  and a distance modulus of 14.52 (8 kpc; Reid 1993).

<sup>e</sup> Bolometric correction to the  $K_0$  magnitudes. The correction is a linear function of  $T_{\text{eff}}$  for giants, and a constant for LPVs and supergiants (see text and Blum et al. 1996b).

<sup>f</sup> Correction to  $A_K$  to account for the range of intrinsic color in the GC stars.  $\Delta A_K = A_K$  (Blum et al. 1996a  $-A_K$  [corrected]; see text).

 $^{g}M_{bol} = K_0 + BC_K + \Delta A_K - DM$ , where DM is the distance modulus to the GC (14.52) for 8 kpc Reid 1993.

<sup>h</sup> Star 72 is included because it was observed and analyzed at high spectral resolution by Ramírez et al. 2000. This star was not used in the star formation history calculation because its  $K_0$  is below the cutoff adopted; see § 3.3.

absorption between the H and K bands and also the measured quantities for CO and  $H_2O$  as described in § 3.2.2: stars with strong CO compared with the comparison I stars can have slightly stronger H<sub>2</sub>O and still be classified as I, while stars with less CO must have very low measured H<sub>2</sub>O to be classified as I. The GC luminosity class assignments (based only on the H<sub>2</sub>O and CO absorption) are consistent with the comparison star luminosities. If we plot the GC and comparison star  $H_2O$  versus  $M_{bol}$ , we see the GC stars fall along two broad tracks delineated by the comparison stars (Fig. 11). To the lower left in this diagram are the warm comparison star IIIs. Going vertically at low H<sub>2</sub>O are the supergiants. The comparison star LPVs run generally to the right to larger H<sub>2</sub>O absorption and to gradually higher luminosity. The GC Is follow on or near the comparison I track, while the cooler GC IIIs essentially fill the region between the warmer comparison IIIs and the cooler comparison LPVs as is expected for cooler IIIs high on the AGB. Several GC stars exist where these broad tracks overlap. In each case, the GC star  $M_{bol}$  is consistent with the comparison luminosities for the corresponding class. The average H<sub>2</sub>O strength for GC stars classified as giants (III in Table 6) was  $10\% \pm 3\%$ . For supergiants, the average was  $3\% \pm 2\%$ . The GC stars are also well separated between I and III classes in the H<sub>2</sub>O versus CO plot (Fig. 6) and are consistent with the positions of the comparison stars.

In summary, we used criteria on the appearance and measured strength of H<sub>2</sub>O and measured CO strength to assign

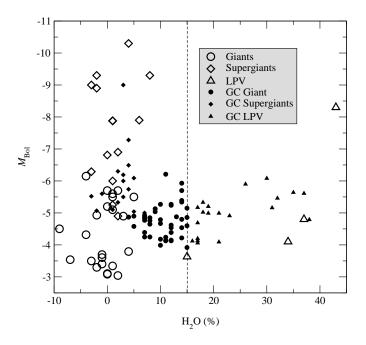


FIG. 11.—Comparison and GC star  $M_{bol}$  vs. H<sub>2</sub>O strength plot. The comparison stars are shown as open symbols, while GC stars are plotted with filled symbols. The assigned GC luminosity classes are consistent with the comparison star luminosity ranges. The plot shows that we have observed later type giants in the GC than are represented in the warmer comparison III sample, and these GC IIIs lie between the comparison III sample, not along the comparison I track, which runs vertically in this diagram (see text for details; see also Fig. 6).

all GC luminosity classes. We then showed that the GC assignments, when quantitatively based on CO, H<sub>2</sub>O, and  $M_{bol}$ , reproduced the classifications by eye and are well matched to the parameters of the comparison stars. The GC giants lie at or below the comparison giant luminosities, while supergiants always lie above the minimum comparison supergiant  $M_{bol}$ . As noted in § 3.2.2, the comparison star sample does not include giants later than type M6 as a consequence of our selection based on  $M_{bol}$  and  $T_{eff}$ : such stars are rare in the solar neighborhood. Nevertheless, the criteria used here for CO and H<sub>2</sub>O indicate that such stars exist in larger numbers in the GC, as would be expected in this dense stellar environment.

### 4. RESULTS: DETERMINATION OF THE STAR FORMATION HISTORY

### 4.1. *The H-R Diagram*

The spectral indexes developed in the preceding section allow us to derive bolometric magnitudes and  $T_{\rm eff}$  for the GC stars.  $T_{\rm eff}$  follows directly from the measured CO index once the luminosity class is chosen (see § 3.3).  $T_{\text{eff}}$  values are given in Table 6. The uncertainty in  $T_{\rm eff}$  is derived by propagating the uncertainty in the measured CO strength through the CO versus  $T_{\rm eff}$  relation. This gives an average uncertainty, for all giants and supergiants in the sample, of 184 K. For the 11 stars in common with the sample given by BSD96 (not including the LPVs), we find  $\Delta T_{\text{eff}} = 32 \pm 156$ K. In addition, our value of  $T_{\rm eff}$  for VR 5-7 differs from that calculated by Ramírez et al. (2000) by only 74 K. The Mbol for a given GC star follows from the intrinsic  $K_0$  magnitude given in Table 2, the distance modulus, and a bolometric correction to the  $K_0$  magnitude. The uncertainty in  $M_{bol}$  is the sum in quadrature of the photometric uncertainty (BSD96) and an 0.4 mag uncertainty due to the uncertainty in the interstellar extinction law (Mathis 1990). The distance modulus is taken as 14.52 (8 kpc; Reid 1993).

The bolometric correction  $(BC_K)$  is derived from the literature for the different luminosity classes given in Table 6. For supergiants and LPVs (Table 6), the BC<sub>K</sub> is the same as given by BSD96: BC<sub>K</sub> = 2.6 for supergiants and 3.2 for LPVs. For the giants, we improve on the work of BSD96 by considering a BC<sub>K</sub>, which is a function of  $T_{\text{eff}}$ . Using the BC<sub>K</sub> as a function of J-K given by Frogel & Whitford (1987), the mean J-K of giants as a function of spectral type from Frogel et al. (1978), the spectral type versus  $T_{\text{eff}}$  from Ramírez et al. (2000) and Dyck et al. (1998), and the  $T_{\text{eff}}$ given in Table 6, we derive BC<sub>K</sub> as a linear function of  $T_{\text{eff}}$ [BC<sub>K</sub> = 2.6 -  $(T_{\text{eff}} - 3800)/1500$ ]. These BC<sub>K</sub> range from 2.8 to 3.2 for the warmest and coolest GC giants.

A second correction was also applied to the GC stellar  $M_{bol}$  values, which relates to the  $A_K$  for each star. BSD96 used mean  $(J-H)_0$  and  $(H-K)_0$  colors to estimate the  $A_K$  for each GC star. As discussed by them, this will lead to values of  $A_K$  that are systematically too high for stars that are intrinsically redder than these colors, and the opposite will be true for stars that are intrinsically bluer. The individual spectrum for each star now allows us to improve upon the corresponding  $A_K$  estimate. Using the same color and  $T_{\text{eff}}$  data described in the preceding paragraph, we estimated a correction (linear with  $T_{\text{eff}}$ ) to the intrinsic colors, and hence to  $A_K$ , for each star. This correction varied from 0.0 to 0.5 mag depending on  $T_{\text{eff}}$  (see Table 6) and was applied in



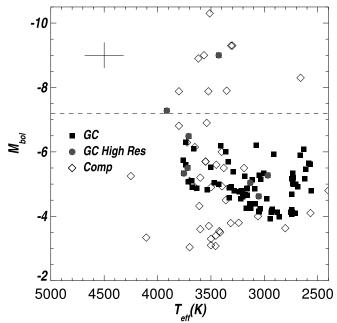


FIG. 12.—H-R diagram for the GC stars (shown as filled squares with typical uncertainty given by the error bar in the upper left-hand corner) and comparison stars ("Comp," *open diamonds*). The GC stars analyzed at high spectral resolution by Carr et al. (2000) and Ramírez et al. (2000) are plotted as filled circles. [See the electronic edition of the Journal for a color version of this figure.]

the sense that makes  $A_K$  less for each star [ $\Delta A_K = -0.11 - (T_{\text{eff}} - 3800)/1730$ ].

Figure 12 shows the H-R diagram for all the GC stars and comparison stars listed in Tables 3 and 6. This figure illustrates that the comparison stars span the same range in  $M_{\rm bol}$  and  $T_{\rm eff}$  as the GC stars do.

Figure 13 shows the H-R diagram for all the GC stars listed in Table 6, with isochrones from Bertelli et al. (1994) and Girardi et al. (2000) overplotted. These isochrones vary in age from 10 Myr to 12 Gyr, with [Fe/H] = 0.0 for all ages. Figure 13 shows that our GC sample spans a wide range in age. As can be seen in the figure, all of the GC giants (those labeled III in Table 6) are AGB stars. They are too luminous to be first ascent giants, which is a consequence of our selection criteria.

Figure 14 again shows the H-R diagram for the GC, but this time overplotted with Bertelli et al. (1994) and Girardi et al. (2000) isochrones having [Fe/H] = -0.2 for all ages. The figure shows that lower metallicity isochrones do not extend to cool enough  $T_{\rm eff}$  to match the GC H-R diagram.

To give a general feel for the SFH represented by Figure 13, Figure 15 shows the H-R diagram against a simulation representing a constant SFH. The observed points appear to span the model parameter space. That there are relatively few model stars at high and low luminosities suggests that higher star formation rates (SFRs) in both the distant and recent past are needed to fit the data. There is an intermediate-age component at  $T_{\rm eff} = 3300$  K,  $M_{\rm bol} = -5.0$  whose position and extent matches the model well. The tightness of this feature suggests that our errors may be overestimated relative to the internal scatter in the points. This is particularly true for the errors in  $M_{\rm bol}$ , which are dominated by the *systematic* uncertainty in the interstellar extinction law.

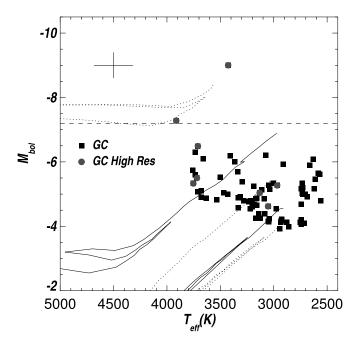


FIG. 13.—GC stars plotted as in Fig. 12, but with [Fe/H] = 0.0 isochrones plotted as well. The isochrones are from Bertelli et al. (1994) for age less than 100 Myr and from Girardi et al. (2000) otherwise. Isochrones are plotted for ages of 10 and 100 Myr and 1, 5, and 12 Gyr. The models have [Fe/H] = 0.0 for all ages, and these appear to better represent the data than models with lower metallicity at older ages ([Fe/H] = -0.2; see Fig. 12c and text). Neither set of isochrones reaches the coolest stars (LPV candidates with  $T_{\rm eff}$  < 2800 K), but the [Fe/H] = 0.0 isochrones extend to cooler temperatures and thus fit more Galactic stars than the [Fe/H] = -0.2 isochrones. Comparison to the isochrones shows that all the GC stars classified as giants (Table 6, III) are AGB stars; they are too luminous to be first-ascent giants. This is a consequence of the selection criteria. The horizontal line segment at  $M_{\text{bol}} = -7.2$  in each panel indicates the approximate observed luminosity above which only supergiants lie (Blum et al. 1996b). [See the electronic edition of the Journal for a color version of this figure.]

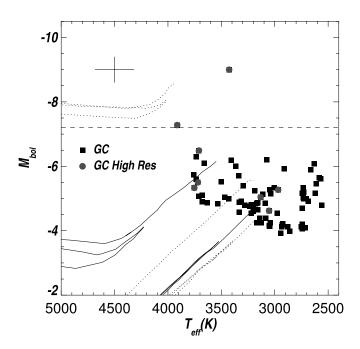


FIG. 14.—Same as for Fig. 12*b*, but with [Fe/H] = -0.2 isochrones plotted. [See the electronic edition of the Journal for a color version of this figure.]

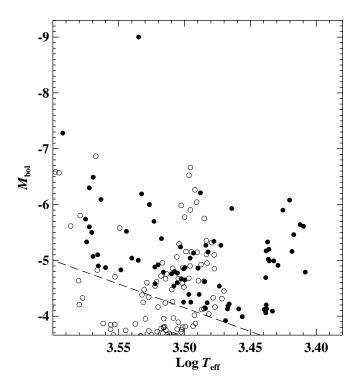


FIG. 15.—H-R diagram for the GC stars (*filled circles*) compared with a model with constant SFR (*open circles*). This figure demonstrates that the models cover the same parameter space as the GC data, including the pronounced intermediate-age feature at  $T_{\rm eff} \approx 3200$  K  $[\log_{10}(T_{\rm eff}) = 3.50]$ ,  $M_{\rm bol} \approx -5.0$ , except for the coolest GC stars  $[\log_{10}(T_{\rm eff}) > 3.45]$  (see text). The model points include objects below our observed magnitude cutoff (the region below the dashed line).

One exception in the model coverage appears to be the coolest AGB stars. Both the Bertelli et al. (1994) and Girardi et al. (2000) models fail to reach the coolest observed  $T_{\rm eff}$ . This is true of both the GC and comparison stars. Thus, only stars with  $T_{\rm eff} > 2800$  K (allowing for the errors in  $T_{\rm eff}$  for the observed stars) were directly included in the basic SFH calculations discussed below. The models attempt to follow evolution along the AGB in a simplified way (Girardi et al. 2000), giving a typical or average locus in the H-R diagram, but real AGB stars pulsate with periods of hundreds of days. These pulsations result in excursions in the H-R diagram of 500-1000 K (Lançon & Mouhcine 2002). As discussed by Lançon & Mouhcine (2002), the effect of pulsations is thus to widen the AGB. For stars that experience excursions within the model temperature range  $(T_{\rm eff} > 2800 \text{ K})$ , the pulsations will be randomly phased, so that differences between the observed location in the H-R diagram and the actual isochrone to which a star would otherwise be associated with are canceled out. For stars that are cooler than the models, we assume that they should be associated with an isochrone inside the model H-R diagram space. Assuming that these stars represent the same fraction of initial mass independent of which isochrone they are really associated with, we simply scale the total star formation by their number. In this case, there are 20 such cool stars (out of 78 used to calculate the SFH); thus we will take the total SFRs to be 1.3 times the amount given by our fit results.

We also did not include IRS 7. This star belongs to the youngest nuclear star burst that is accounted for by the Krabbe et al. (1995) model.

BSD96 estimated ages for a number of the coolest, most luminous stars, which are also shown in Figures 12, 13, and 14. The effect of a reduced  $A_K$  for these stars (compared with BSD96) has lowered their luminosity. This tends to increase the estimated age. In particular, for IRS 9, 12N, 23, 24, and 28 (see Table 6), BSD96 estimated the mass for these stars from their luminosity. This resulted in correspondingly young ages (100–200 Myr) from model isochrones. The reduced luminosity determined in the present work (due to the decrease in  $A_K$ ) and comparison to different model isochrones as adopted herein (Fig. 13) suggest a somewhat older age for these stars (roughly 500 Myr to a few gigayears). Given the preceding discussion, it is not possible to estimate the age of any particular star to great precision.

## 4.2. Star Formation History Calculation

We have used the results of Figure 12 to derive the SFH implied by these observations of the GC cool stars. The calculation was carried out using Olsen's (1999) implementation of the method described by Dolphin (1997), with some modifications. In brief, we constructed a set of models describing the expected distribution of stars in the H-R diagram within specified age bins, assuming a particular metallicity, slope of the initial mass function (IMF), and constant SFR ( $1 M_{\odot} \text{ yr}^{-1}$ ) within the bin and accounting for observational errors and incompleteness. We chose the best model SFH for the GC by fitting the observed data to a linear combination of the star formation within these bins. This fit was determined through the maximum likelihood analysis described below.

## 4.2.1. Model Parameters

Two choices of sets of age bins and metallicities were used. Models had either four age bins (model A: 10-100 Myr, 100 Myr to 1 Gyr, 1-5 Gyr, and 5-12 Gyr) or three age bins (model 1: 10-50 Myr, 50 Myr to 3 Gyr, and 3-12 Gyr). For both sets of age bins, models were run with all stars at solar [Fe/H] (model A with four age bins and model 1 with three age bins) and then again with solar [Fe/H] for the

younger stars and [Fe/H] = -0.2 in the oldest bin (model B with 4 age bins, model 2 with three age bins). To explore the effect of the IMF on the derived SFH, we also computed models with a power-law slope  $\alpha = -2.35$  (Salpeter 1955) and with a slightly flatter slope  $\alpha = -2$  for stars with masses greater than 1  $M_{\odot}$ . The matrix of models is given in Table 7.

The SFH is also constrained by the amount of mass inferred from dynamical models in the GC. Recent models have been computed by McGinn et al. (1989), Sellgren et al. (1990), Haller et al. (1996), Genzel et al. (1996), Saha, Bicknell, & McGregor (1996), Ghez et al. (1998), and Genzel et al. (2000). The most detailed models, which include proper motion and radial velocities (Ghez et al. 1998; Genzel et al. 2000), require a concentrated mass of approximately  $3 \times 10^6 M_{\odot}$  (presumably a supermassive black hole) that dominates the distribution within less than 1 pc. These models also predict approximately  $6 \times 10^6 M_{\odot}$ total mass enclosed within a radius of 2 pc (for  $R_0 = 8$  kpc). Finally, new observations of single stars orbiting the black hole (Schödel et al. 2002; Ghez et al. 2003) require slightly higher black hole masses of  $\approx 4 \pm 1 \times 10^6 M_{\odot}$ .

If the black hole itself was not built up from stars and stellar remnants in the formation of the nucleus, this leaves roughly  $1-3 \times 10^6 M_{\odot}$  in distributed mass that arises from the luminous stellar population and the accumulated stellar remnants integrated over the lifetime of the nucleus.

We built in the constraint on dynamical mass by limiting the low-mass end of the IMF. The constraint is taken to be that the black hole mass plus the stellar mass formed minus the mass loss from stellar winds equals the total dynamical mass. We discuss the effects of mass loss on the present-day mass below (§ 5). This is the simplest choice since our H-R diagram is not sensitive to stars below about 1  $M_{\odot}$ . We discuss a possible cause for this low-mass cutoff below in § 5. We began with the same IMF as Miralda-Escudé & Gould (2000), including their corrections, with masses spanning the range 0.08–120  $M_{\odot}$ . This IMF is based on Salpeter's mass function for the more massive stars (Salpeter 1955); i.e.,  $dN/dM \propto m^{-\alpha}$  with  $\alpha = -2.35$ . For lower masses, the IMF is flatter, as determined from bulge star counts at 6°

TABLE 7Summary of Models Fitted

Model	$\chi^2_\lambda$	$P_{\lambda}$ (%)	$M_T \ (10^6  M_\odot)$	$\sigma_{M_T}$	Age (Gyr)	[Fe/H]	SFR $(10^{-4} M_{\odot} \mathrm{yr}^{-1})$	$\sigma_{ m SFR}$
A	231.9	58.9	9.9	3.0	0.01-0.10	0.0	33.3	9.7
					0.10 - 1.00	0.0	2.9	1.9
					1.00 - 5.00	0.0	4.5	1.6
					5.00-12.00	0.0	10.8	4.1
A, IMF slope varied	232.4	54.8	16.4	4.3	0.01 - 0.10	0.0	28.3	8.5
					0.10 - 1.00	0.0	3.4	2.4
					1.00 - 5.00	0.0	7.3	2.4
					5.00-12.00	0.0	18.4	7.3
B	233.3	13.7	16.1	4.1	0.01 - 0.10	0.0	30.2	9.7
					0.10 - 1.00	0.0	3.8	1.8
					1.00 - 5.00	0.0	0.00	1.8
					5.00-12.00	-0.2	22.1	5.7
1	238.1	11.0	9.07	2.4	0.01-0.50	0.0	10.7	2.4
					0.50 - 3.00	0.0	3.1	1.3
					3.00-12.00	0.0	8.7	2.7
2	238.5	2.2	12.1	1.8	0.01-0.50	0.0	12.3	2.4
					0.50-3.00	0.0	0.00	0.10
					3.00-12.00	-0.2	12.8	2.0

projected distance from the GC (Zoccali et al. 2000; exponent -2.0 for  $1 M_{\odot} > M > 0.7M_{\odot}$  and -1.65 for  $M < 0.7M_{\odot}$ ). For the lowest masses, there is a correction to the exponent due to binaries (Miralda-Escudé & Gould 2000). In the end, we found that cutting off the mass function at 0.7  $M_{\odot}$  resulted in models that are consistent with the present-day dynamical mass.

# 4.2.2. Model Calculations

Within each age bin, we calculated the distribution of stars over a fine grid in  $T_{\text{eff}}$  and  $M_{\text{bol}}$  directly from 20 isochrones spaced linearly in age. We used the Girardi et al. (2000) isochrones to construct the models with ages  $\geq 63$  Myr and the Bertelli et al. (1994) isochrones for younger models; we interpolated the isochrones in age and metallicity following the procedure described in Olsen (1999). We then convolved the grid with a two-dimensional Gaussian kernel having a shape representing the typical errors in  $T_{\text{eff}}$  and  $M_{\text{bol}}$  and multiplied the grid with a surface representing our estimate of the completeness as a function of  $T_{\text{eff}}$  and  $M_{\text{bol}}$  (i.e., the models were transformed to the observational plane).

After gridding the observed H-R diagram to the same resolution as the models and selecting an area to exclude the likely Mira variable stars, we searched for the linear combination of models producing the highest likelihood of describing the observed distribution of stars. This search was done by using the Numerical Recipes routine amoeba (Press et al. 1992) to minimize the Poisson maximum likelihood parameter  $\chi_{\lambda}^2 = 2 \sum_i m_i - n_i + n_i \ln(n_i/m_i)$  (e.g., Mighell 1999), where  $m_i$  is the number of stars predicted by the model in the *i*th bin of the H-R diagram and  $n_i$  is the number of observed stars in the bin. The virtues of this parameter are discussed extensively by Dolphin (2002)-the most important of which is that it minimizes to the correct solution when presented with a data set that sparsely samples the range of possible measurements (i.e., the Poisson regime), which the standard  $\chi^2$  does not.

### 4.2.3. Uncertainties and Goodness of Fit

The size of the errors in  $M_{bol}$  and  $T_{eff}$ , the size of our data set, and the fact that we are studying only the most luminous members of the GC population impose some limitations on our ability to discriminate model SFHs. First, we selected the age bins so as to roughly divide the H-R diagram into equal areas and to distinctly separate the footprints of adjacent age bins, given our errors. Because the isochrones bunch together at larger ages, the age bins necessarily grow wider with age, with a corresponding decrease in age resolution in the derived star formation histories. Second, we specified the age-metallicity relation in advance. While the maximum likelihood procedure described above in principle has the ability to *derive* the metallicity distribution and agemetallicity relation directly from the data, the degeneracy between age and metallicity in this region of the H-R diagram is severe. This difficulty is compounded by our sizable errors in  $M_{\rm bol}$  and  $T_{\rm eff}$  and the relatively small size of our sample. Finally, as mentioned above, we assumed the IMF, since our data do not span a large enough range in mass at a given age to allow it to be a free parameter.

For each model in Table 7, we calculated the uncertainties in the derived SFRs through Monte Carlo simulations. We produced 100 Monte Carlo samples, each containing 59 stars, by drawing randomly from the *observed* data set while allowing any particular star to be drawn any number of times (a technique referred to as "bootstrapping"). For each sample, we then derived the star formation history just as was done for the original data set. The uncertainty in a given SFR ( $\sigma_{SFR}$ ) reported in Table 7 is the 1  $\sigma$  standard deviation of the distribution of SFRs in the corresponding age bin for the 100 Monte Carlo samples.

To address the separate question of whether the data are a likely representation of the models listed in Table 7, we ran a different set of Monte Carlo simulations. For these simulations, we produced 10,000 samples, each containing 59 stars drawn randomly from the *fitted* models (not from the data). We then assembled the distribution of minimum values of the  $\chi^2_{\lambda}$  parameter by refitting the model to each Monte Carlo sample. The column labeled  $P_{\lambda}$  in Table 7 shows the percentage of runs that had higher  $\chi^2_{\lambda}$  when fitting the Monte Carlo sample to the model than that obtained by a fit of the data to the model. Thus, small values of  $P_{\lambda}$  represent poorer fits of the models. This is so because Monte Carlo samples drawn from the "right" model should, on average, produce as many fits with  $\chi^2_{\lambda}$  above as below that for the fit to the data: values near 50% are achieved by the average data set drawn randomly from the model probability distribution.

### 4.2.4. Model Results

Examining Table 7, we find that the models with only three age bins are significantly worse than those with four and so are not discussed further. Model A, in which we assume solar metallicity throughout, exclude the coolest stars and the most luminous one, and account for the uncertainty in the extinction law, fits the data as well as it does the average data set drawn randomly from the model. Model B, which is identical to model A except that we assume [Fe/H] = -0.2 for ages greater than 5 Gyr, is an unlikely fit to the data.

As suggested by Figures 13 and 14, age and metallicity are difficult to disentangle using only the tip of the AGB. However, our models do produce better fits to an SFH with purely solar [Fe/H] from 0.01–12 Gyr. This may be understood through consideration of Figure 16. The left-hand panel shows model A compared with the data, while the right-hand panel shows model B. The colors cyan, magenta, vellow, and gray represent the model distributions scaled to the best-fit SFR. Darker regions indicate higher number density of stars. In the case of model A, the footprints of the different age bins align nicely with concentrations of observed data points. In model B, the oldest age bin moves to higher temperature and luminosity, becoming degenerate with the solar metallicity and younger stars and forcing it to fit a larger number of points. As a result, model B does not fit the coolest low-luminosity stars well; indeed, these stars are difficult to fit with anything other than a greater than 12 Gyr old solar metallicity model. However, greater than 12 Gyr old stars with  $[Fe/H] \leq -0.6$  are less luminous than our  $T_{\text{eff}}$ -dependent  $M_{\text{bol}}$  limit. This means that our SFH may not account for a potential very old (>12 Gyr), metal-poor component that might have fewer luminous AGB stars.

The resulting SFH for models A and B, showing the SFR in age bins of 10–100 Myr, 100 Myr to 1 Gyr, 1–5 Gyr, and 5–12 Gyr, are given in Figures 17 and 18, respectively. The

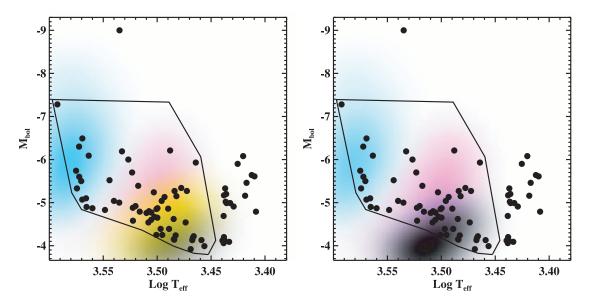
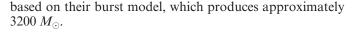


FIG. 16.—Comparison of the best-fitting star formation history with solar metallicity throughout (*left-hand panel*, model A, Table 7) and with [Fe/H] = -0.2 for ages  $\geq 5$  Gyr (*right-hand panel*, model B, Table 7). Darker regions represent higher number density. Cyan corresponds to ages between 10 and 100 Myr, magenta to 100 Myr to 1 Gyr, yellow to 1–5 Gyr, and gray to 5–12 Gyr. The purely solar metallicity model fits the data better (see text). Model B (*right*) has a best-fit solution with no star formation in the third (1–5 Gyr) bin (this is why no yellow region appears; see Fig. 18). However, a low-metallicity ( $[Fe/H] \leq -0.6$ ), very old (>12 Gyr) component would not be detected by our sample. The dark polygons in each panel represent the area of the model and observational parameter space used in the fits. The coolest stars are not accounted for by the models and so were not used in the fits (i.e., those stars outside the polygon); neither was the brightest star, IRS 7 (age < 10 Myr).

total initial mass represented by these SFHs are shown in Figures 19 and 20, respectively. In Figure 17 we have also plotted a point corresponding to the model of Krabbe et al. (1995) that we take as the average SFR over the last 10 Myr



## 5. DISCUSSION

A number of investigators have discussed the SFH in the GC. Lebofsky et al. (1982b) used the presence of young M supergiants to infer a recent ( $\leq 10$  Myr) burst of star formation. Sellgren et al. (1987) later reclassified

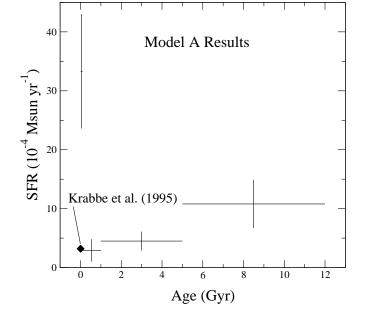


FIG. 17.—SFH for the GC (see Table 7). The crosses represent the results to the SFH fits to the H-R diagram (Fig. 13) for model A with solar [Fe/H] throughout. Model A provides the best fit (see text). The age bins correspond to the horizontal width of the crosses and are 10–100 Myr, 100 Myr to 1 Gyr, 1–5 Gyr, and 5–12 Gyr. The vertical height of each cross is the 1  $\sigma$  error in the SFR for the respective bin. The filled diamond represents the starburst model from Krabbe et al. (1995) averaged over 10 Myr and is placed at 5 Myr along the age axis (i.e., it is the youngest point in the plot).

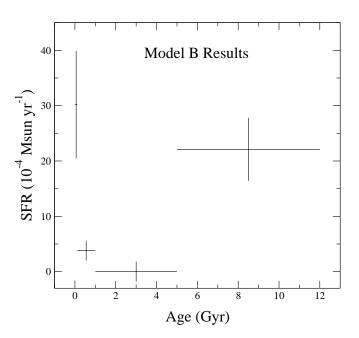


FIG. 18.—Same as Fig. 17, but for model B with [Fe/H] = -0.2 in the oldest age bins. Model A (Fig. 15*a*) provides a better fit (see text).

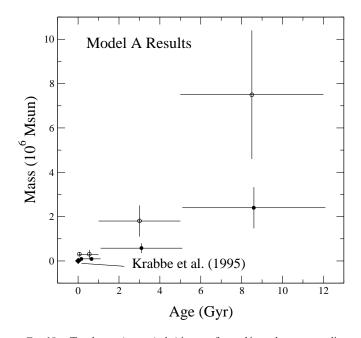


FIG. 19.—Total mass (*open circles*) in stars formed in each corresponding age bin (10–100 Myr, 100 Myr to 1 Gyr, 1–5 Gyr, and 5–12 Gyr) for model A of Fig. 17 (see also Table 7). The Krabbe et al. (1995) result for a ~5 Myr old burst, which produces  $3200 M_{\odot}$ , is shown. The filled circles represent the present-day mass for the same star formation history accounting for the mass loss due to stellar winds. The total present-day mass is consistent with the dynamical mass determinations (see text). The filled circles have been shifted by 0.1 Gyr in the figure for clarity.

some of the same stars used by Lebofsky et al. (1982b) as luminous AGB stars, indicating that intermediate ages ( $\gtrsim 100$  Myr) were present as well. Genzel et al. (1994) discussed the SFH in terms of a constant SFR and noted that such a model produces too few young blue supergiants if it is adjusted to match the number of older latetype stars. Krabbe et al. (1995) modeled the most recent

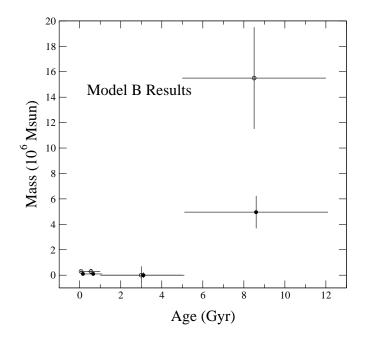


FIG. 20.—Same as Fig. 19, but for model B results

epoch of star formation in the nuclear cluster, producing the observed blue supergiants with a burst of  $10^{3.5} M_{\odot}$ over the last  $\sim$ 7 Myr. This point is shown in Figure 17. Apparently, the GC is currently in a more quiescent state than in the recent past. Davidge et al. (1997) argued for an old population to dominate the number counts within 6'' (0.23 pc) of the nuclear cluster on the basis of nearinfrared photometry. Philipp et al. (1999) used surface brightness fitting and photometry of individual stars to discuss the relative contributions of young and old stars over a larger ( $\sim 11' \times 11'$ ) area. Mezger et al. (1999) used the same data presented by Philipp et al. (1999) to further constrain the star formation history. Recently, van Loon et al. (2003) used photometric spectral energy distributions of a large sample of stars in the Galactic bulge to investigate the stellar populations there. Their results are qualitatively similar to those presented here, although they investigate a much larger area and have poor angular resolution in the central few parsecs. BSD96 used near-infrared spectra similar to those presented here to identify young, intermediate-age, and old stars. In this paper we build on this earlier work by computing the SFH directly from the observed stars using the  $M_{\rm bol}$  and  $T_{\rm eff}$  determined from the individual stellar spectra.

# 5.1. Stellar Remnants

Some fraction of the mass within the GC is due to massive stellar remnants (neutron stars or black holes from initially massive stars) that have migrated there from farther out (Morris 1993; Miralda-Escudé & Gould 2000) by dynamical friction with the low-mass stellar population. To the extent that some of these massive objects were formed outside the 2 pc radius covered by our data, that mass is unaccounted for in our SFH model. Morris (1993) has computed migration times as a function of mass and finds that a remnant of 10  $M_{\odot}$  could migrate in to the center from about 4 pc radius in a Hubble time. Miralda-Escudé & Gould (2000) find a similar result. The most massive remnants, black holes, will be able to migrate inward from the largest radius. Their large mass and the volume over which they are drawn mean that they will dominate this extra mass component (Morris 1993).

Morris (1993) explored a range of models and found that a total mass of remnants in the range 0.5 to  $8 \times 10^6 M_{\odot}$ could have migrated by dynamical friction to the central parsec (Morris's models for similar IMFs as used here produce masses  $\leq 5 \times 10^6 M_{\odot}$ ). The current estimate for the mass of the central black hole are within this range, which suggests that a significant fraction of the black hole mass could be made of stellar remnants or that a large number of remnants form a dark cluster surrounding the black hole, Miralda-Escudé & Gould (2000) estimate that 25,000 7  $M_{\odot}$ black hole remnants could have settled in the central parsec over a Hubble time; these remnants would still be in the stellar cluster because the depletion timescale for capture by the central black hole is 30 Gyr. Given the somewhat lower estimate of remnant mass due to migration by Miralda-Escudé & Gould (2000) and the estimate of  $\sim 2 \times 10^6 M_{\odot}$  in a luminous stellar cluster by Schödel et al. (2002) within a 2 pc radius of the center, we expect on order of at least a few times 10<sup>6</sup>  $M_{\odot}$  in stars to have formed in the inner 2 pc radius. If the Galaxy is 10 Gyr old, then this implies a time averaged SFR of  $\gtrsim 2 \times 10^{-4} M_{\odot} \text{ yr}^{-1}$ .

# 5.2. Depletion of CO Absorbers in the Central Parsec

Sellgren et al. (1990) showed that the CO absorption strength seen through apertures centered on unresolved starlight was weaker inside a radius of 15" (0.58 pc) than outside this radius (see also Genzel et al. 1996). We now know the weakness of the CO feature is due in part to dilution by very young stars (Genzel et al. 1996, for example). Genzel et al. (1996) also showed that the brightest resolved sources with strong CO (i.e., luminous AGB stars or supergiants) were absent from the inner 5''(0.19 pc). Both Sellgren et al. (1990) and Genzel et al. (1996) concluded it was possible that the atmospheres of such stars might be destroyed by collisions with lower mass stars in the GC, leading to a deficiency of COstrong stars. Bailey & Davies (1999) made detailed calculations of collision probabilities in the GC and concluded that collisions between giants and lower mass stars was unlikely to explain the putative missing stars because such collisions (which form a common envelope system) were ultimately ineffective in expelling the giant's envelope on a timescale shorter than the evolutionary timescale. However, Davies et al. (1998) did find that collisions between giants and binaries might be effective in removing the giant atmospheres in a short enough time to be observable (again through the development of common-envelope systems). In any case, the deficiency of resolved AGB or M supergiant stars is actually well concentrated to the center (Genzel et al. 1996,  $R \leq 5'' = 0.2$ pc) and should thus not affect the SFH estimates for the larger area studied in this work  $(R \sim 2')$ .

A consequence of the migration of massive remnants into the central parsec is the relaxation of the resulting dark cluster with the lower mass stars that exist there. The model of Miralda-Escudé & Gould (2000) predicts that stars older than a few gigayears will be pushed to larger radii, forming a distribution with a larger core radius (1-2 pc) and lower core density than they would otherwise have. The implication of this prediction is that the present data set would not be sensitive to the oldest epochs of star formation in the GC if a significant fraction of the  $\sim 1~M_{\odot}$  tracers have been removed from the inner 2 pc, and so it might underestimate the total SFH. On the other hand, if the extent of the dynamical redistribution of the low-mass stars is toward the low end of the range predicted by Miralda-Escudé & Gould (2000; core radius  $\sim 1$  pc), then we would not expect this to be a large effect. The total amount of mass deduced from dynamical models also constrains the SFH, and we discuss below the possibility that this constraint coupled with our models may provide evidence that the dynamical friction effects are seen in the GC.

# 5.3. The Star Formation History

Figure 17 indicates that there is significant ongoing star formation in the central few parsecs but that the bulk of stars (roughly 75% by mass) formed at earlier times (Fig. 19). This is in agreement with earlier work based on near-infrared number counts (Genzel et al. 1994; Mezger et al. 1999), and for the range of parameters discussed in § 3 and listed in Table 7, this conclusion holds. For the oldest stars, we are sampling just the very tip of the AGB; hence to observe any stars in such a short-lived phase requires a large mass to have originally formed. The details change by roughly a factor of 2, depending on which model SFH is chosen. The goodness-of-fit criterion,  $P_{\lambda}$ , suggests the uniform metallicity case (model A) is preferred. If true, it could suggest that the nucleus formed largely from enriched material produced in the early formation of the bulge. The purely solar [Fe/H] is also consistent with the narrow distribution of [Fe/H] from high-resolution spectra (Ramírez et al. 2000). Although the range of ages considered by Ramírez et al. (2000) is not as large as the data set presented here, the high-resolution data sample stars with ages up to ~5 Gyr (Fig. 13). In § 3 we noted that the current data set may not be sensitive to metal poor populations older than 12 Gyr.

The total mass represented by the SFH in Figure 17 is shown in Figure 19 and is  $9.9 \pm 3.0 \times 10^6 M_{\odot}$  for model A. This is about 3–6 times larger than the most detailed models (not counting the central black hole mass).

However, mass loss during the lifetime of stars from about 1 to 120  $M_{\odot}$  will reduce the cumulative final mass in the cluster. If all stars with 120  $M_{\odot} > M > 1 M_{\odot}$  are taken to have their remnant mass at the present time, then, using the mapping of initial mass to final mass given by Morris (1993), the present mass in stars is reduced by about 68%(i.e., we infer the present mass in stars remaining in the cluster to be 0.32 times the total mass formed over all times) to a total of  $3.2 \pm 1 \times 10^6 M_{\odot}$ , which is consistent with the dynamical models cited in § 4.2.1. The total present-day mass from the SFH depicted in Figure 17, but including mass loss, is shown in Figure 19. We have implicitly assumed that all the tracers of the SFH lie within a true radius of 2 pc, but they are actually distributed in a projected radius of 2 pc. Given the steepness of the stellar cluster radial density distribution ( $\leq 0.5$  pc core radius), the overestimate is likely to be small. The mass lost through stellar winds could be expelled from the region and/or recycled into new generations of stars. This is the maximal mass loss since not all stars are yet old enough to have reached their final mass, although most of the star formation has occurred at earlier times.

Morris (1993) argued that the IMF in the CMZ should be slanted toward higher masses than the Salpeter (1955) mass function. Figer et al. (1999) derived a flatter mass function in the nuclear young cluster, the Arches, 30 pc from the GC. However, it is not clear if this is representative of the initial mass function or dynamical effects (Portegies Zwart et al. 2002). An IMF slanted toward higher masses increases the total mass derived here since the IMF must still produce the same number of low-mass stars that form the majority in Figure 13. For the same low-mass cutoff of 0.7  $M_{\odot}$ , the flatter IMF model yields a total mass of  $16.4 \pm 5.2 \times 10^6 M_{\odot}$ . This model loses more mass through stellar winds, resulting in a presentday mass of  $3.9 \pm 1.2 \times 10^6 M_{\odot}$ . This is not too different than model A discussed above. Higher values of the lower mass cutoff can reduce the total mass formed and thus the present-day mass, but the cutoff is already near the limit of the mass of stars that we observe in the H-R diagram (about 1  $M_{\odot}$ ). Apart from this, our models are not very sensitive to the details of the IMF.

Our derived SFH may provide evidence that a significant number of low-mass stars have been removed from the central few parsecs, as suggested by Miralda-Escudé & Gould (2000). Our models, which trace the initial mass formed, require that we cut off the lower end of the mass function (below 0.7  $M_{\odot}$ ) in order to produce a final mass that is consistent with the dynamical mass in central few parsecs. It is possible that this mass is actually formed but subsequently removed by dynamical friction as massive remnants migrate toward the center (Morris 1993; Miralda-Escudé & Gould 2000). A better test of the dynamical relaxation effect predicted by Miralda-Escudé & Gould (2000), as they point out, is to compute the radial distribution of the many fainter low-mass stars. Alternately, the cutoff at 0.7  $M_{\odot}$  may represent a bias to higher masses forming in the GC, as argued by Morris (1993) and Figer et al. (1999).

Our technique of deriving  $M_{\rm bol}$  and  $T_{\rm eff}$  from lowresolution spectra is preferred over broadband photometric analyses because of the large scatter in  $A_K$  in the GC (BSD96) and the large variation in intrinsic colors of the giant stars (which can be corrected for with the spectroscopically derived  $T_{\rm eff}$ ). Newly commissioned multiobject infrared spectrometers on large telescopes are now available and will be used in the near future to make these observations. Obtaining such observations to 1.5–2 mag deeper than present will also improve the SFH calculation (e.g., Fig. 15) since the lower parts of the older isochrones will be more densely populated. The observational errors on  $T_{\rm eff}$  and  $M_{\rm bol}$  still set a limit on the final age resolution, however.

Our sample includes a number of luminous stars that require a substantial amount of very recent star formation (10–100 Myr). The high SFR at later times is reflected in Figures 15 and 17. The relatively large number of luminous AGB and supergiant stars in this region of the H-R diagram requires significant recent star formation, perhaps in the form of one or more concentrated bursts. Another way to see this is by considering Figure 15. There is a relative paucity of luminous stars in the part of the H-R diagram covered by the youngest isochrones for a constant SFR compared with the cooler, less luminous AGB stars. Our models do not rule out high SFRs at early times (in concentrated bursts); the lower mean SFRs result from the wider age bins. The SFR in the youngest (and narrowest) age bin is similar to the in fall rate of gas into the 2 pc molecular ring (see § 1). This suggests that star formation might be relatively efficient in the GC ( $\sim$ 50%–100%), although the current SFR from Krabbe et al. (1995) is a factor of 10 lower than in our most recent bin.

In § 1 we noted that the CMZ was forming about 0.5  $M_{\odot}$  yr<sup>-1</sup>. If we consider the SFR in the CMZ and GC per unit area (taking a 200 pc radius disklike distribution for the CMZ), then the GC has formed stars at a prodigious rate over its history. The CMZ normalized SFR is  $4 \times 10^{-6} M_{\odot}$  yr<sup>-1</sup> pc<sup>-2</sup>. In the GC, the average SFR is  $\sim 8 \times 10^{-4} M_{\odot}$  yr<sup>-1</sup>, and taking a radius of 2 pc gives  $6 \times 10^{-5} M_{\odot}$  yr<sup>-1</sup> pc<sup>-2</sup> with a peak of  $2.6 \times 10^{-4} M_{\odot}$  yr<sup>-1</sup> pc<sup>-2</sup>.

Sjouwerman et al. (1999) have used luminous OH/IR (maser sources) stars as tracers of star formation on larger scales (up to 50 pc) in the GC region. They find evidence of significant star formation at an epoch  $\gtrsim 1$  Gyr ago. Narayanan, Gould, & DePoy (1996) and Wood, Habing, & McGregor (1998) also find evidence for massive AGB, hence intermediate-age, stars, indicating significant star formation on a similar timescale and spatial scale. Frogel, Tiede, & Kuchinski (1999) studied the *K*-band number counts of stars in the inner Galaxy and found that the younger (i.e., intermediate-age) population may extend out to a degree from the nucleus. There is a feature in the H-R diagram of Figure 13 that coincides with this age range. It is the relatively dense group of stars centered at  $T_{\rm eff} = 3300$ K,  $M_{\rm bol} = -5.0$ . This feature represents significant intermediate-age star formation ( $\sim 1-2$  Gyr), although not a significant fraction of the total mass. The number of stars that trace out tracks near 1 Gyr in age is suggestive of a true "burst," although our models do not have the time resolution to conclusively limit the duration of this star formation activity. Van Loon et al. (2003) have discussed the properties of the stellar populations in the inner Galaxy using broadband photometric indexes from the DENIS and ISOGAL surveys. These surveys, whose angular resolution is more appropriate for studies on large scales, give results that are broadly consistent with those presented here. In particular, they find that the bulk of stars in the inner Galaxy are old and not metal-poor, that there has been significant star formation at intermediate ages, and that current SFRs are relatively high. The correspondence of star formation tracers in the central parsecs and on larger scales suggests that star formation in the GC may be influenced by processes in the inner Galaxy at large. The supply of gas to the central few parsecs may be linked to the stellar bar, which is thought to be a mechanism to funnel star-forming material to the inner Galaxy (see § 1).

## 6. SUMMARY

We have presented an H-R diagram for a sample of 79 cool and luminous M-type stars in the central few parsecs of the Galaxy. The sample is based on a magnitude-limited *K*-band data set presented by Blum et al. (1996a). The  $T_{\rm eff}$  and  $M_{\rm bol}$  were derived from CO and H<sub>2</sub>O molecular absorption features in  $\lambda/\Delta\lambda \sim 550$ –1200 *H*- and *K*-band spectra.

The H-R diagram was used to derive the star formation history for the GC. Our sample of stars is too small to independently constrain all the parameters in a detailed SFH (e.g., the slope of the initial mass function and the chemical enrichment history), and thus our SFH is not strictly unique. However, we find that the bulk of stars in the GC formed at early times ( $\gtrsim 5$  Gyr ago) for a range of model parameters. There is also evidence for significant recent star formation ( $\leq$  few gigayears ago). Such recent star formation activity coincides in time with evidence from other evolved stars at larger radii in the inner Galaxy (>50 pc) and suggests a connection between star formation in the central parsec and on larger scales (presumably through gas input to the region). The age resolution of our sample is not great because of observational errors on the derived  $M_{\rm bol}$ and  $T_{\rm eff}$ , the fact that the oldest isochrones are not well separated along the asymptotic giant branch, and the relatively small number of old, luminous stars, which trace the majority of the derived mass. Our best-fitting models require a cutoff in the IMF below a solar mass (at  $\sim 0.7 M_{\odot}$ ) in order to produce a present-day mass in the central few parsecs, which is consistent with existing enclosed masses derived from dynamical models. This "cutoff" might be evidence that mass-segregation effects are at work in the GC, as has been predicted previously, or might instead point to a bias toward high-mass star formation. Finally, we find better fits to the data with models that have [Fe/H] = 0.0 at all ages. This is consistent with earlier work at high spectral resolution that showed that stars between 10 Myr and  $\sim$ 5 Gyr in the GC have solar metallicity.

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REFERENCES

- Allen, D. A., Hyland, A. R., & Hillier, D. J. 1990, MNRAS, 244, 706 Bailey, V. C., & Davies, M. B. 1999, MNRAS, 308, 257 Barthes, D. 1998, A&A, 333, 647

- Becklin, E. E., & Neugebauer, G. 1975, ApJ, 200, L71 Bedding, T. R., & Zijlstra, A. A. 1998, ApJ, 506, L47
- Benson, P. J., & Little-Marenin, I. R. 1996, ApJS, 106, 579
- Bertelli, G., Bressan, A., Chiosi, C., Fagotto, F., & Nasi, E. 1994, A&AS,
- 106, 275
- Binney, J., Gerhard, O. E., Stark, A. A., Bally, J., & Uchida, K. I. 1991, MNRAS, 252, 210

- Blitz, L., & Spergel, D. N. 1991, ApJ, 379, 631 Carr, J. S., Sellgren, K., & Balachandran, S. C. 2000, ApJ, 530, 307
- Davidge, T. J., Simons, D. A., Rigaut, F., Doyon, R., & Crampton, D.
- 1997, AJ, 114, 2586 Davies, M. B., Blackwell, R., Bailey, V. C., & Sigurdsson, S. 1998, MNRAS, 301, 745
- DePoy, D. L., Atwood, B., Byard, P., Frogel, J. A., & O'Brien, T. 1993, Proc SPIE, 1946, 667
- DePoy, D. L., Gregory, B., Elias, J., Montane, A., Pérez, G., & Smith, R. 1990, PASP, 102, 1433 Dolphin, A. E. 1997, NewA, 2, 397
- . 2002, MNRAS, 332, 91
- Dumm, T., & Schild, H. 1998, NewA, 3, 137
- Dwek, E., et al. 1995, ApJ, 445, 716 Dyck, H. M., Benson, J. A., van Belle, G. T., & Ridgway, S. T. 1996, AJ, 111, 1705
- Dyck, H. M., van Belle, G. T., & Thompson, R. R. 1998, AJ, 116, 981
- Elias, J. H., Frogel, J. A., & Humphreys, R. M. 1985, ApJS, 57, 91

- Emas, J. H., Froget, J. A., & Humphreys, K. M. 1985, ApJS, 57, 91
  Fernandez-Villacanas, J. L., Rego, M., & Cornide, M. 1990, AJ, 99, 1961
  Figer, D. F., Kim, S. S., Morris, M., Serabyn, E., Rich, R. M., & McLean, I. S. 1999, ApJ, 525, 750
  Forrest, W. J., Shure, M. A., Pipher, J. L., & Woodward, C. E. 1987, in AIP Conf. Proc. 155, The Galactic Center, ed. C. H. Townes & D. C. Backer (New York: AIP), 153
  Frogel L A. Person S. F. Matthaws, K. & Aaronson M. 1078, ApJ
- Frogel, J. A., Persson, S. E., Matthews, K., & Aaronson, M. 1978, ApJ, 220, 75

- Frogel, J. A., Tiede, G. P., & Kuchinski, L. E. 1999, AJ, 117, 2296 Frogel, J. A., & Whitford, A. E. 1987, ApJ, 320, 199 Genzel, R., Hollenbach, D., & Townes, C. H. 1994, Rep. Prog. Phys., 57, 417
- Genzel, R., Pichon, C., Eckart, A., Gerhard, O. E., & Ott, T. 2000, MNRAS, 317, 348Genzel, R., Thatte, N., Krabbe, A., Kroker, H., & Tacconi-Garman, L. E.
- 1996, ApJ, 472, 153
- Ghez, A. M., Klein, B. L., Morris, M., & Becklin, E. E. 1998, ApJ, 509, 678 Ghez, A. M., et al. 2003, ApJ, 586, L127 Girardi, L., Bressan, A., Bertelli, G., & Chiosi, C. 2000, A&AS, 141, 371
- Güsten, R. 1989, in IAU Symp. 136, The Center of the Galaxy, ed.
- Gusten, R., 1939, in TAO Symp. 130, The Center of the Galaxy, ed.
  M. Morris (Dordrecht: Kluwer), 89
  Güsten, R., Genzel, R., Wright, M. C. H., Jaffe, D. T., Stutzki, J., & Harris, A. I. 1987, ApJ, 318, 124
  Haller, J. 1992, Ph.D. thesis, Univ. Arizona

- Haller, J. W., Rieke, M. J., Rieke, G. H., Tamblyn, P., Close, L., & Melia, F. 1996, ApJ, 456, 194

- F. 1996, ApJ, 456, 194
  Holtzman, J. A., et al. 1993, AJ, 106, 1826
  Humphreys, R. M. 1978, ApJS, 38, 309
  Jackson, J. M., Geis, N., Genzel, R., Harris, A. I., Madden, S., Poglitsch, A., Stacey, G. J., & Townes, C. H. 1993, ApJ, 402, 173
  Kent, S. M. 1992, ApJ, 387, 181
  Kleinmann, S. G., & Hall, D. N. B. 1986, ApJS, 62, 501
  Krakhe, A., Genzel, B., Dravata, S., & Patrain, M. 1001, A. J. 202, 110

- Krabbe, A., Genzel, R., Drapatz, S., & Rotaciuc, V. 1991, ApJ, 382, L19 Krabbe, A., et al. 1995, ApJ, 447, L95
- Lambert, D. L., Gustafsson, B., Eriksson, K., & Hinkle, K. H. 1986, ApJS, 62,373
- Lafon, J.-P. J., & Berruyer, N. 1991, A&A Rev., 2, 249
- Lançon, A., & Mouhcine, M. 2002, A&A, 393, 167
- Lançon, A., & Rocca-Volmerange, B. 1992, A&AS, 96, 593
- Lebofsky, M. J., Rieke, G. H., Deshpande, M. R., & Kemp, J. C. 1982a, ApJ, 263, 672
- Lebofsky, M. J., Rieke, G. H., & Tokunaga, A. T. 1982b, ApJ, 263, 736

- Lee, Y. 1992, AJ, 104, 1780 Libonate, S., Pipher, J. L., Forrest, W. J., & Ashby, M. L. N. 1995, ApJ, 439 202
- Liszt, H. S., & Burton, W. B. 1980, ApJ, 236, 779
- Lindqvist, M., Habing, H. J., & Winnberg, A. 1992, A&A, 259, 118 Little, S. J., & Price, S. D. 1985, AJ, 90, 1812
- Luck, R. E., & Bond, H. E. 1980, ApJ, 241, 218
- 1982a, ApJ, 263, 215 1982b, ApJ, 256, 177 1989, ApJS, 71, 559
- Massey, P., Johnson, K. E., & Degioia-Eastwood, K. 1995, ApJ, 454, 151 Mathis, J. S. 1990, ARA&A, 28, 37
- McGinn, M. T., Sellgren, K., Becklin, E. E., & Hall, D. N. B. 1989, ApJ, 338 824
- McWilliam, A., & Rich, R. M. 1994, ApJS, 91, 749 Mezger, P. G., Zylka, R., Philipp, S., & Launhardt, R. 1999, A&A, 348, 457

- <sup>457</sup>
  Mighell, K. 1999, ApJ, 518, 380
  Miralda-Escudé, J., & Gould, A. 2000, ApJ, 545, 847
  Moneti, A., Glass, I. S., & Moorwood, A. F. M. 1994, MNRAS, 268, 194
  Morris, M. 1993, ApJ, 408, 496
  Morris, M., & Serabyn, E. 1996, ARA&A, 34, 645
  Mulder, W. A., & Liem, B. T. 1986, A&A, 157, 148
  Najarro, F., Hillier, D. J., Kudritzki, R. P., Krabbe, A., Genzel, R., Lutz, D. Drapatz, S. & Geballe, T. R. 1994, A&A, 285, 573

- D., Drapatz, S., & Geballe, T. R. 1994, A&A, 285, 573 Najarro, F., Krabbe, A., Genzel, R., Lutz, D., Kudritzki, R. P., & Hillier, D. J. 1997, A&A, 325, 700 Narayanan, V. K., Gould, A., & DePoy, D. L. 1996, ApJ, 472, 183 Oliva, E., & Origlia, L. 1992, A&A, 254, 466 Olsen, K. A. G. 1999, AJ, 117, 2244

- Ott, T., Eckart, A., & Genzel, R. 1999, ApJ, 523, 248 Paczyński, B. 1970, Acta Astron., 20, 47

- Perryman, M. A. C., et al. 1997, A&A, 323, L49
   Philipp, S., Zylka, R., Mezger, P. G., Duschl, W. J., Herbst, T., & Tuffs, R. J. 1999, A&A, 348, 768
   Pogge, R. W., et al. 1998, Proc. SPIE, 3354, 414
   Portoria Zararté, S. E. McLing, L. McMillar, S. L. W., & Int, P. 2002
- Pogge, R. W., et al. 1998, Proc. SPIE, 3354, 414
  Portegies Zwart, S. F., Makino, J., McMillan, S. L. W., & Hut, P. 2002, ApJ, 565, 265
  Press, W. H., Teulkolsky, S. A., Vetterling, W. T., & Flannery, B. P., 1992, Numerical Recipes in C (2d ed.; Cambridge: Cambridge Univ. Press)
  Ramírez, S. V., DePoy, D. L., Frogel, J. A., Sellgren, K., & Blum, R. D. 1997, AJ, 113, 1411
  Ramírez, S. V., Sellgren, K., Carr, J. S., Balachandran, S. C., Blum, R., Terndrup, D. M., & Steed, A. 2000, ApJ, 537, 205

- Reid, M. J. 1993, ARA&A, 31, 345 Richichi, A., Ragland, S., Stecklum, B., & Leinert, C. 1998, A&A, 338, 527 Sadler, E. M., Rich, R. M., & Terndrup, D. M. 1996, AJ, 112, 171 Saha, P., Bicknell, G. V., & McGregor, P. J. 1996, ApJ, 467, 636

467,645

- Salpeter, E. E. 1955, ApJ, 121, 161 Sanders, R. H. 1999, in ASP Conf. Ser. 186, The Central Parsecs of the Galaxy, ed. H. Falcke (San Francisco: ASP), 250 Schödel, R., et al. 2002, Nature, 419, 694
- Sellgren, K., Hall, D. N. B., Kleinmann, S. G., & Scoville, N. Z. 1987, ApJ,
- 317,881 Sellgren, K., McGinn, M. T., Becklin, E. E., & Hall, D. N. 1990, ApJ, 359,
- 112 Sjouwerman, L. O., Habing, H. J., Lindqvist, M., van Langevelde, H. J., & Winnberg, A. 1999, in ASP Conf. Ser. 186, The Central Parsecs of the Galaxy, ed. H. Falcke (San Francisco: ASP), 379
  Smith, V. V., & Lambert, D. L. 1985, ApJ, 294, 326

Strecker, D. W., Erickson, E. F., & Witteborn, F. C. 1978, AJ, 83, 26
 Tamblyn, P., Rieke, G. H., Hanson, M. M., Close, L. M., McCarthy, D. W., & Rieke, M. J. 1996, ApJ, 456, 206

Tamura, M., Werner, M. W., Becklin, E. E., & Phinney, E. S. 1996, ApJ,

Tanner, A., et al. 2003, The Central 300 Parsecs of the Milky Way (Astron. Nachr. Suppl., Vol. 324, S1) Terndrup, D. M. 1988, AJ, 96, 884 Tiede, G. P., Frogel, J. A., & Terndrup, D. M. 1995, AJ, 110, 2788

Tollestrup, E. V., Becklin, E. E., & Capps, R. W. 1989, AJ, 98, 204

- van Belle, G. T., Dyck, H. M., Thompson, R. R., Benson, J. A., & Kannappan, S. J. 1997, AJ, 114, 2150
  van Leeuwen, F., Feast, M. W., Whitelock, P. A., & Yudin, B. 1997, MNRAS, 287, 955
  van Loon, J., et al. 2003, MNRAS, 338, 857

- Weiland, J. L., et al. 1994, ApJ, 425, L81
  Wood, P. R., Habing, H. J., & McGregor, P. J. 1998, A&A, 336, 925
  Zoccali, M., Cassisi, S., Frogel, J. A., Gould, A., Ortolani, S., Renzini, A., Rich, R. M., & Stephens, A. W. 2000, ApJ, 530, 418